

Topological properties of tilings

From structure to spectrum

Eric Akkermans



Virtual Workshop on new approaches to quasi-periodic
spectral and topological analysis, **May 2021**

Is there a relation between
structure and spectrum in aperiodic
tilings ?

A equivalent of Bloch theorem for
tilings

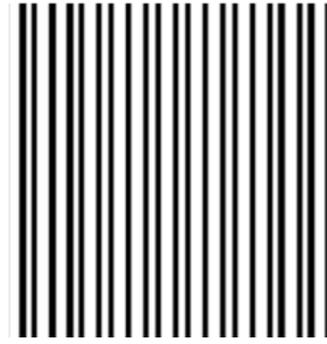
1. E. Akkermans, Y. Don, J. Rosenberg and C. L. Schochet, *Relating Diffraction and Spectral Data of Aperiodic Tilings: Towards a Bloch theorem*, *J. Geom. Phys.* **165**, 104217 (2021).

Outline

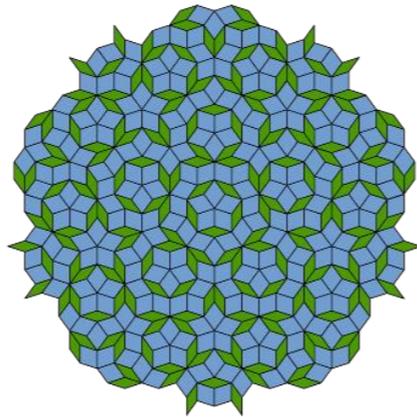
- 1 Prologue
- 2 Cut and Project Tilings and Windings
- 3 Substitution Tilings and Čech Cohomology
- 4 Bloch Theorem for Aperiodic Tilings
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in Fractals and Random Tilings
- 6 Epilogue

Tilings

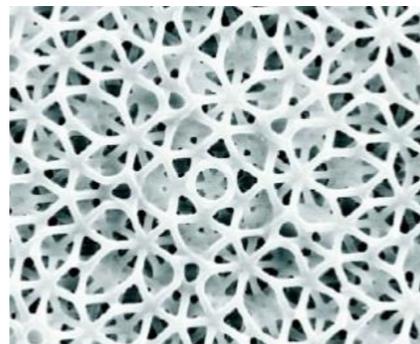
1D



2D

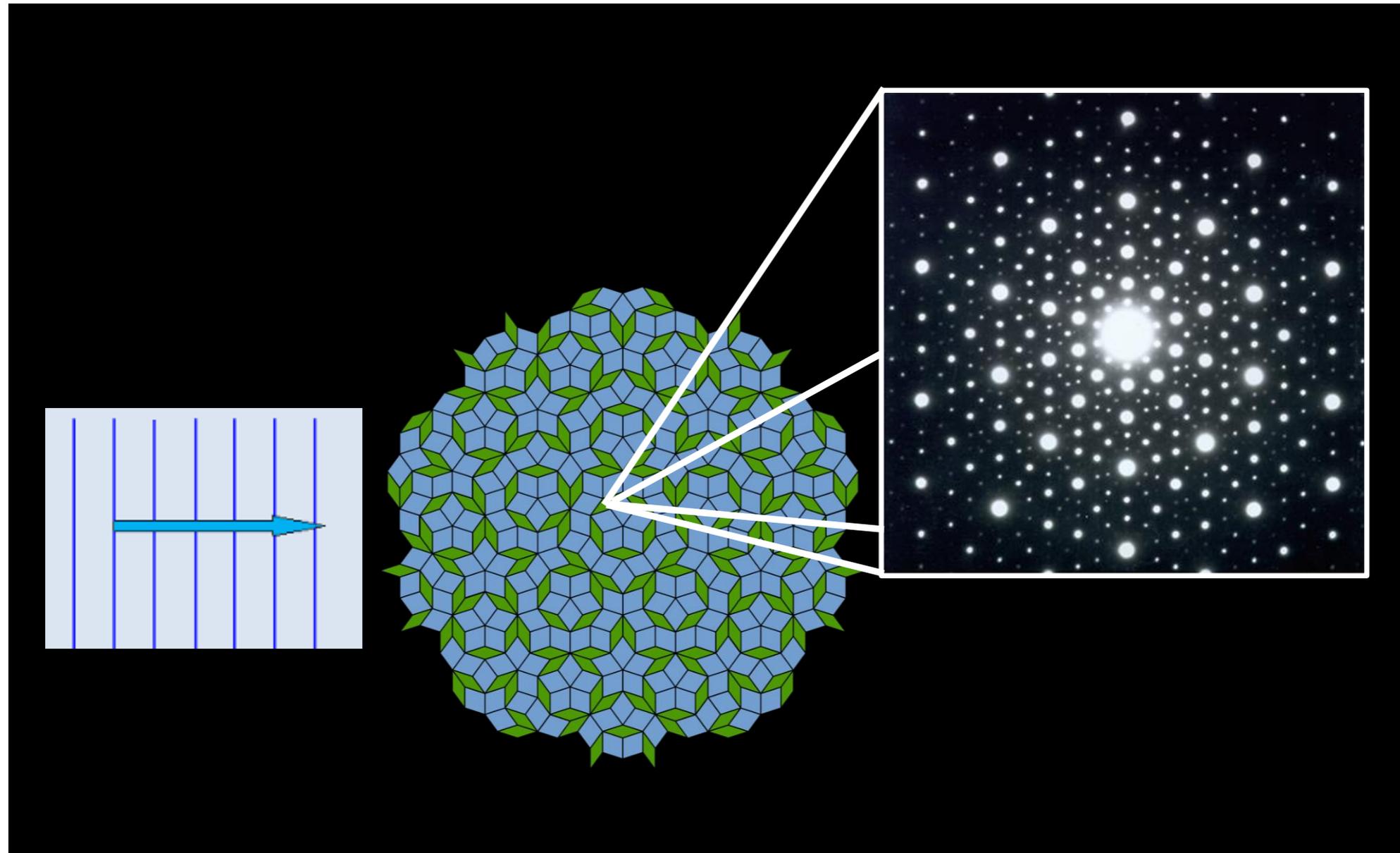


3D

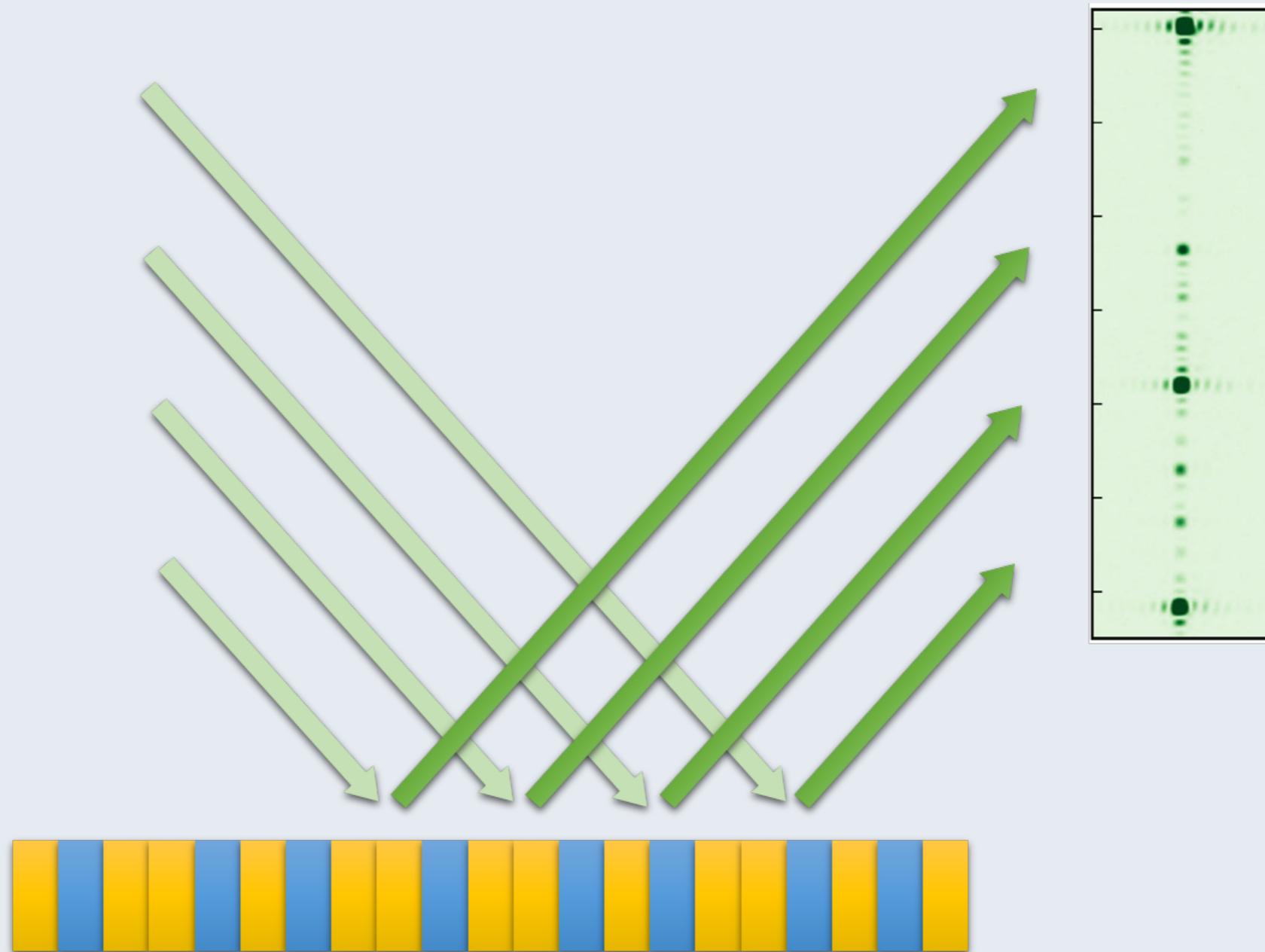


How to Characterize Tilings – Structure?

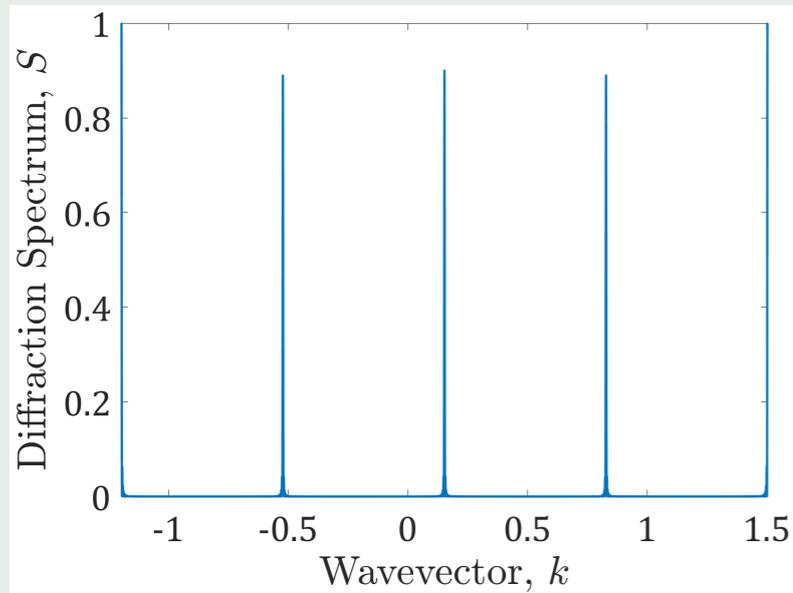
Diffraction pattern - Structure



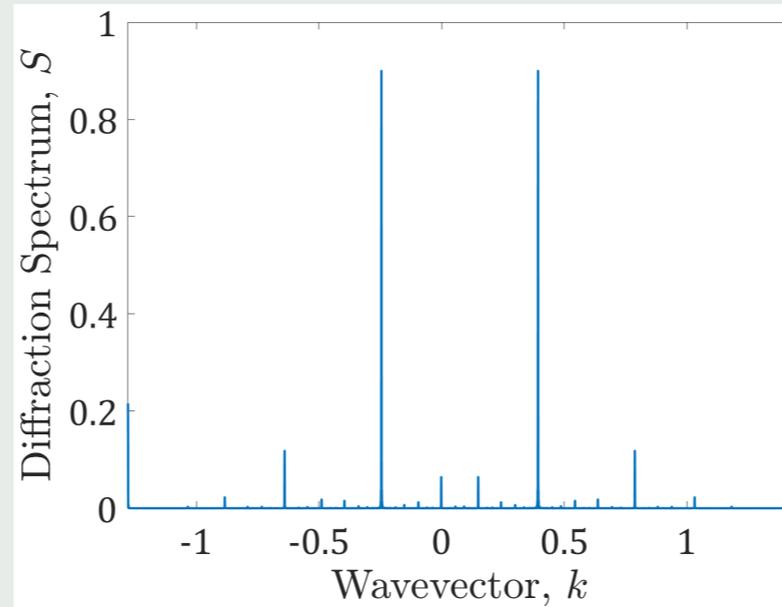
Diffraction (X-ray) pattern



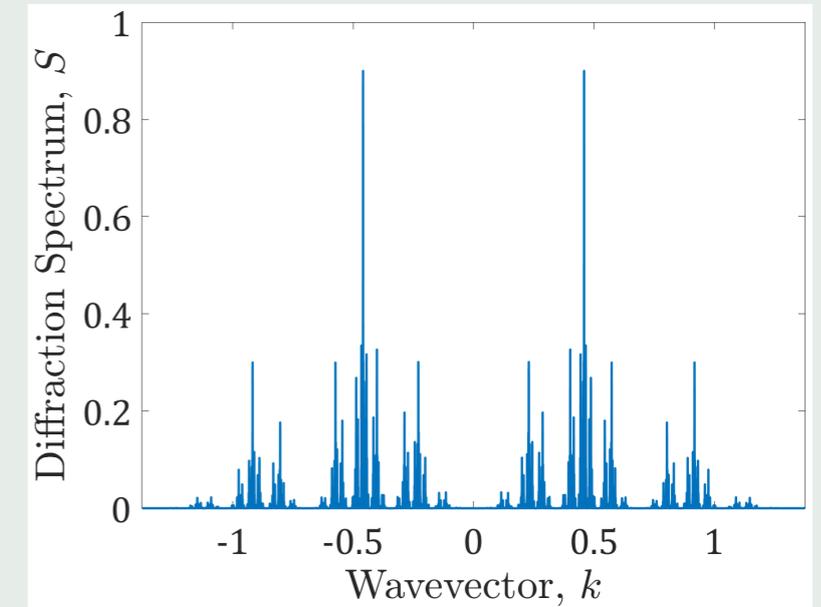
Periodic



Quasiperiodic



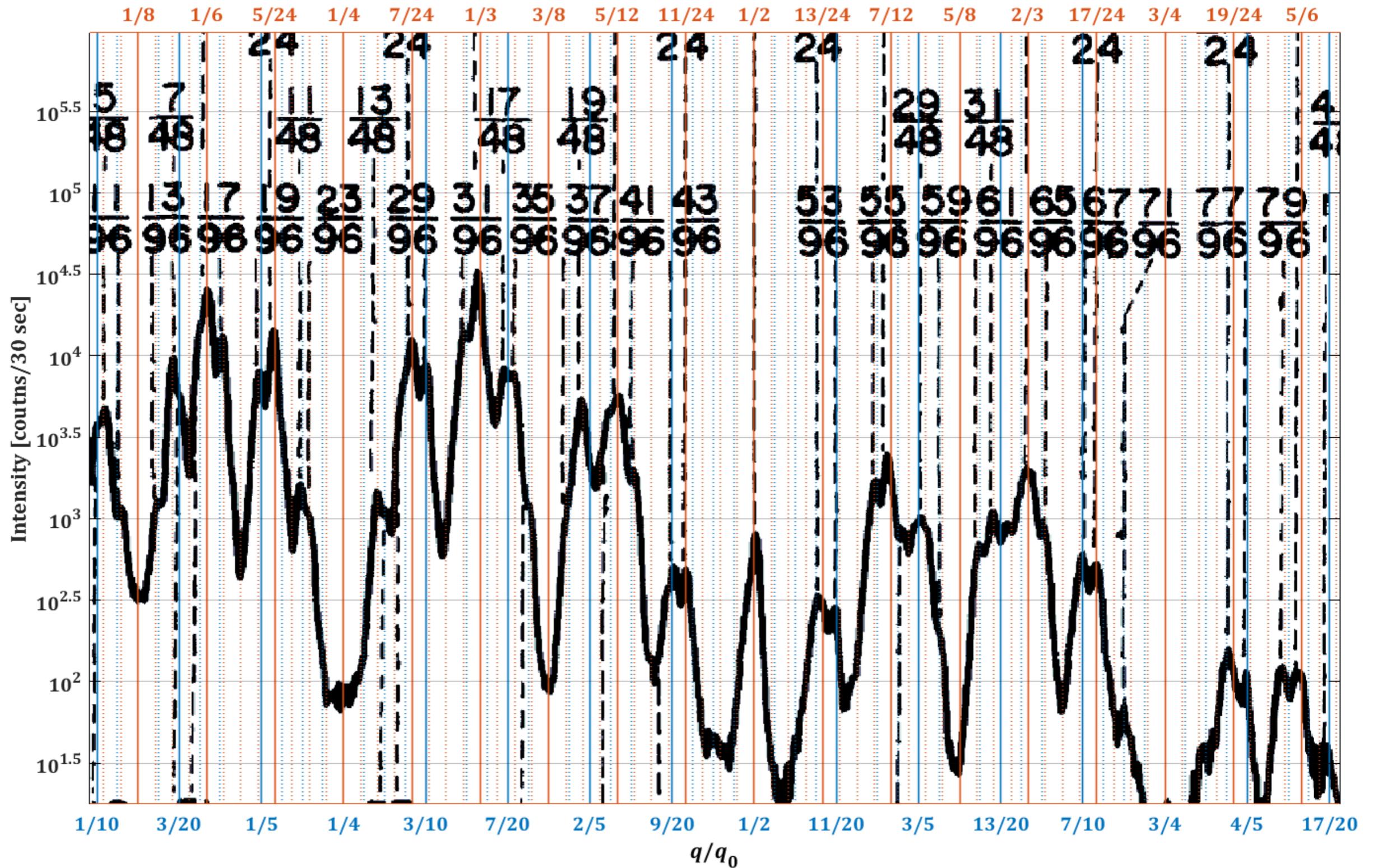
Aperiodic



Fibonacci

Thue-Morse

Existence of a Bragg peaks (PP) diffraction pattern is often unclear (e.g. Thue-Morse)



F. Axel and H. Terauchi, *Phys. Rev. Lett.* **66**, 2223–2226 (1991)

How to Characterize Tilings – Spectrum?

How to Characterize Tilings – Spectrum?

Spectrum & Integrated Density of States

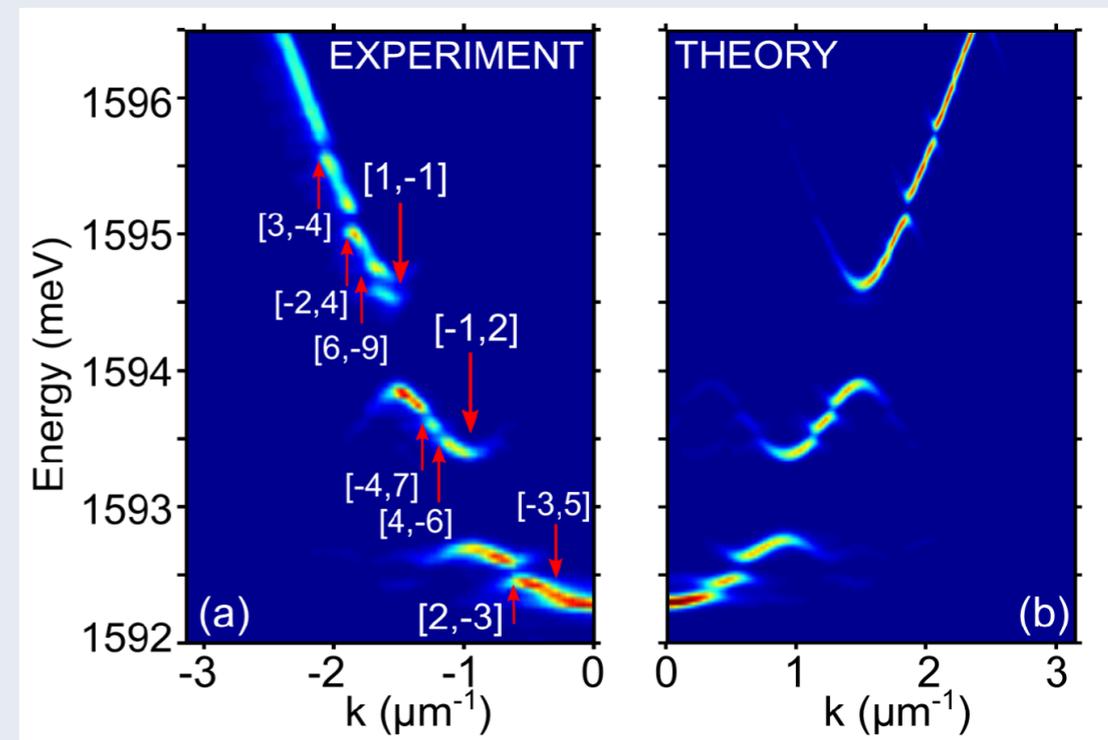
- Solve a Hamiltonian $H(E)$ (Fibonacci)

$$H\psi(x) = E\psi(x)$$



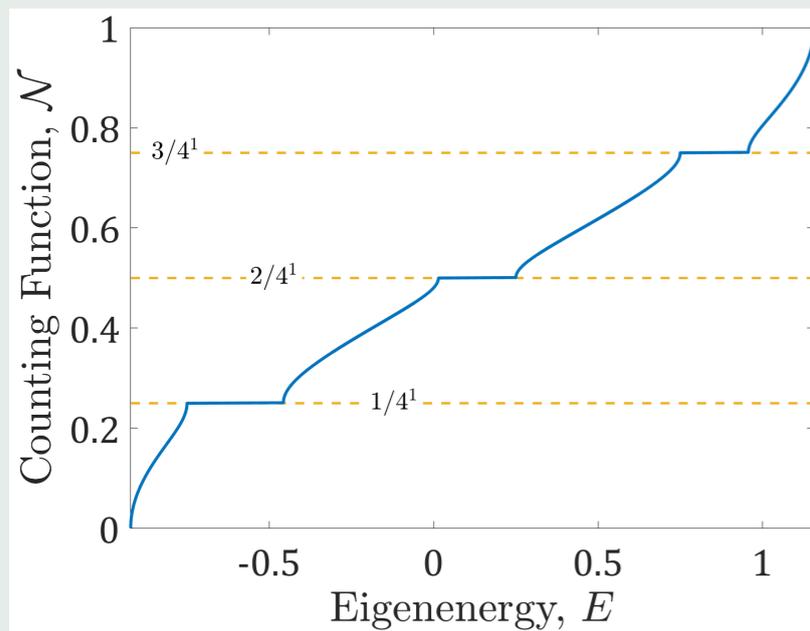
- Find the spectrum:
 - dispersion $E(k)$
 - integrated density of states

$$H(E) \rightarrow \begin{cases} \rho(E) & \text{DOS} \\ \mathcal{N}(E) & \text{IDOS} \end{cases}$$

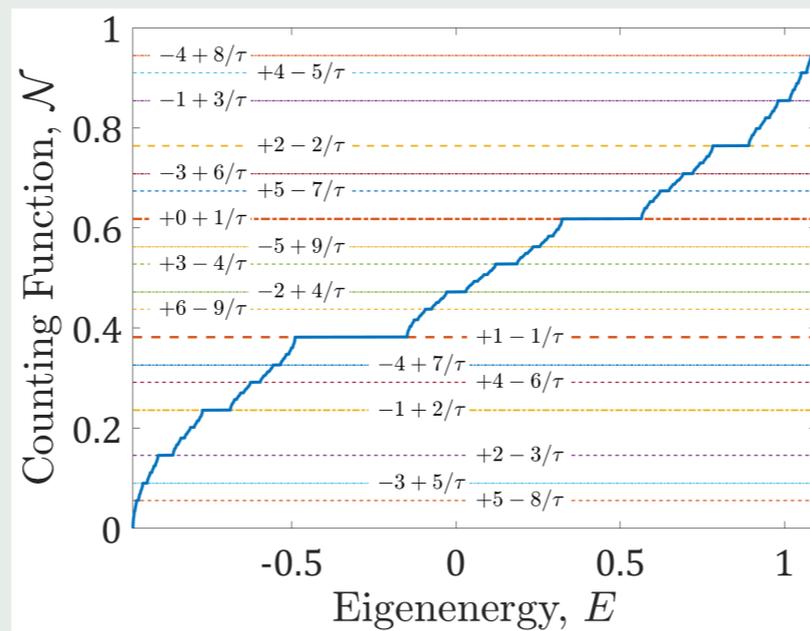


How to Characterize Tilings – Spectrum?

Periodic

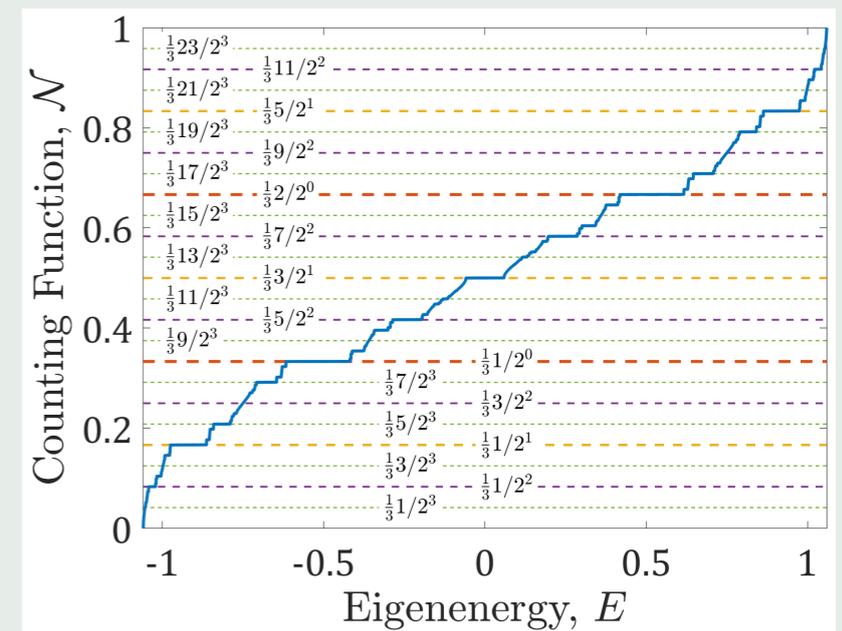


Quasiperiodic



Fibonacci

Aperiodic



Thue-Morse

Correspondence between Structure and Spectrum?

Bloch theorem

Periodic case: we know the connection

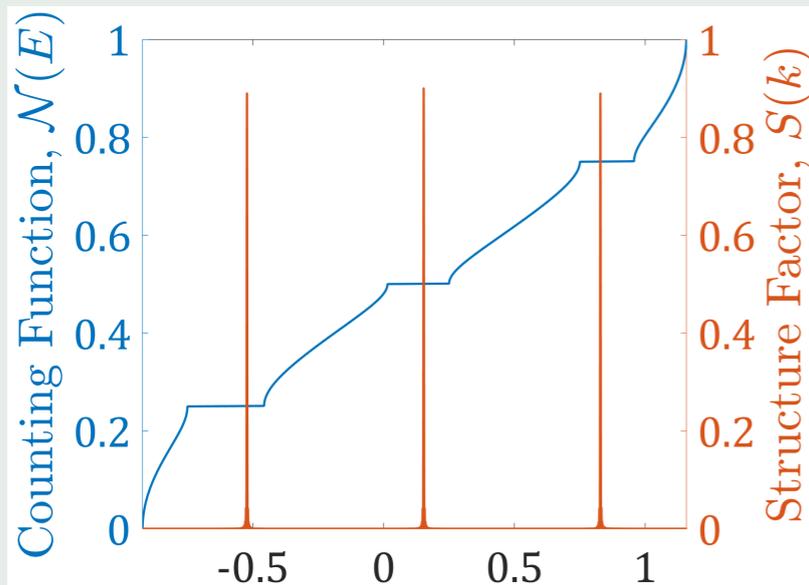
finite # of peaks $\xleftrightarrow[\text{correspondence}]{1 \text{ to } 1}$ finite # of gaps

Aperiodic case: this is not necessarily true

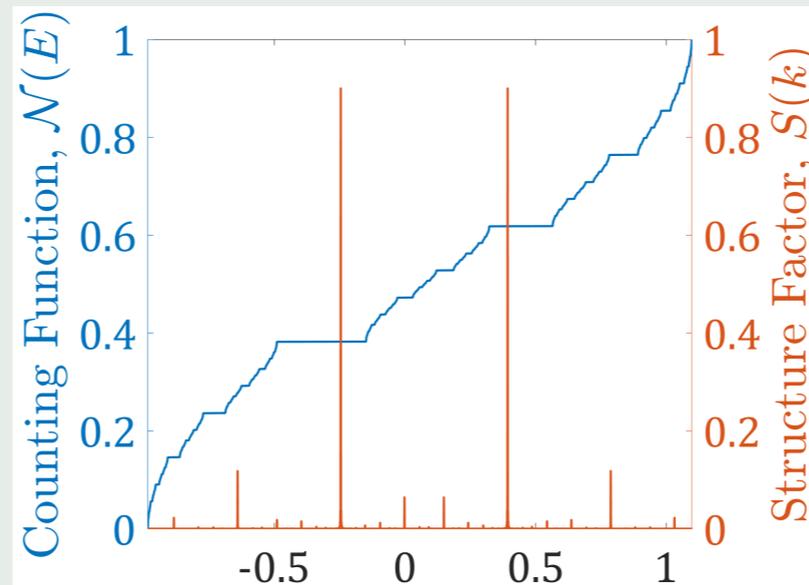
- We show that—at least for one family—there is a connection

Correspondence between Structure and Spectrum?

Periodic

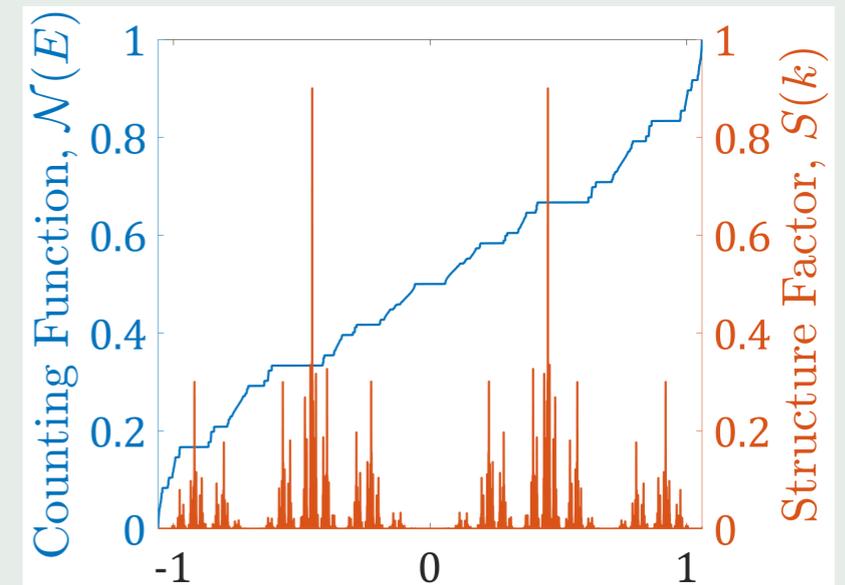


Quasiperiodic



Fibonacci

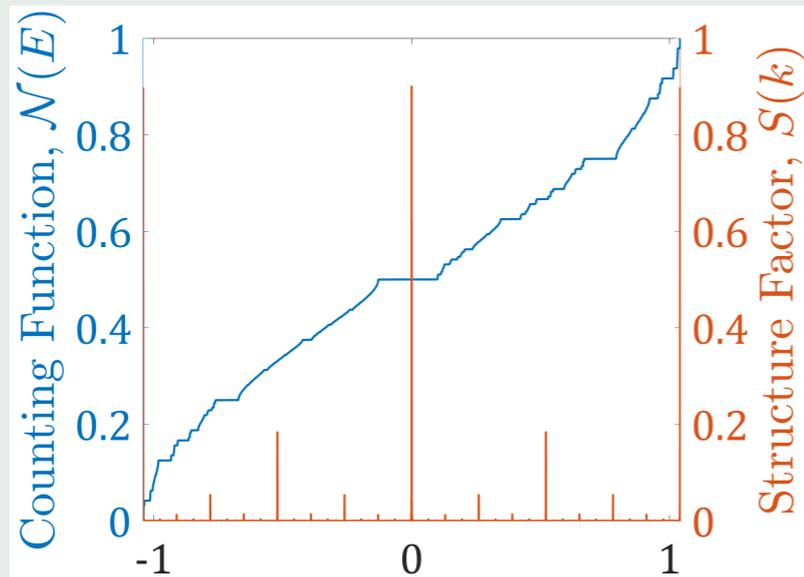
Aperiodic



Thue-Morse

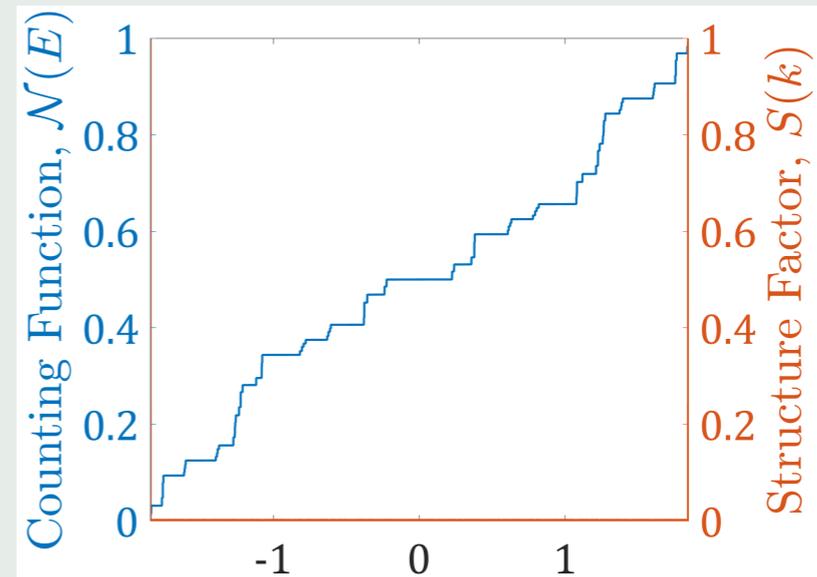
Correspondence between Structure and Spectrum?

Aperiodic



Period Doubling

Aperiodic



Rudin-Shapiro

Showing the connection - Finding the tools to discriminate between tilings

The tool: topological invariants

We use the Čech cohomology \check{H}^1

- to calculate Bragg peaks
- to compute topological numbers
- to show correspondence

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How to Construct Aperiodic Tilings?

Problem

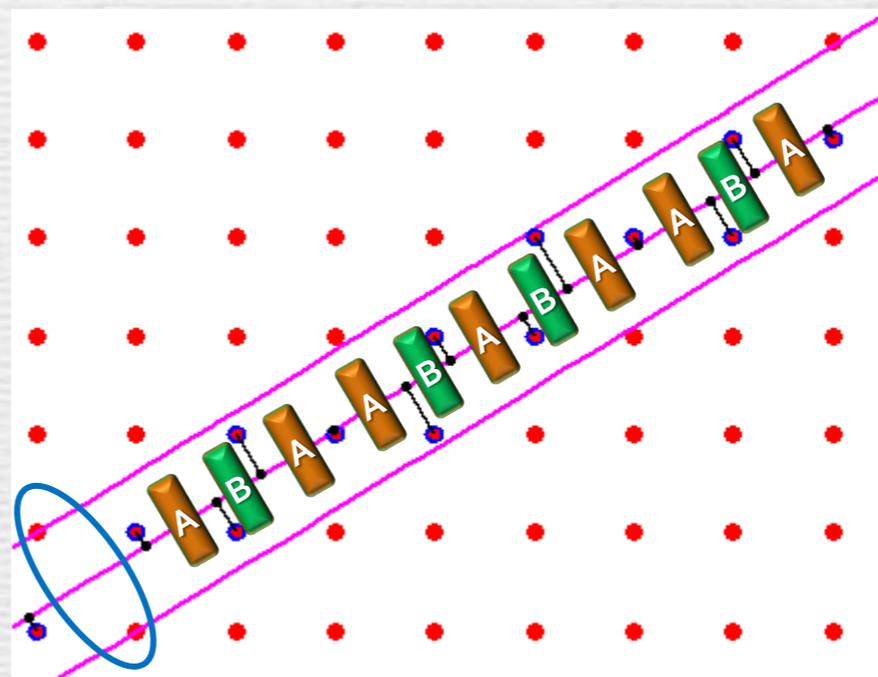
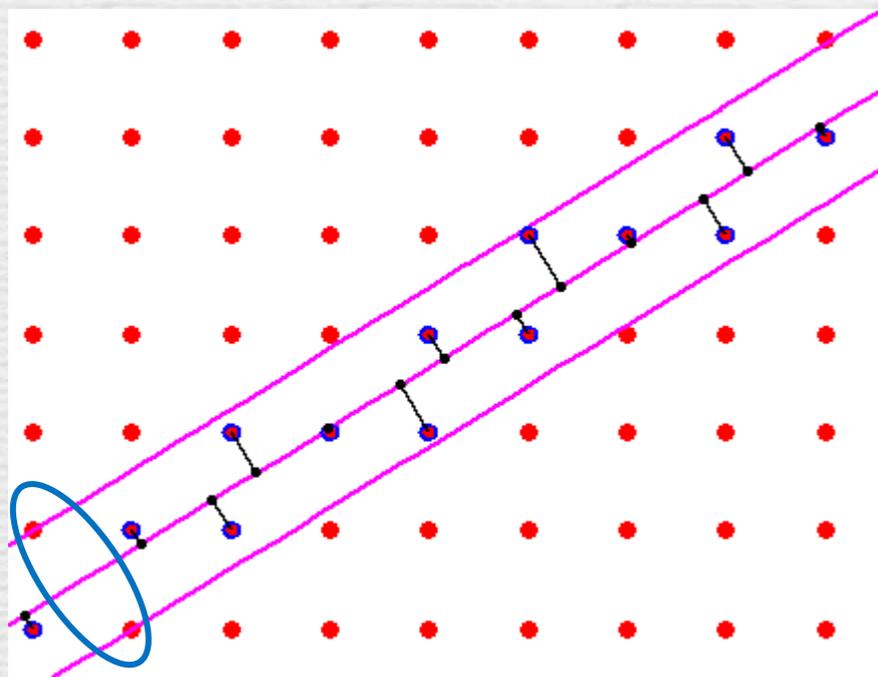
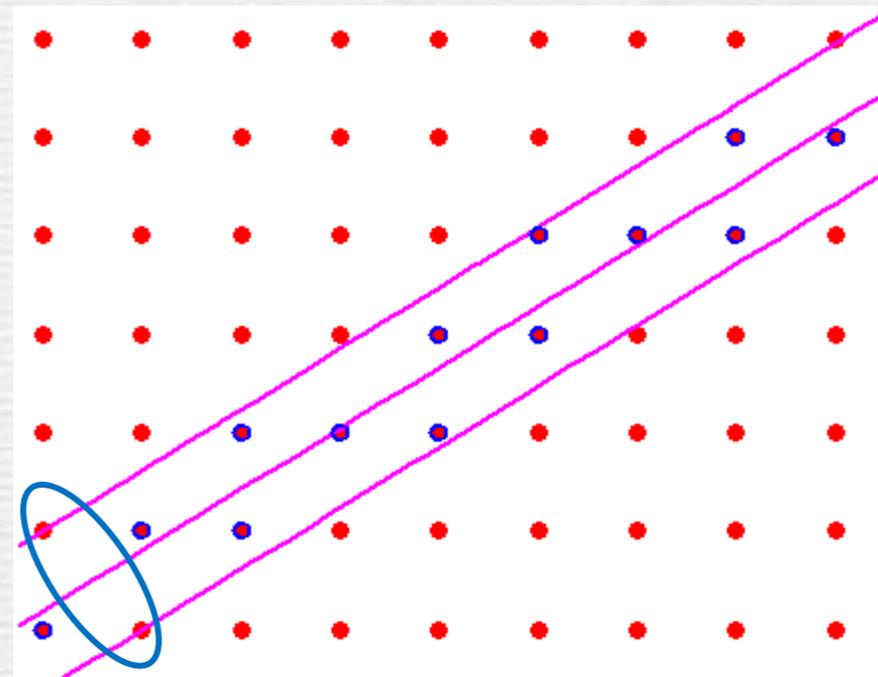
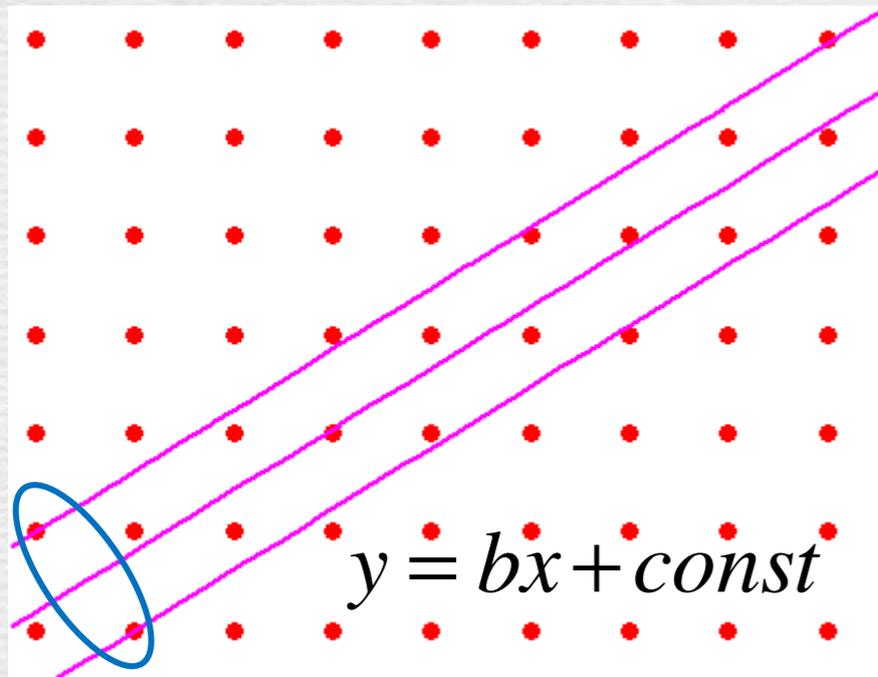
How to tile a space without repeating a pattern indefinitely?

Solution

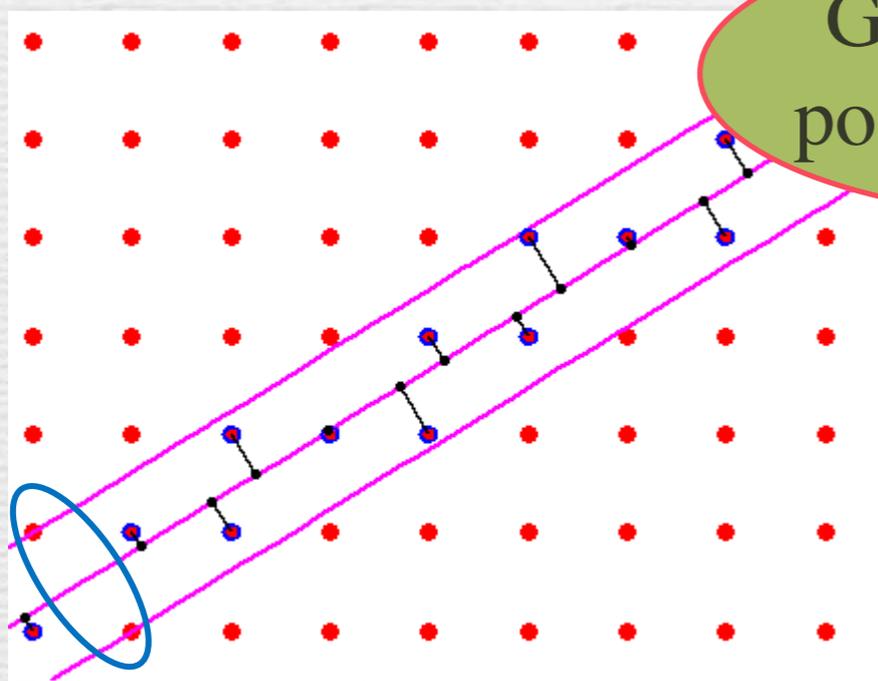
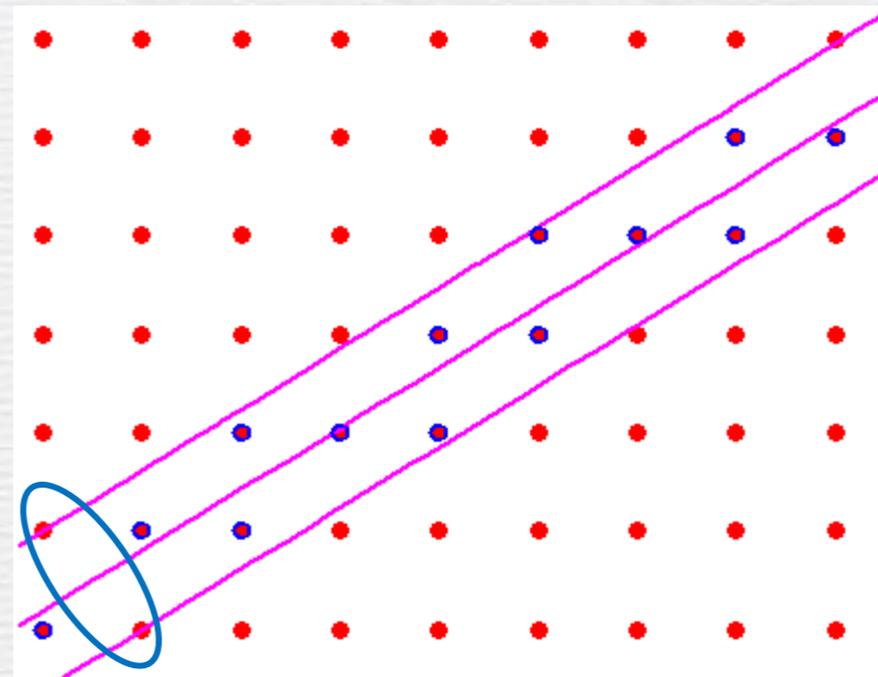
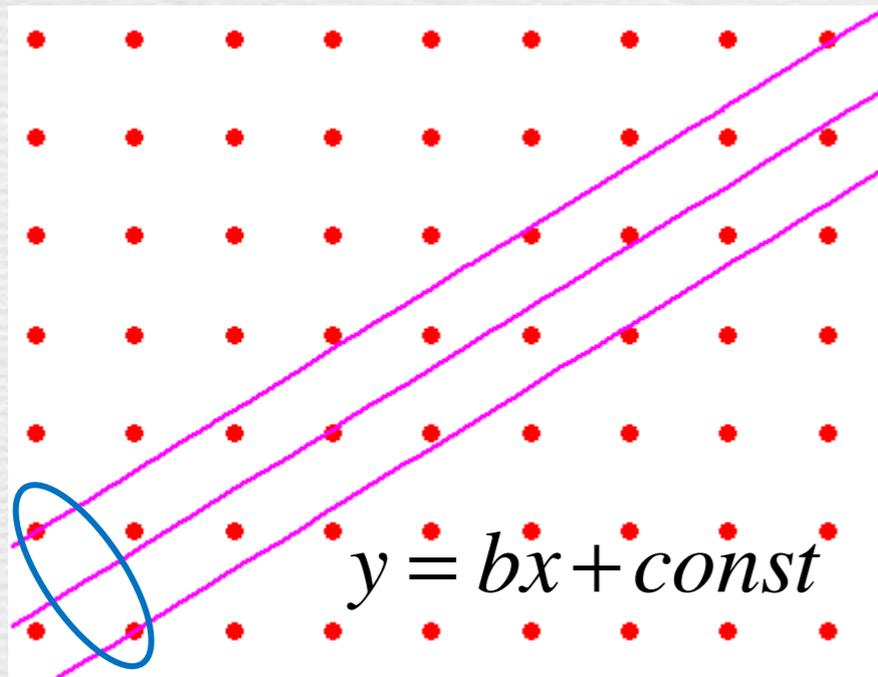
There are several methods

- We start with the Cut and Project (C&P) scheme

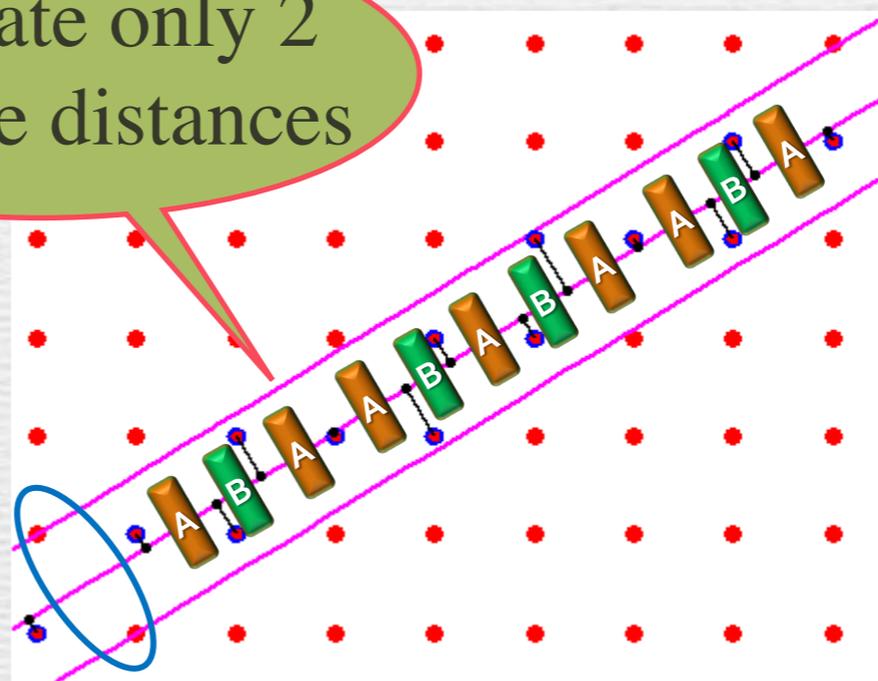
Start from a 2D lattice $L = \mathbb{Z}^2$



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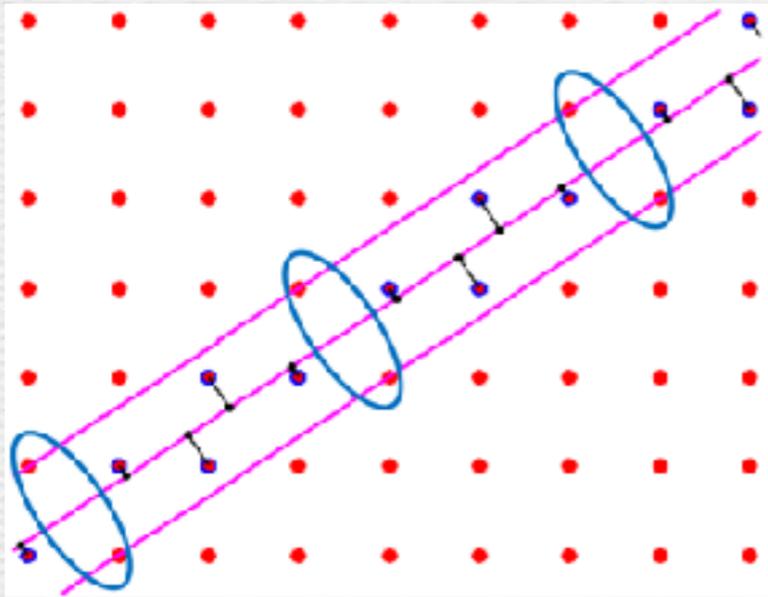


Generate only 2 possible distances



A B A A B A B A A B A ...

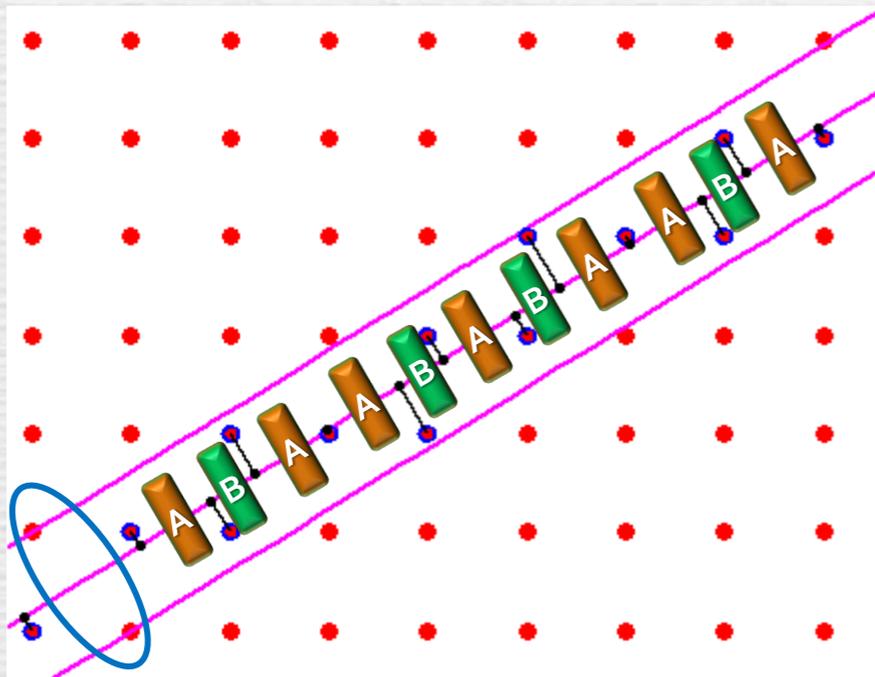
For a rational slope : periodic superlattice



$$y = \frac{2}{3}x + \text{const}$$



For an irrational slope : quasi-periodic structure



$$y = \tau^{-1}x + \text{const}$$



$$\tau = \frac{(1 + \sqrt{5})}{2}$$

golden mean

Different ways to build tiling chains

- Characteristic function
- Cut & Project

Characteristic function

$$\chi_n = \text{sign} \left[\cos(2\pi n \tau^{-1} + \phi) - \cos(\pi \tau^{-1}) \right]$$

$$\begin{aligned} -1 &= \text{B} \\ +1 &= \text{A} \end{aligned}$$

Characteristic function

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$$F_N(\phi) = [\chi_1 \chi_2 \cdots \chi_n \cdots \chi_N] \iff \text{A B A A B A B A A B A A B} \cdots$$

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The angle ϕ is a (legitimate) degree of freedom.

ϕ is known as a phason

$$\tau = \frac{\sqrt{5} + 1}{2} \approx 1.62$$

Characteristic function

$$\chi_n = \text{sign} \left[\cos(2\pi n \tau^{-1} + \phi) - \cos(\pi \tau^{-1}) \right]$$

ϕ is an innocuous and thus ignored modulation phase.

For an infinite Fibonacci chain :

$$\phi_\infty = 3\pi\sigma = 3\pi\tau^{-1}$$

Define instead

$$\chi_n = \text{sign} \left[\cos(2\pi n \tau^{-1} + \phi_\infty + \Delta\phi) - \cos(\pi \tau^{-1}) \right]$$

$$\tau = \frac{(1 + \sqrt{5})}{2}$$

golden mean

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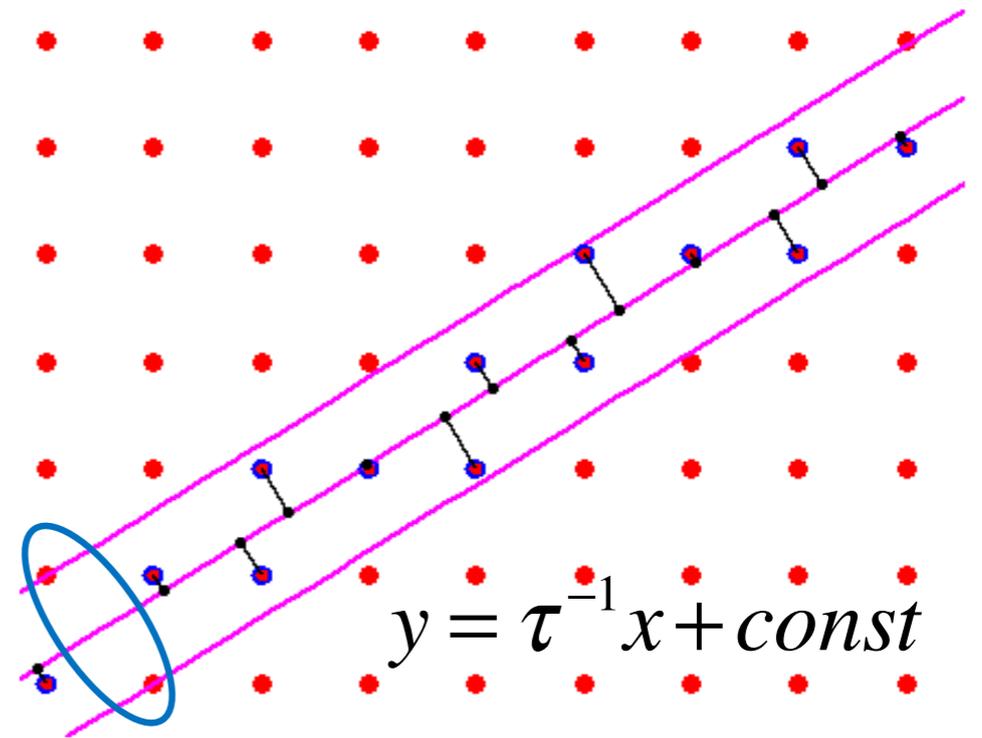
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golden mean

C&P method

Is it possible to give a meaning to $\Delta\phi$ using the C&P method ?



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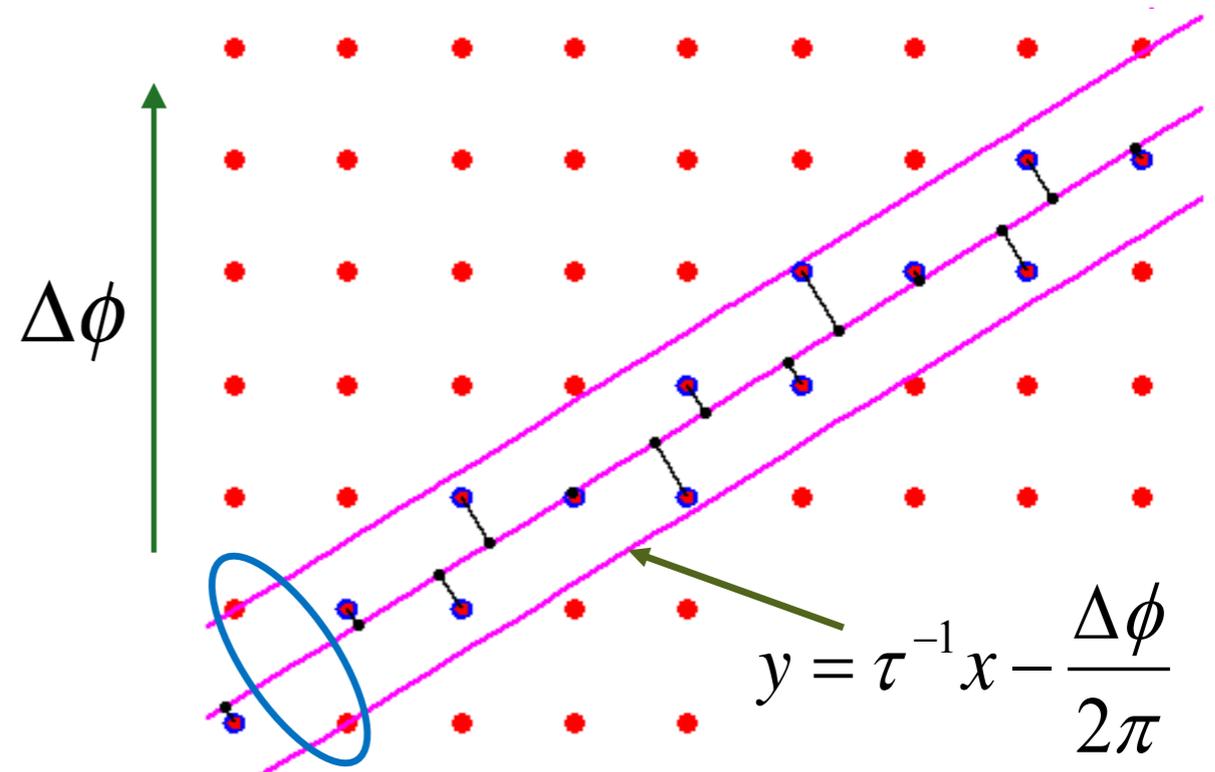
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golden mean

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Is it possible to give a meaning to $\Delta\phi$ using the C&P method ?



We understand the meaning of $\Delta\phi$

Meaning of the phason ϕ : a gauge field

A gauge degree of freedom

- Take a characteristic function

$$\chi(n, \phi) = \text{sign} \left[\cos \left(2\pi n \lambda_1^{-1} + \phi \right) - \cos \left(\pi \lambda_1^{-1} \right) \right]$$



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A finite segment of size N

A gauge degree of freedom

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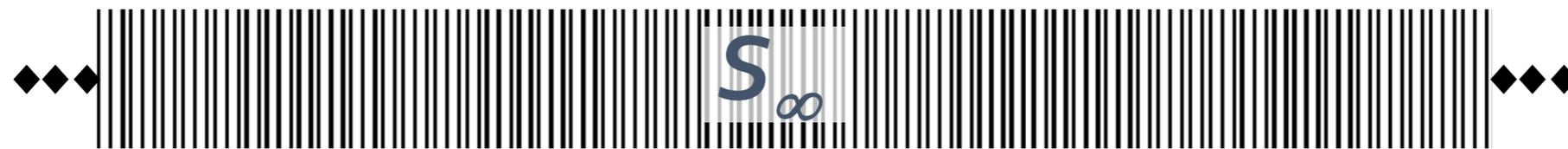


Discrete gauge: Choice of origin

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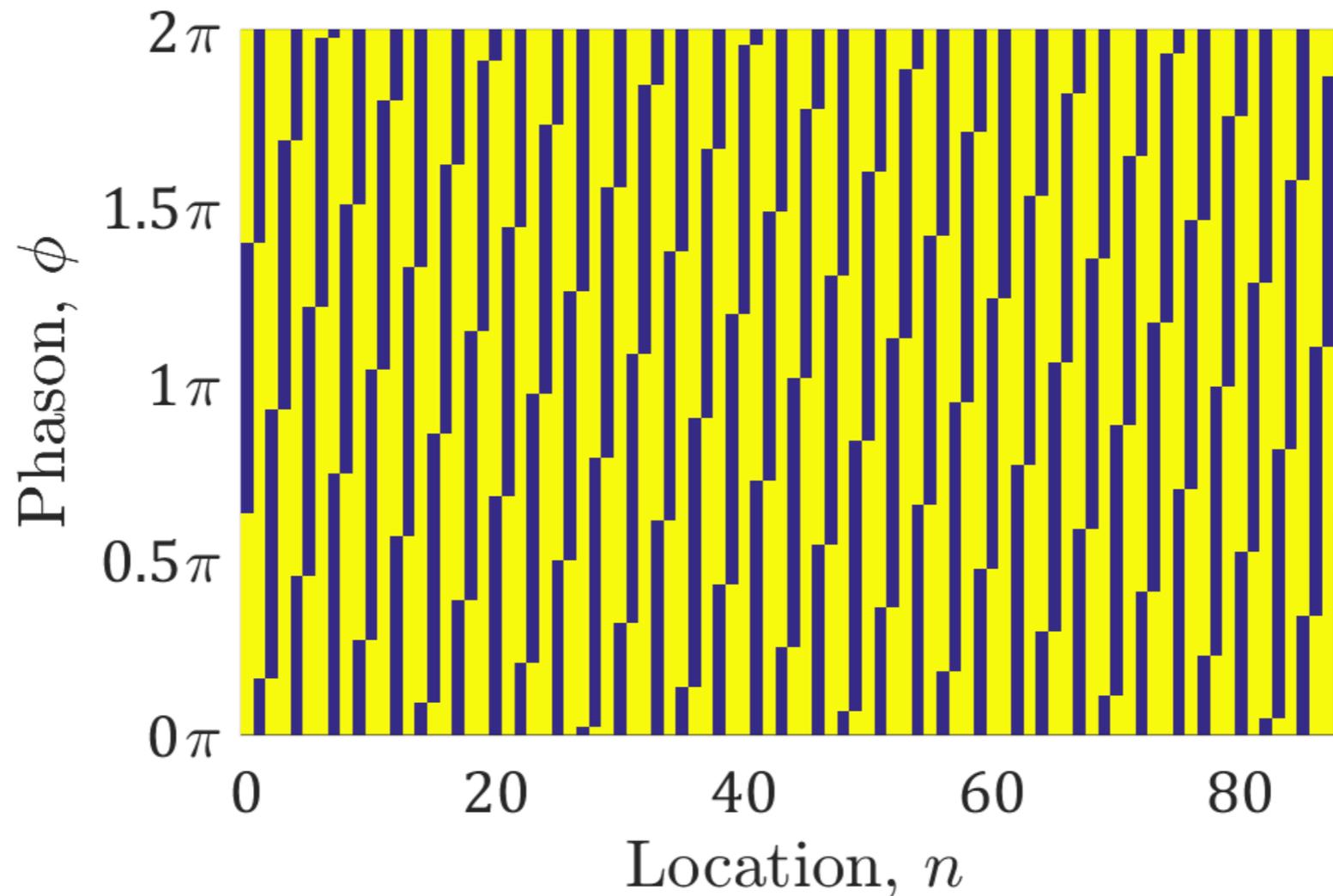
$$F_N(\phi) = [\chi_1 \chi_2 \cdots \chi_n \cdots \chi_N] \iff \text{A B A A B A B A A B A A B} \cdots$$

A torus

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with $n = 0 \dots F_N - 1$ and $[0, 2\pi] \ni \phi$



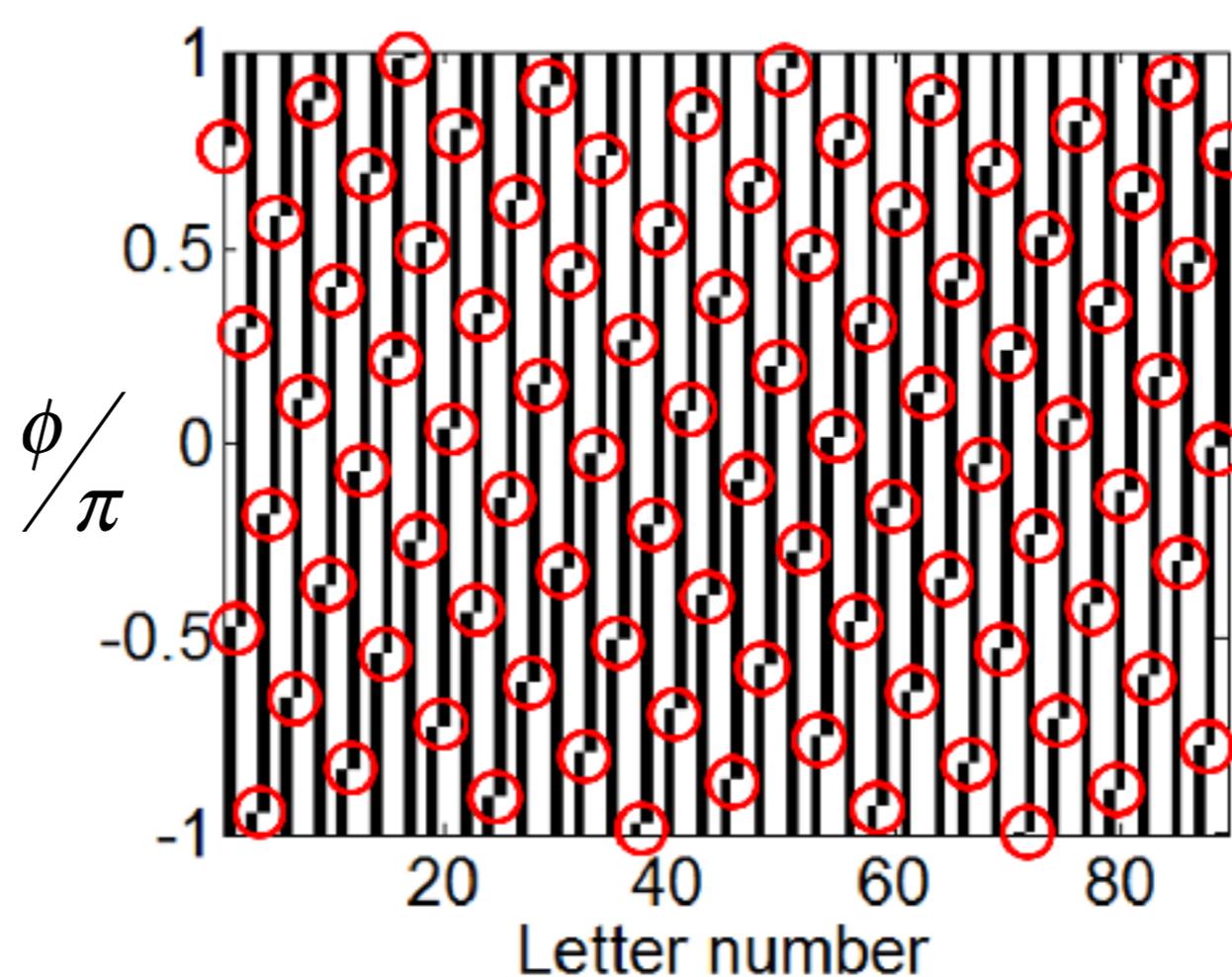
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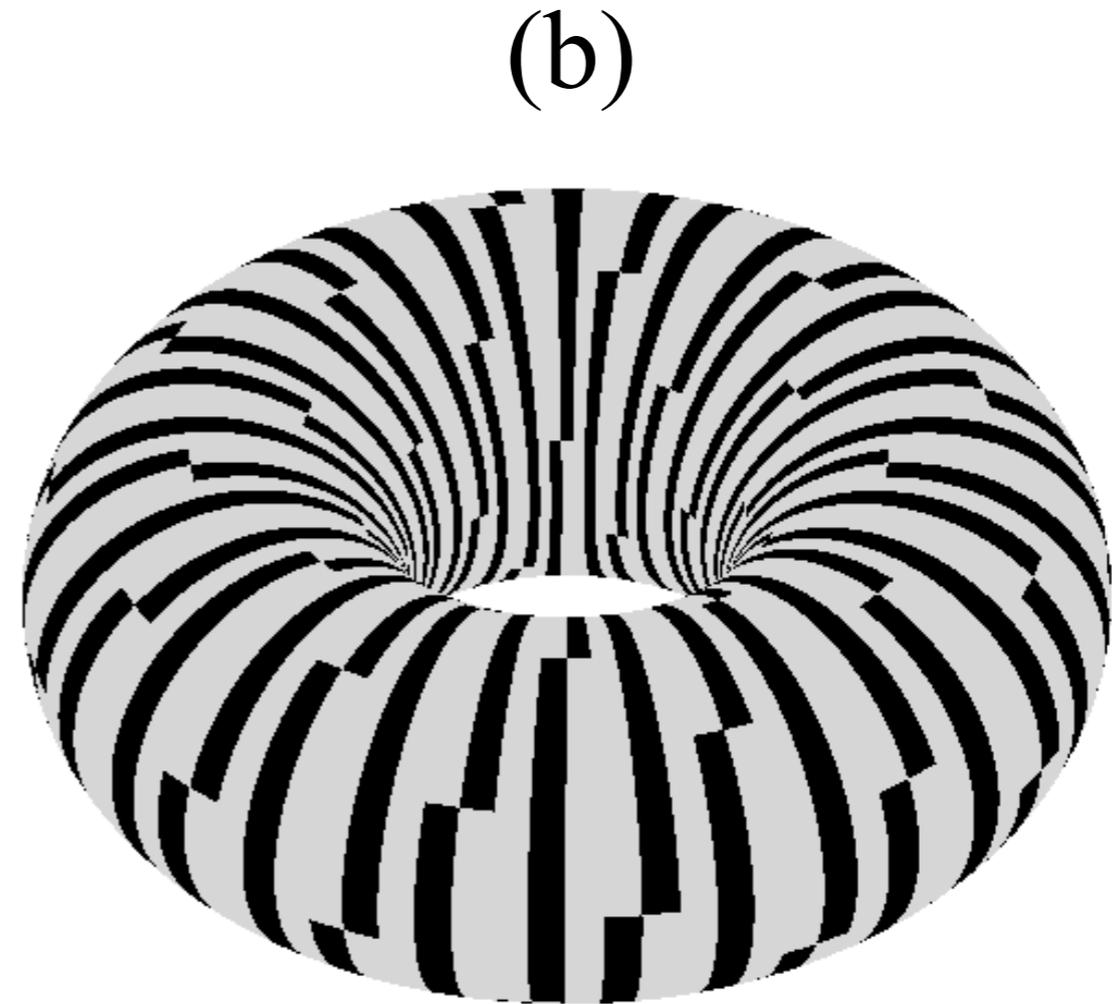
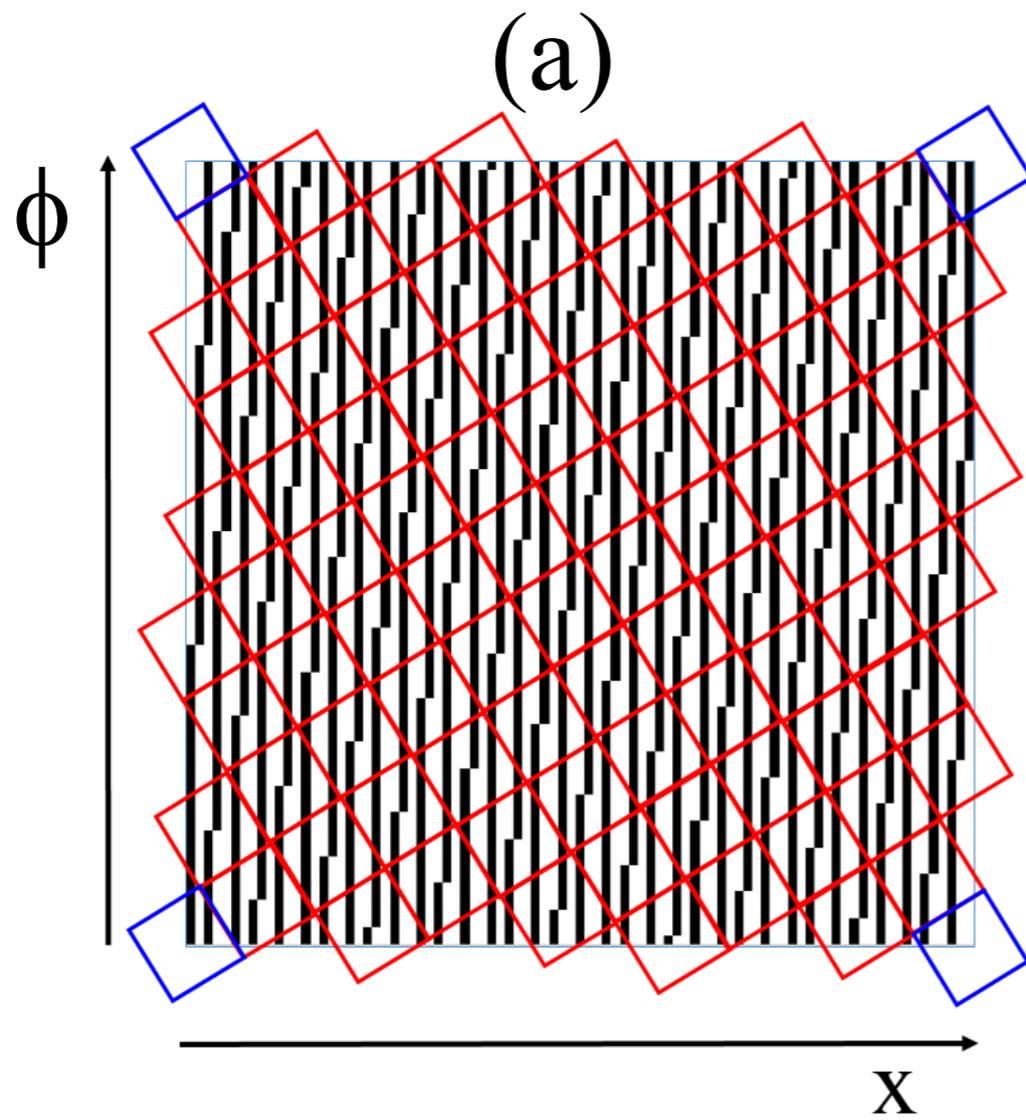
with $n = 0 \dots F_N - 1$ and $[0, 2\pi] \ni \phi$



Scanning ϕ generates **local** structural changes.



A torus



A gauge degree of freedom

- Take a characteristic function

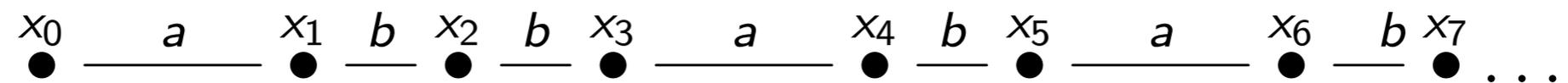
$$\chi(n, \phi) = \text{sign} \left[\cos \left(2\pi n \lambda_1^{-1} + \phi \right) - \cos \left(\pi \lambda_1^{-1} \right) \right]$$

Are there spectral consequences
of these structural properties ?

Almost No...

Atomic distributions - Structure factor

- Distributions of identical atoms in $1D$
- Use language of **tilings**: two tiles (letters) a and b
- Distribution of letters underlies distribution of atoms



- Define atomic density

$$\rho(x) = \sum_k \delta(x - x_k)$$

The diffraction pattern of the infinite chain $\rho(x) = \sum_k \delta(x - x_k)$ is given by

$$g(\xi) = \int_{-\infty}^{+\infty} dx \rho(x) e^{-i\xi x} = \sum_k e^{-i\xi x_k}$$

with structure factor

$$S(\xi) = |g(\xi)|^2$$

Definition

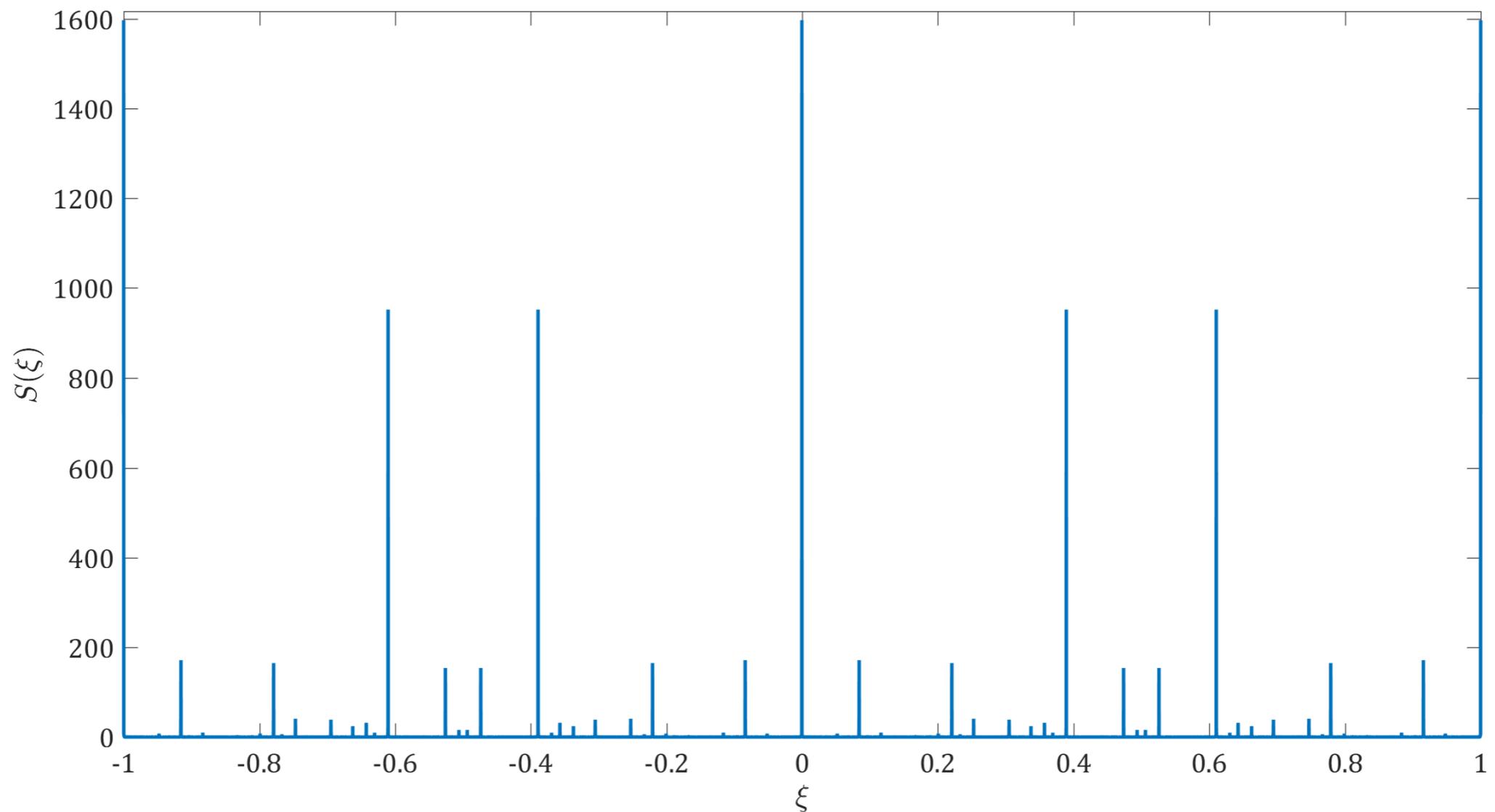
Diffraction spectrum has a **Bragg peak** (atomic part) at ξ_B iff

$$\xi_B x_c \xrightarrow{c \rightarrow \infty} 0 \pmod{2\pi}$$

for $\{x_c\}_{c=1}^{\infty} \subset \{x_k\}_{k=1}^{\infty}$, so that $g(\xi_B) \rightarrow \delta(\xi - \xi_B)$

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The extra phase - Winding numbers

- Take a characteristic function

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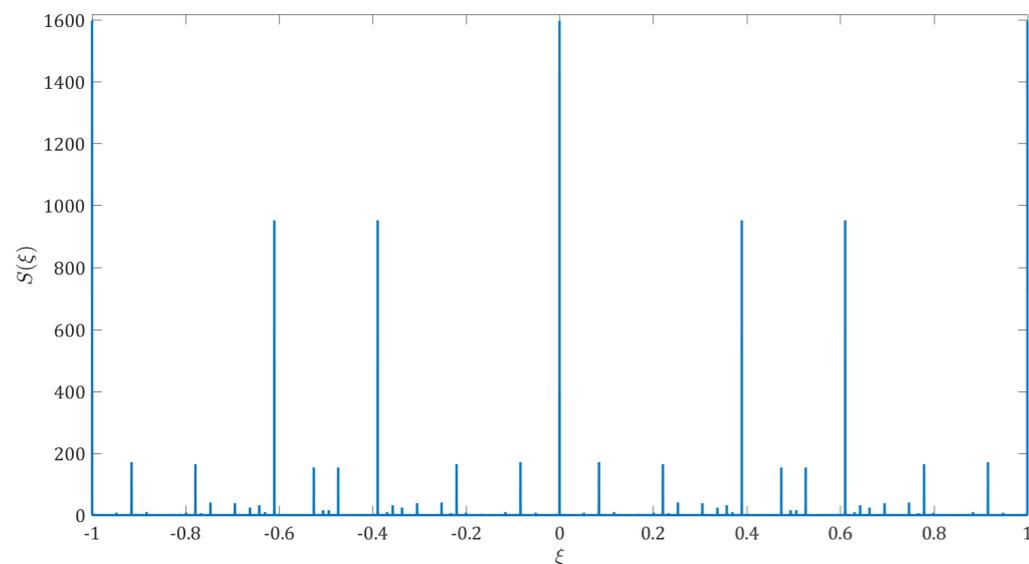
with $n = 0 \dots F_N - 1$ and $[0, 2\pi] \ni \phi \rightarrow \phi_\ell = \frac{2\pi}{F_N} \ell$

- Discrete Fourier transform w.r.t. n

$$g(\xi, \phi) = \sum_{n=0}^{F_N-1} \omega^{-\xi n} \chi(n, \phi), \quad \omega = e^{\frac{2\pi i}{F_N}}$$

- Structure factor S and phase θ

$$S(\xi, \phi) = |g(\xi, \phi)|^2$$



Structure factor is ϕ - independent

- Take a characteristic function

$$\chi(n, \phi) = \text{sign} \left[\cos \left(2\pi n \lambda_1^{-1} + \phi \right) - \cos \left(\pi \lambda_1^{-1} \right) \right]$$

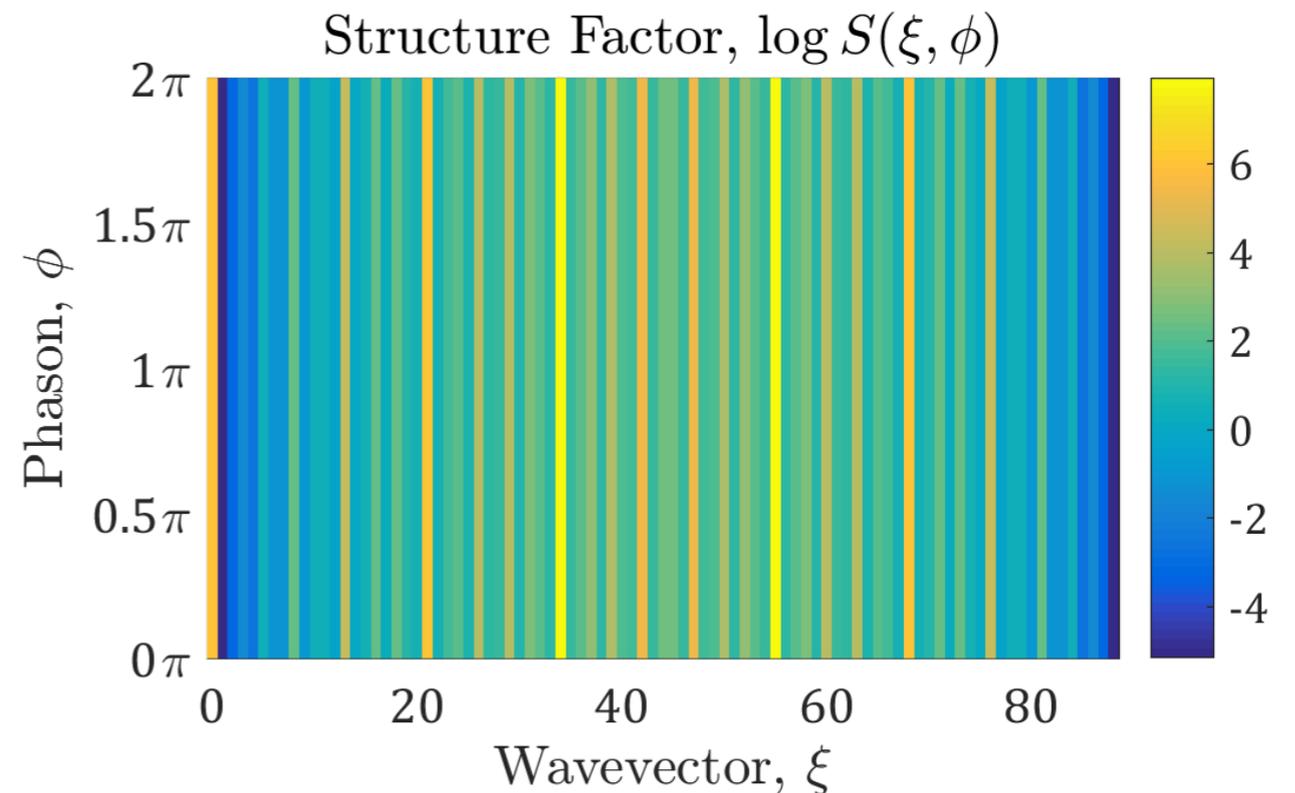
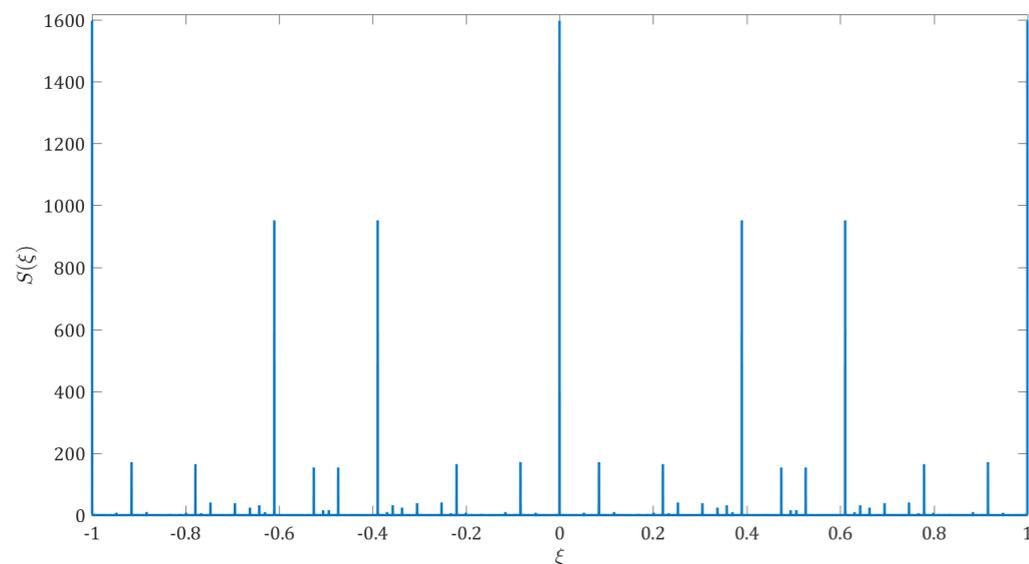
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- Structure factor S and phase θ

$$S(\xi, \phi) = |g(\xi, \phi)|^2, \quad \theta(\xi, \phi) = \arg g(\xi, \phi)$$

usually disregarded

The extra phase - Winding numbers

- Take a characteristic function

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- Structure factor S and phase θ

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- Winding number at ξ_0

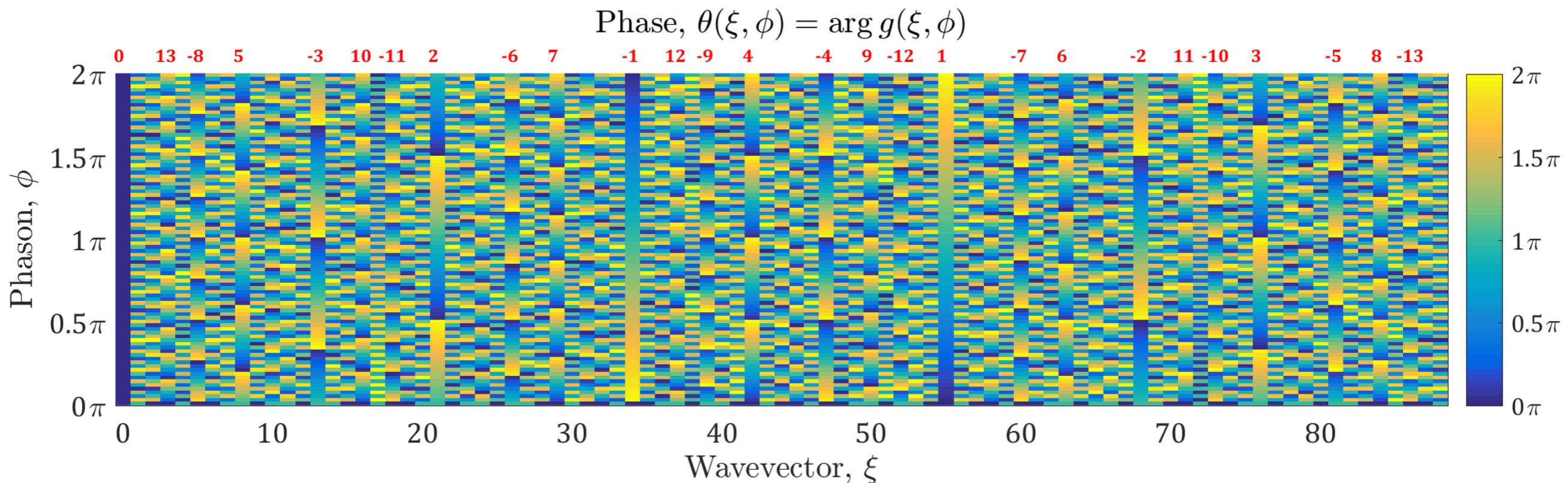
$$W_{\xi_0} = \frac{1}{2\pi} \int_0^{2\pi} \frac{\partial \theta(\xi = \xi_0, \phi)}{\partial \phi} d\phi$$

The extra phase - Winding numbers

$$g(\xi, \phi) = \sum_{n=0}^{F_N-1} \omega^{-\xi n} \chi(n, \phi), \quad \omega = e^{\frac{2\pi i}{F_N}}$$

Structure factor S and phase θ

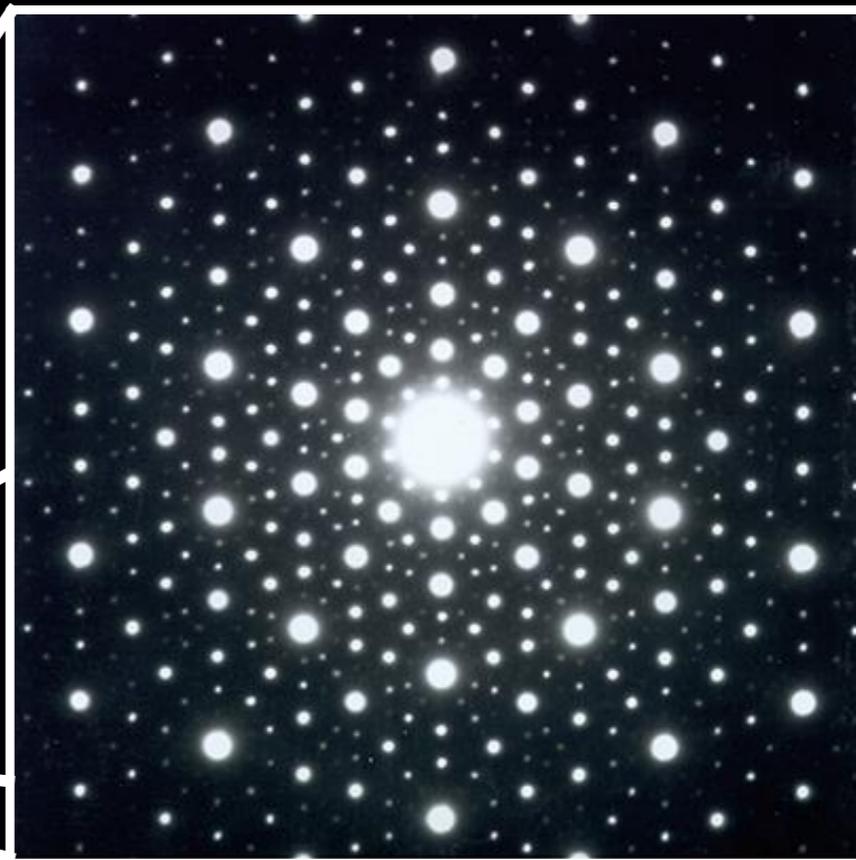
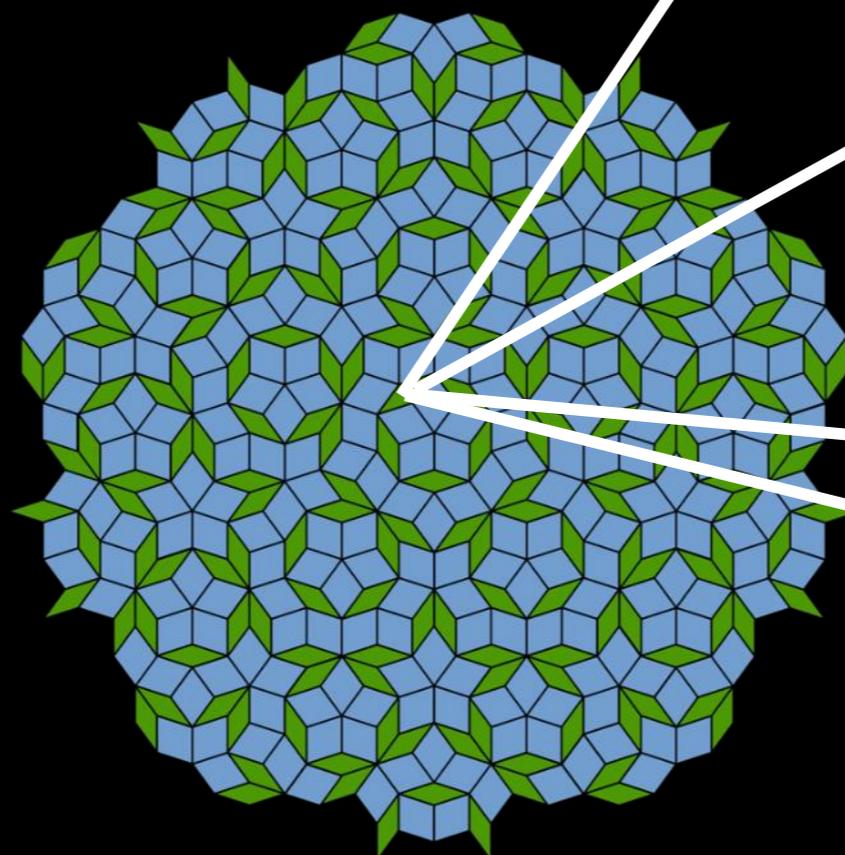
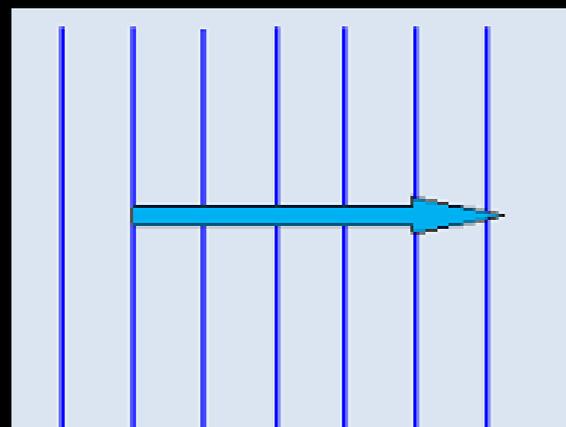
$$S(\xi, \phi) = |g(\xi, \phi)|^2, \quad \theta(\xi, \phi) = \arg g(\xi, \phi)$$



These are topological numbers!

Measuring the structural winding numbers

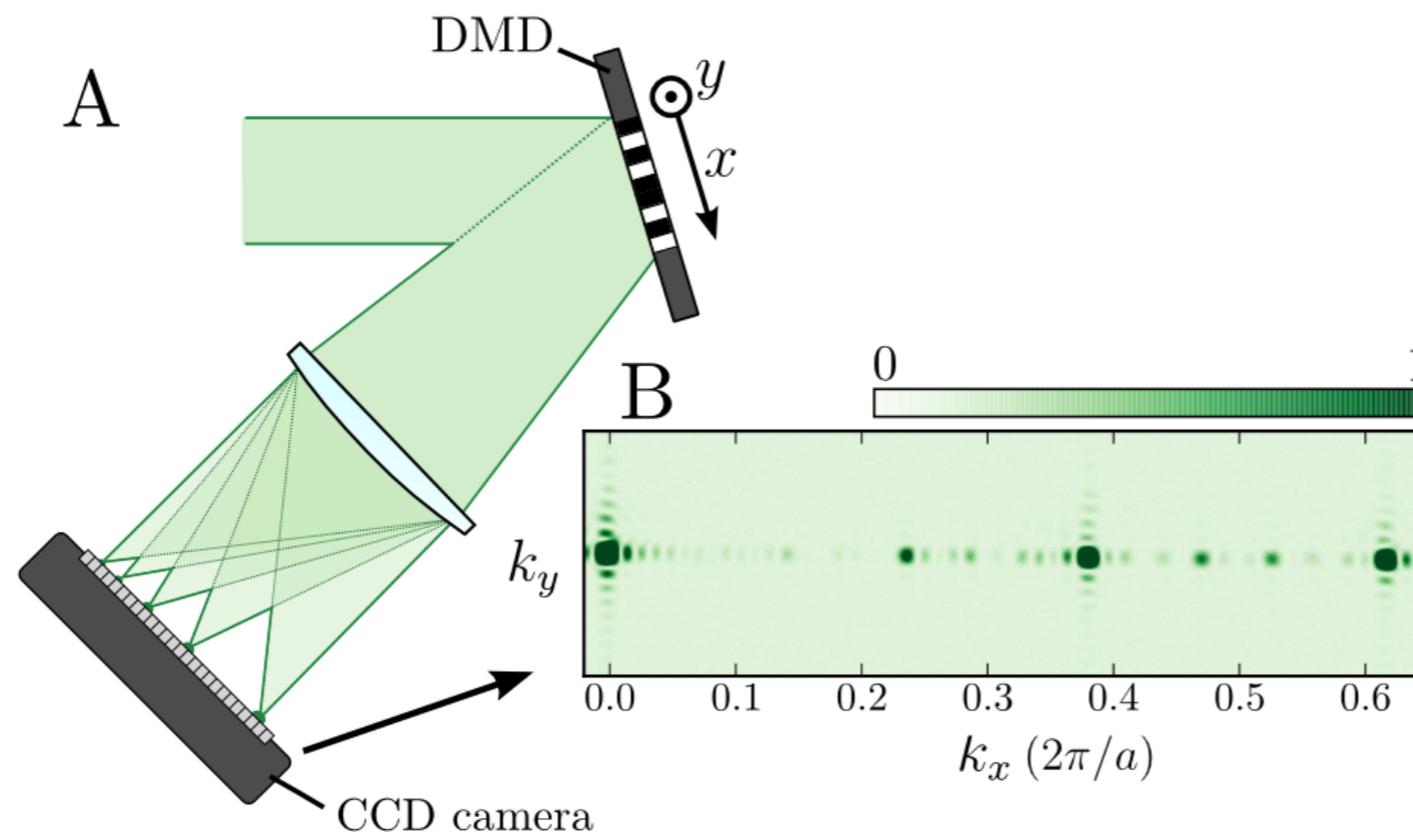
A. Dareau, E. Levy, F. Gerbier and J. Beugnon and E.A 2017



A diffraction measurement of winding numbers

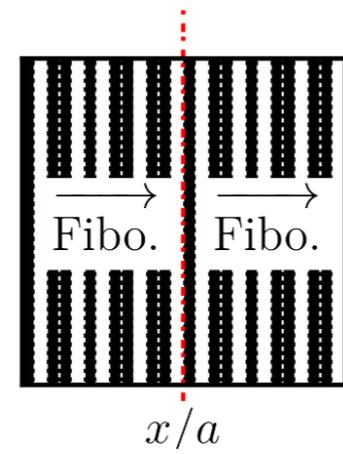


Fibonacci finite string



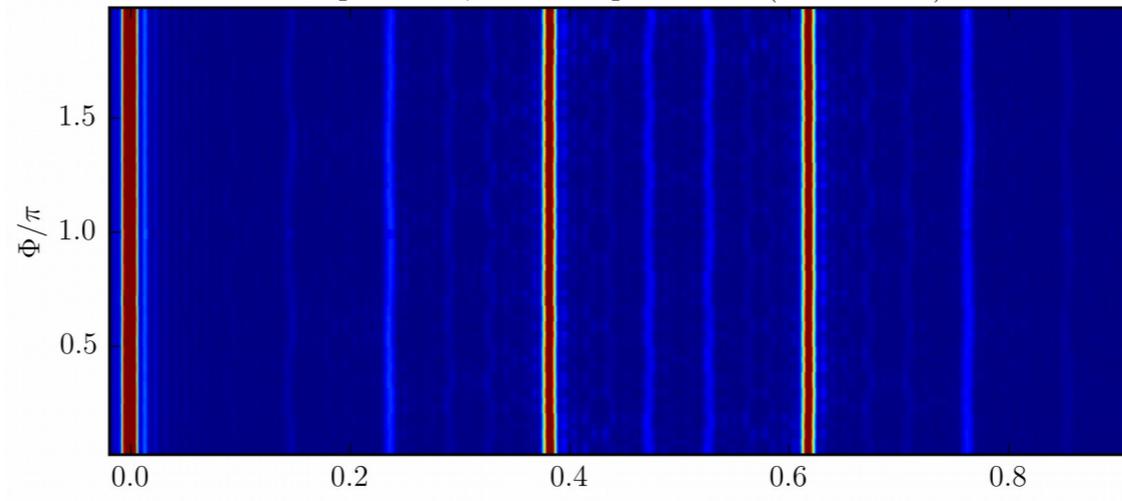
A. Dureau, E. Levy, F. Gerbier and J. Beugnon and E.A 2017

DMD Pattern



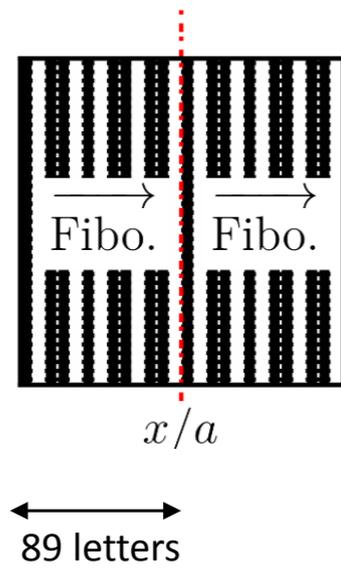
89 letters

Experiment, no artif. palindrom (linear scale)

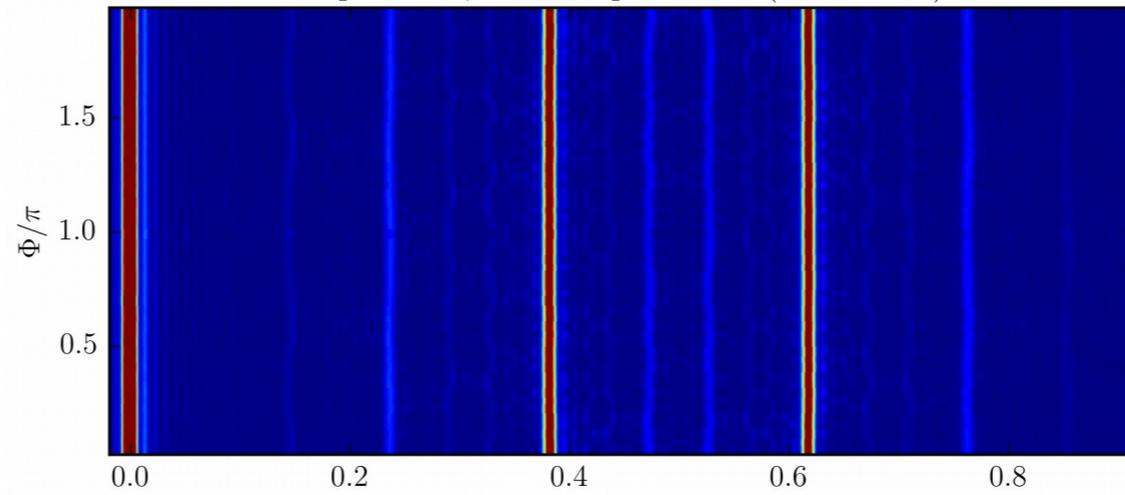


No effect of ϕ

DMD Pattern

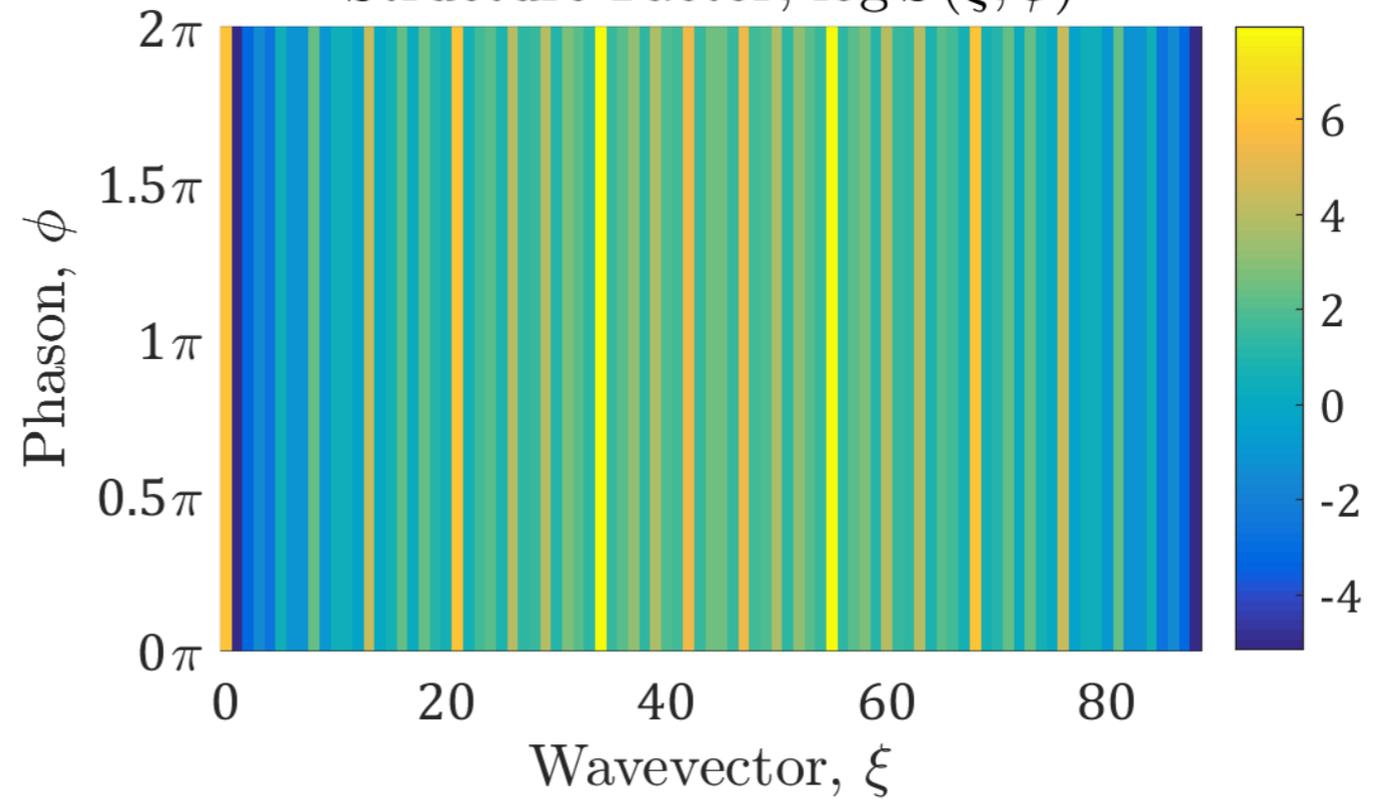


Experiment, no artif. palindrom (linear scale)

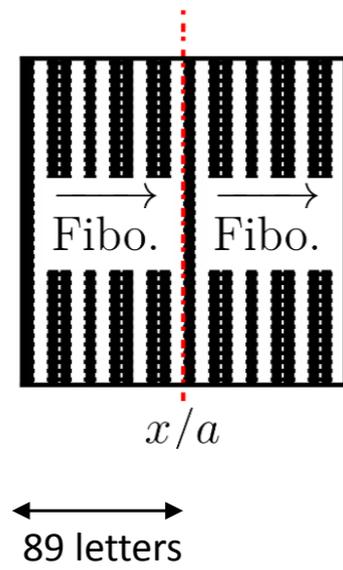


No effect of ϕ

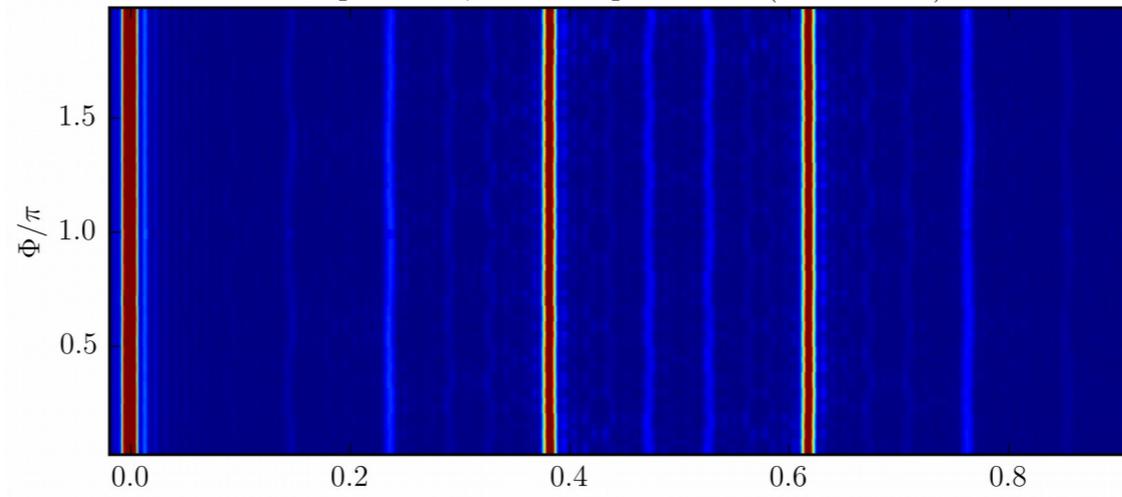
Structure Factor, $\log S(\xi, \phi)$



DMD Pattern



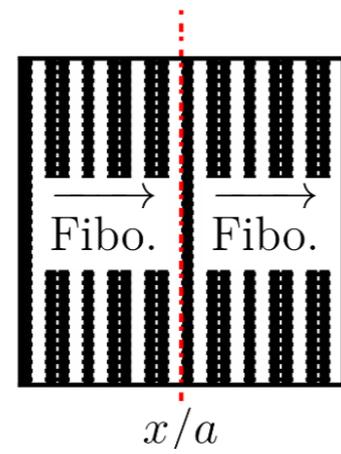
Experiment, no artif. palindrom (linear scale)



No effect of ϕ

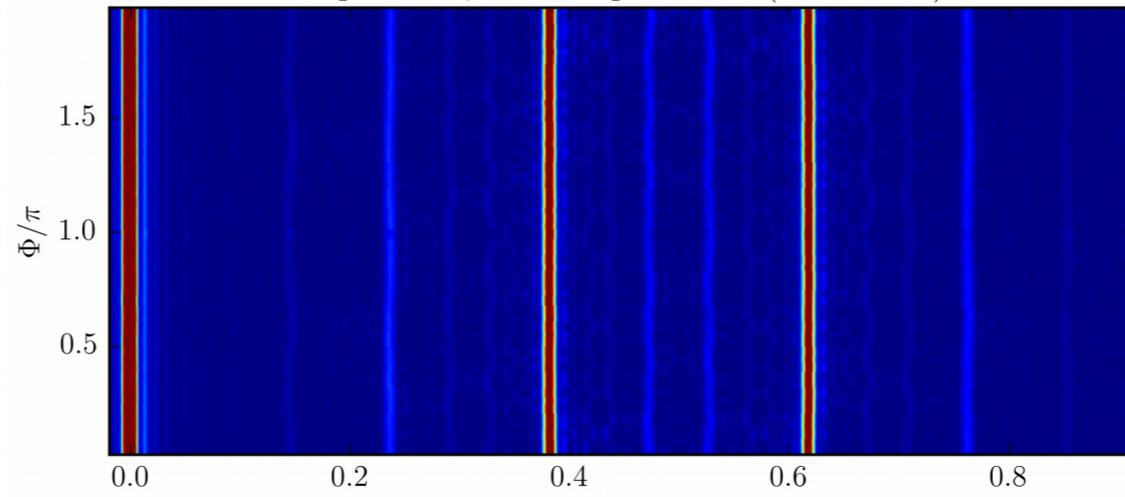
Creating edge states

DMD Pattern

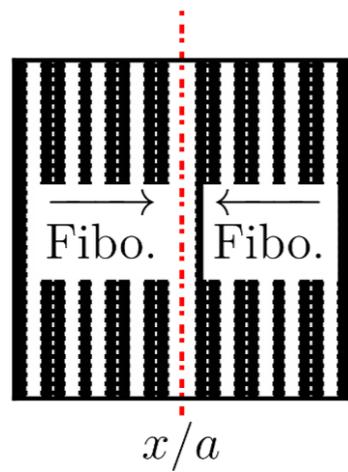


89 letters

Experiment, no artif. palindrom (linear scale)

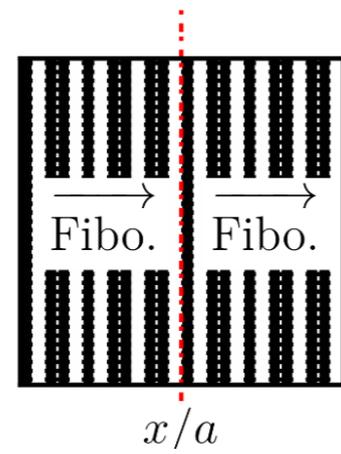


DMD Pattern



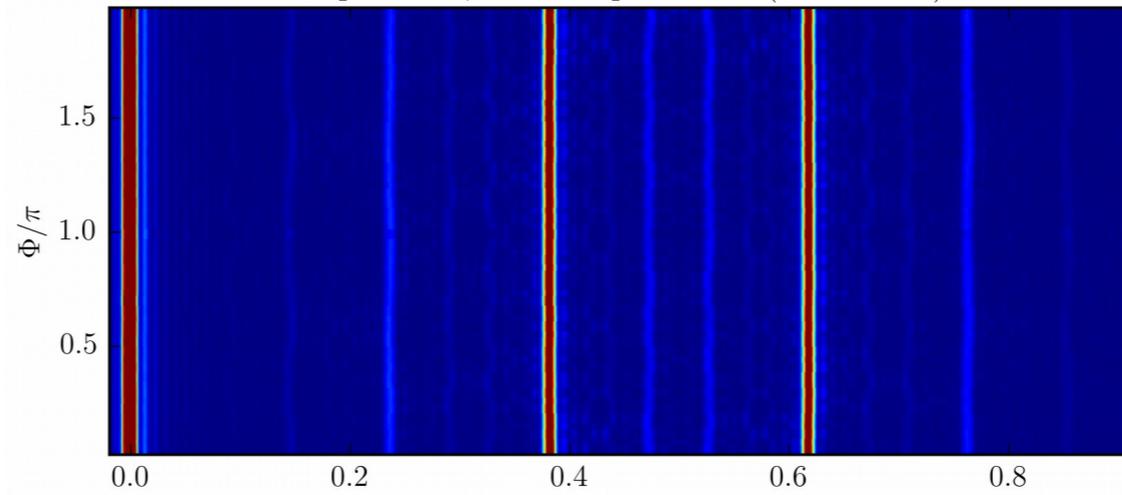
89 letters

DMD Pattern



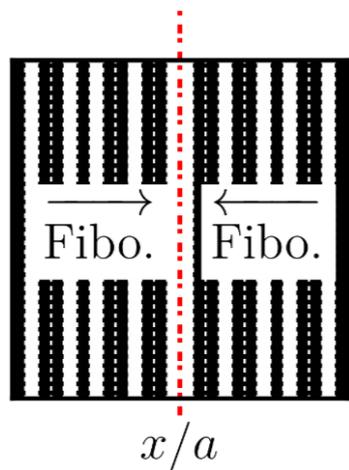
89 letters

Experiment, no artif. palindrom (linear scale)



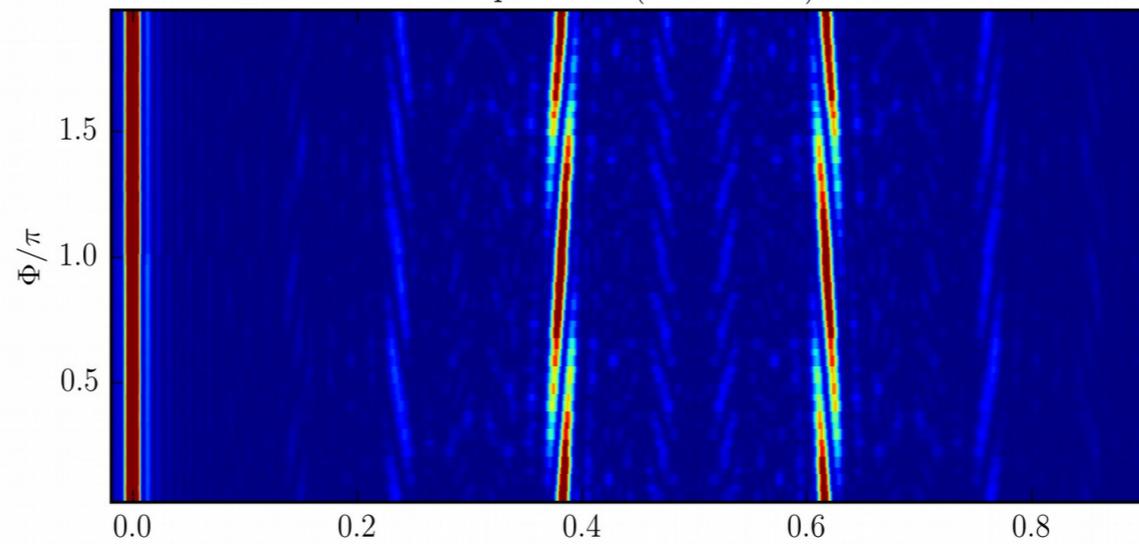
No effect of ϕ

DMD Pattern



89 letters

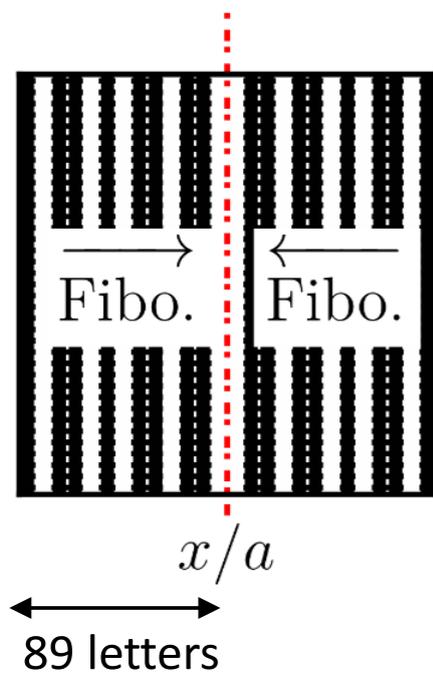
Experiment (linear scale)



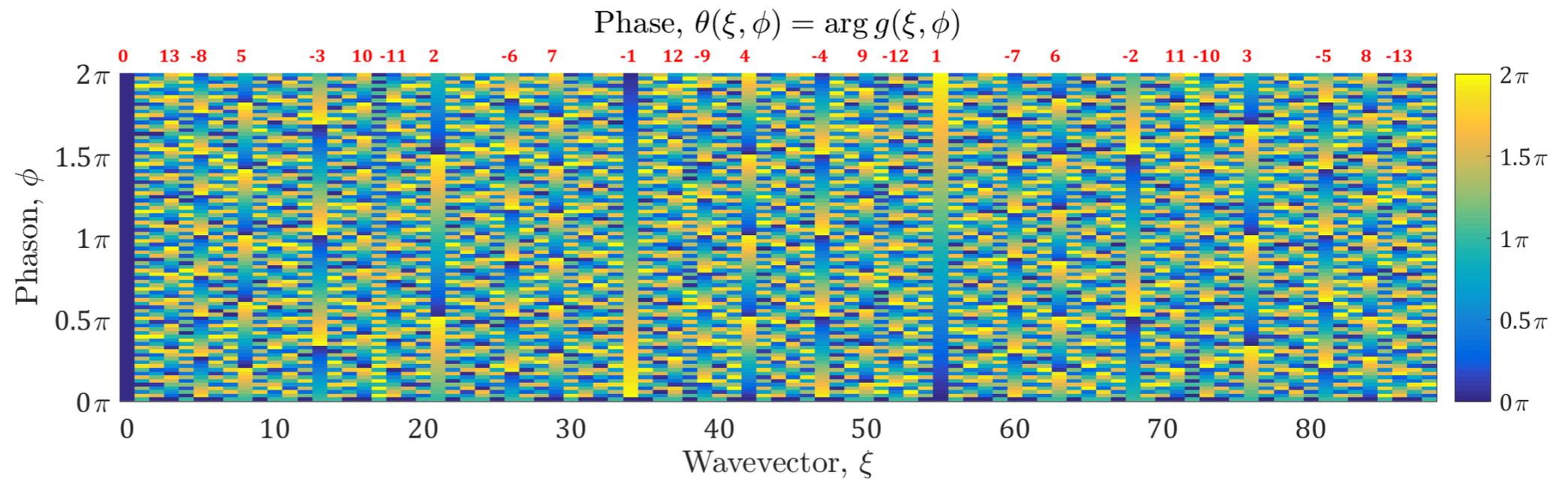
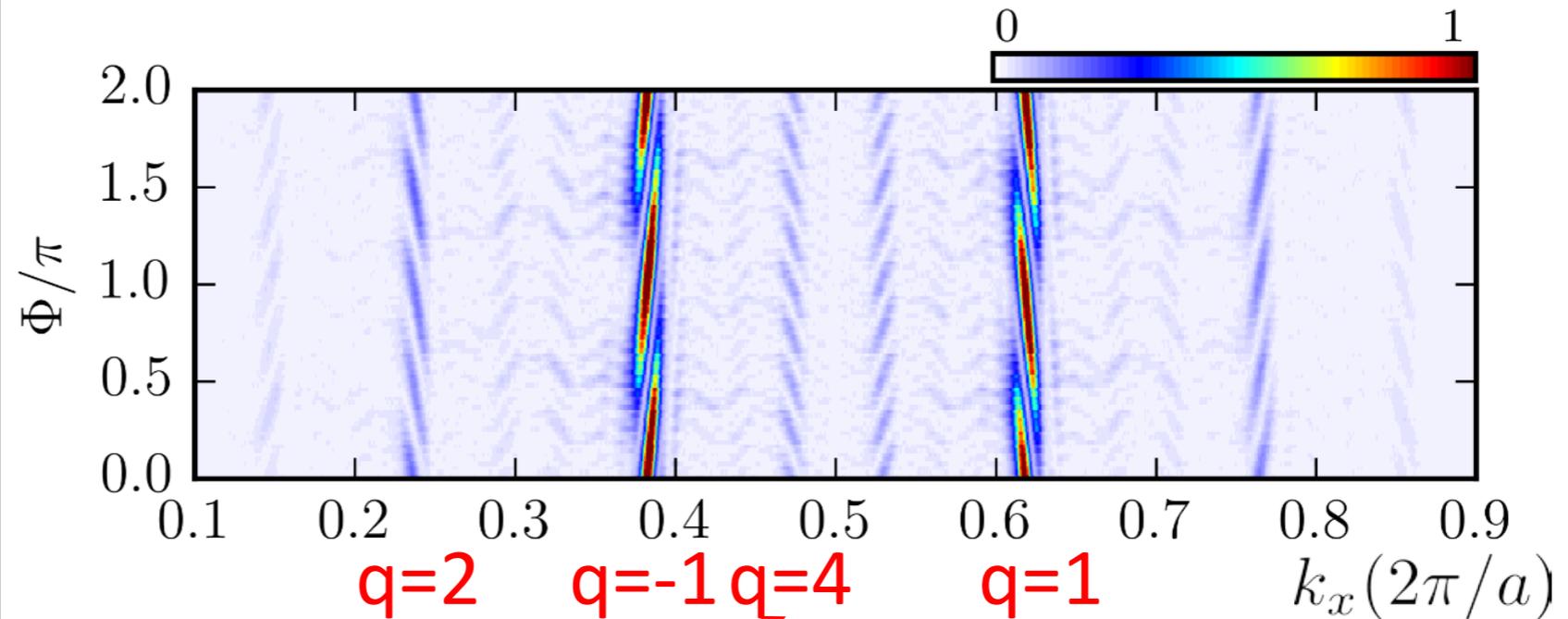
There is an effect of the phason ϕ

A diffraction measurement of winding numbers

DMD Pattern



Diffraction pattern



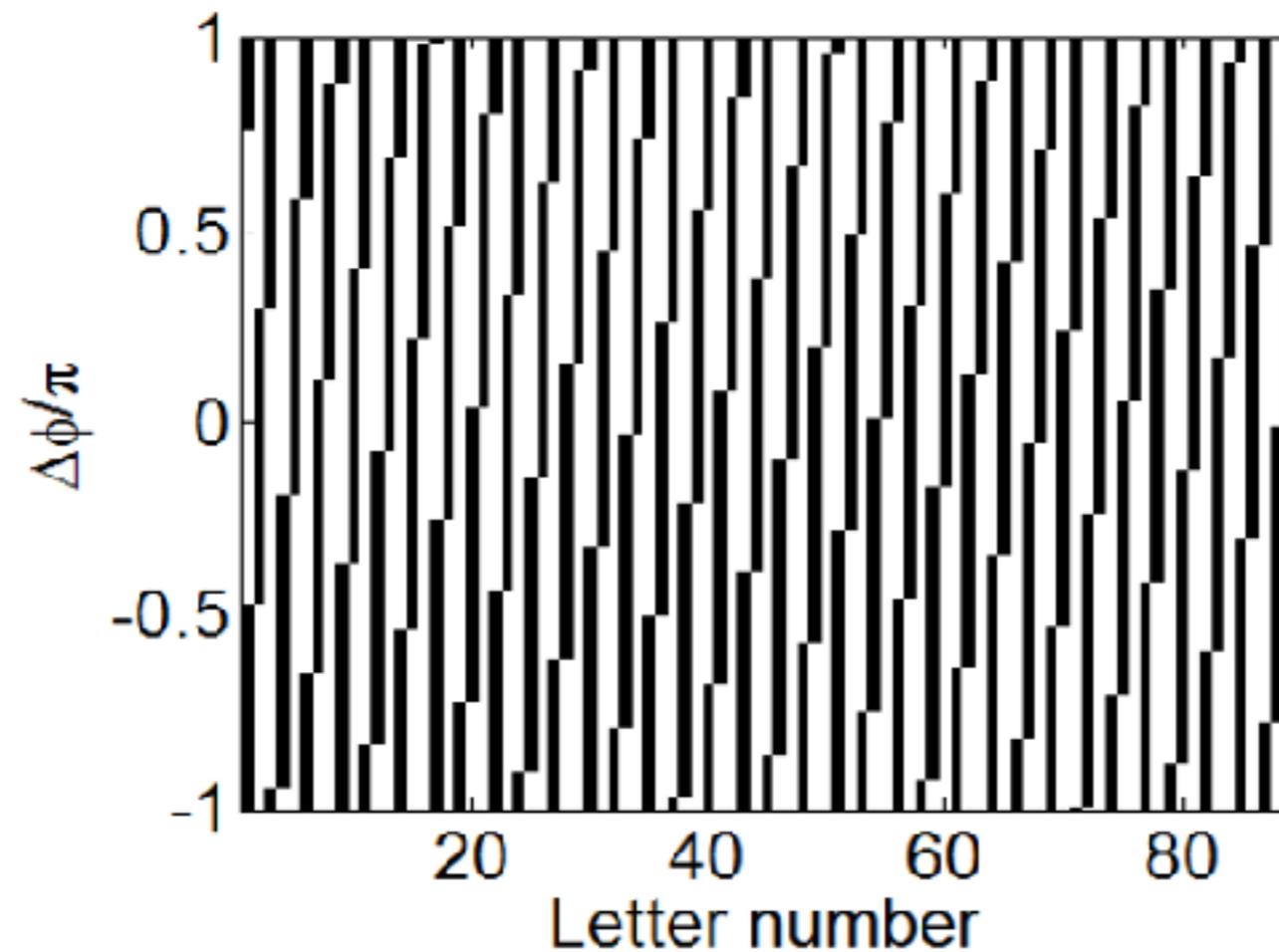
$$W_{\xi_0} = \frac{1}{2\pi} \int_0^{2\pi} \frac{\partial \theta(\xi = \xi_0, \phi)}{\partial \phi} d\phi$$

2D diffraction experiment

Instead of



consider all realisations

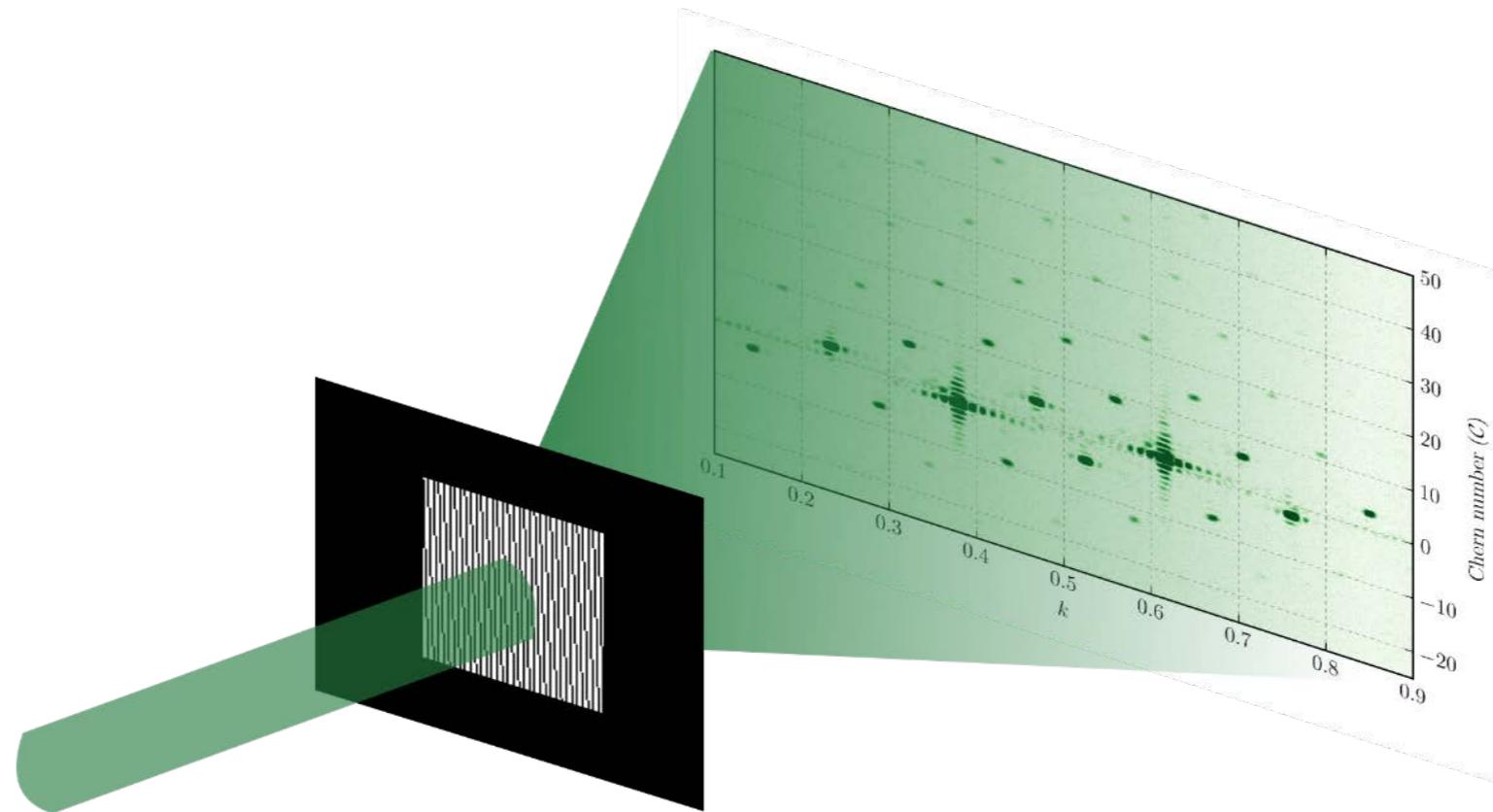


2D diffraction experiment

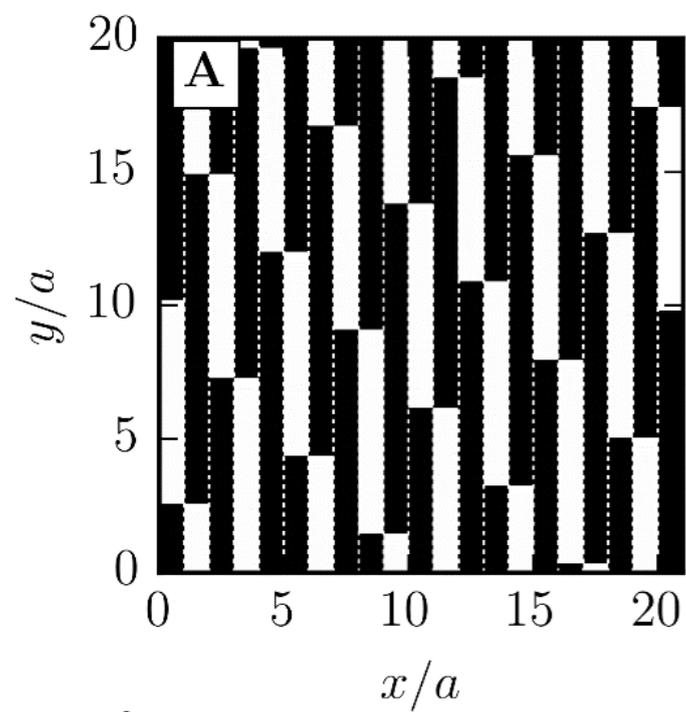
Instead of



consider all realisations

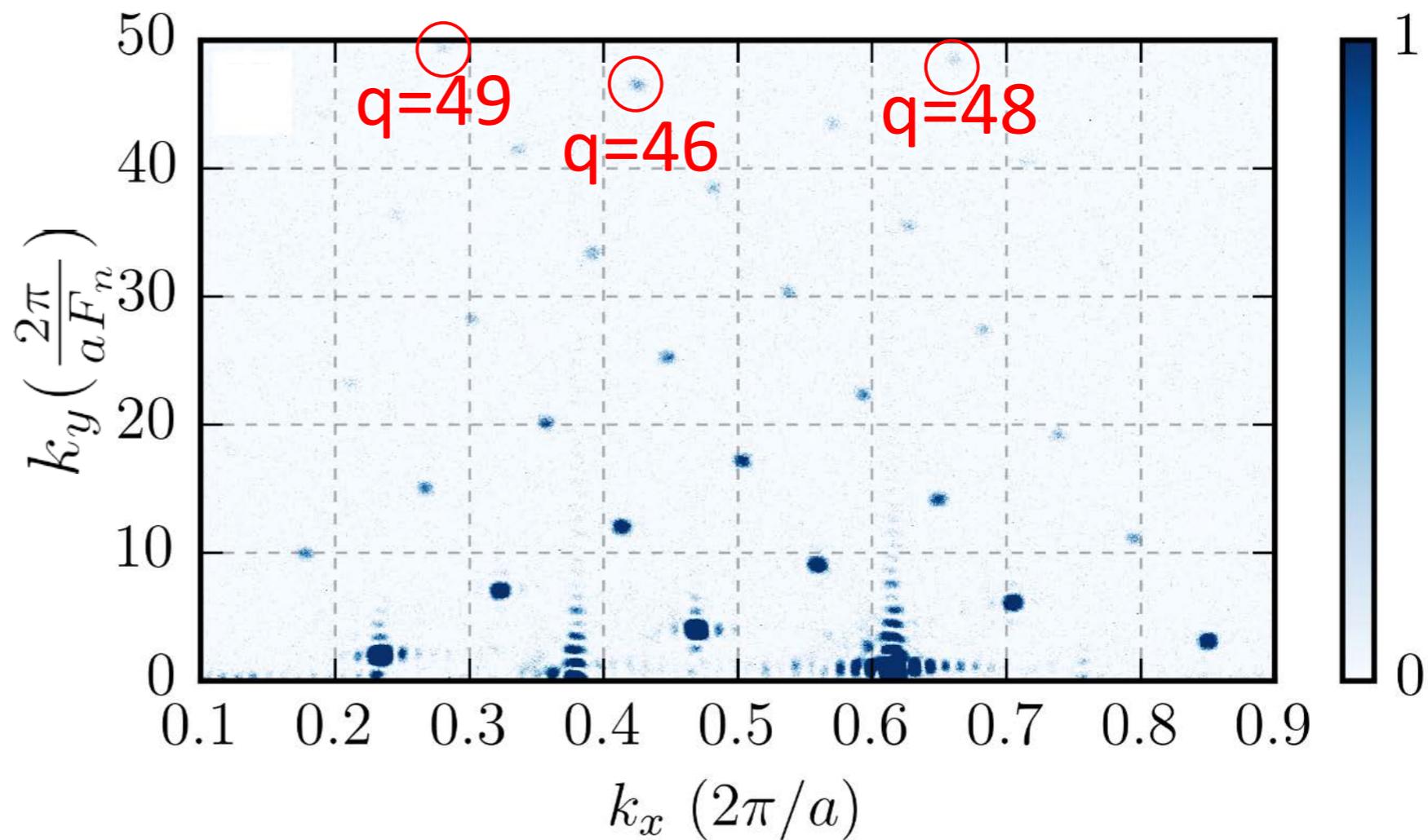


DMD Pattern



y axis is associated with Φ

Diffraction pattern



Spectral Features

So far, we presented structural features
culminating in topological winding numbers

What about spectral ones?

How to Characterize Tilings – Spectrum?

Spectrum & Integrated Density of States

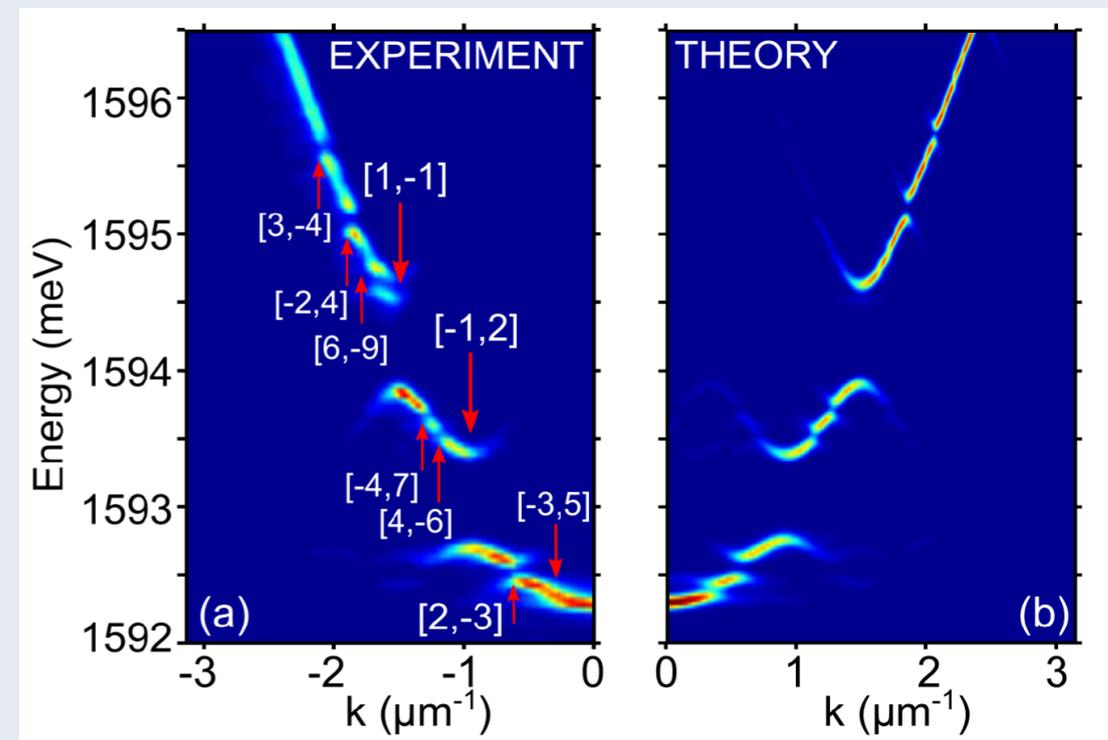
- Solve a Hamiltonian $H(E)$ (Fibonacci)

$$H\psi(x) = E\psi(x)$$



- Find the spectrum:
 - dispersion $E(k)$
 - integrated density of states

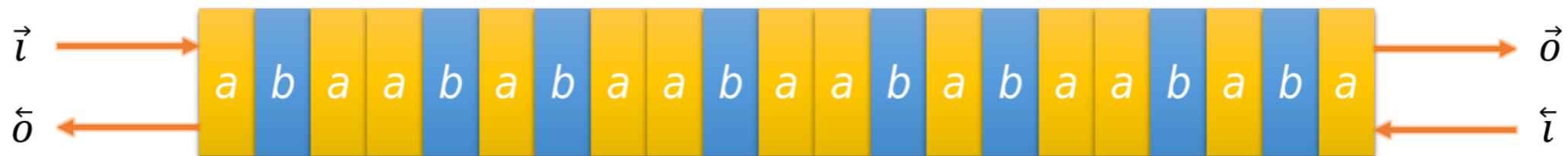
$$H(E) \rightarrow \begin{cases} \varrho(E) & \text{DOS} \\ \mathcal{N}(E) & \text{IDOS} \end{cases}$$



Scattering formalism

- Take 1D wave system of size L bounded by two semi-infinite free systems

$$\underbrace{\begin{array}{c|c} \vec{l} & \vec{o} \\ \hline \vec{o} & \vec{l} \end{array}}_{\mathcal{S}} \Rightarrow \begin{pmatrix} \vec{o} \\ \vec{o} \end{pmatrix} = \begin{pmatrix} \vec{r}(k) & t(k) \\ t(k) & \vec{r}(k) \end{pmatrix} \begin{pmatrix} \vec{l} \\ \vec{l} \end{pmatrix} \equiv \mathcal{S} \begin{pmatrix} \vec{l} \\ \vec{l} \end{pmatrix}$$



Scattering formalism

- The \mathcal{S} -matrix is diagonalized to

$$\mathcal{S} \mapsto \begin{pmatrix} e^{i\phi_1} & 0 \\ 0 & e^{i\phi_2} \end{pmatrix} \Rightarrow \det \mathcal{S} = e^{2i\delta(k)}$$

with $\delta(k) = \frac{1}{2} (\phi_1(k) + \phi_2(k))$

Scattering formalism

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with $\delta(k) = \frac{1}{2} (\phi_1(k) + \phi_2(k))$

- Find density of modes with Krein-Schwinger formula ,

$$\varrho(k) - \varrho_0(k) = \frac{1}{2\pi} \operatorname{Im} \frac{d}{dk} \ln \det \mathcal{S}(k)$$

- The **normalized IDOS** is given by

$$\mathcal{N}(\nu) - \mathcal{N}_0(\nu) = \frac{1}{2\pi} \text{Im} \log \det \mathcal{S}(\nu, \phi)$$

independent of ϕ

- The **normalized IDOS** is given by

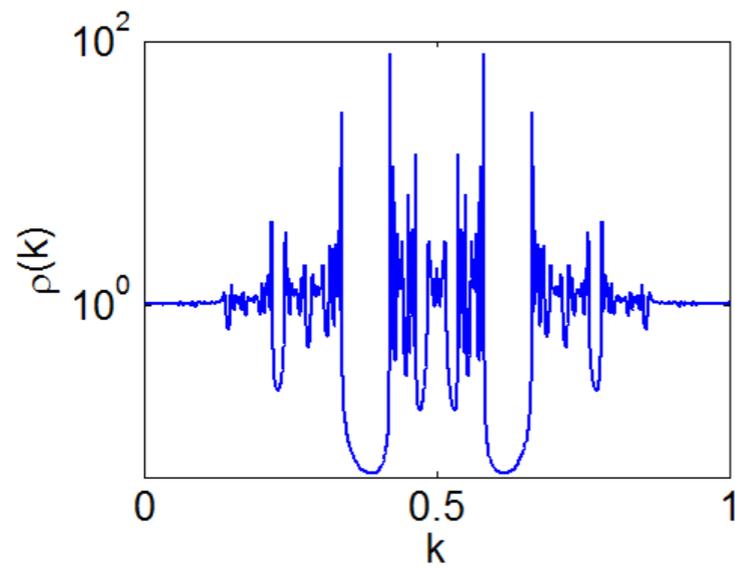
$$\mathcal{N}(\nu) - \mathcal{N}_0(\nu) = \frac{1}{2\pi} \text{Im} \log \det \mathcal{S}(\nu, \phi)$$

independent of ϕ

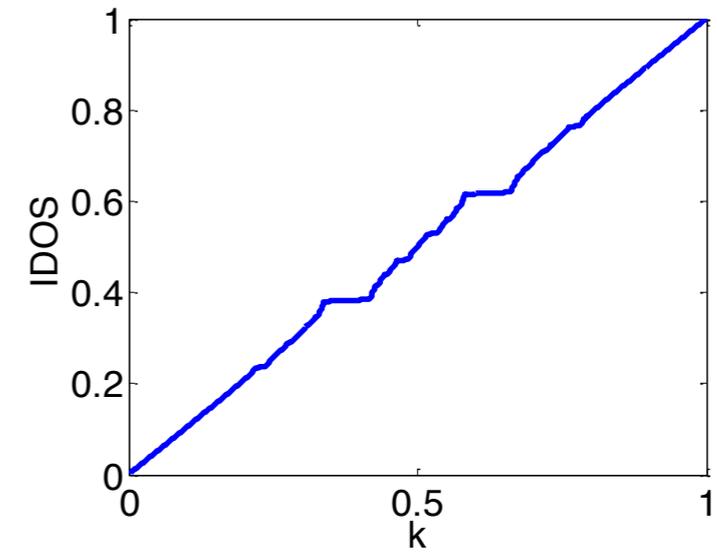
- For C&P, **gaps** appear at

$$\mathcal{N}_{\text{gap}} = p + q s \pmod{1}$$

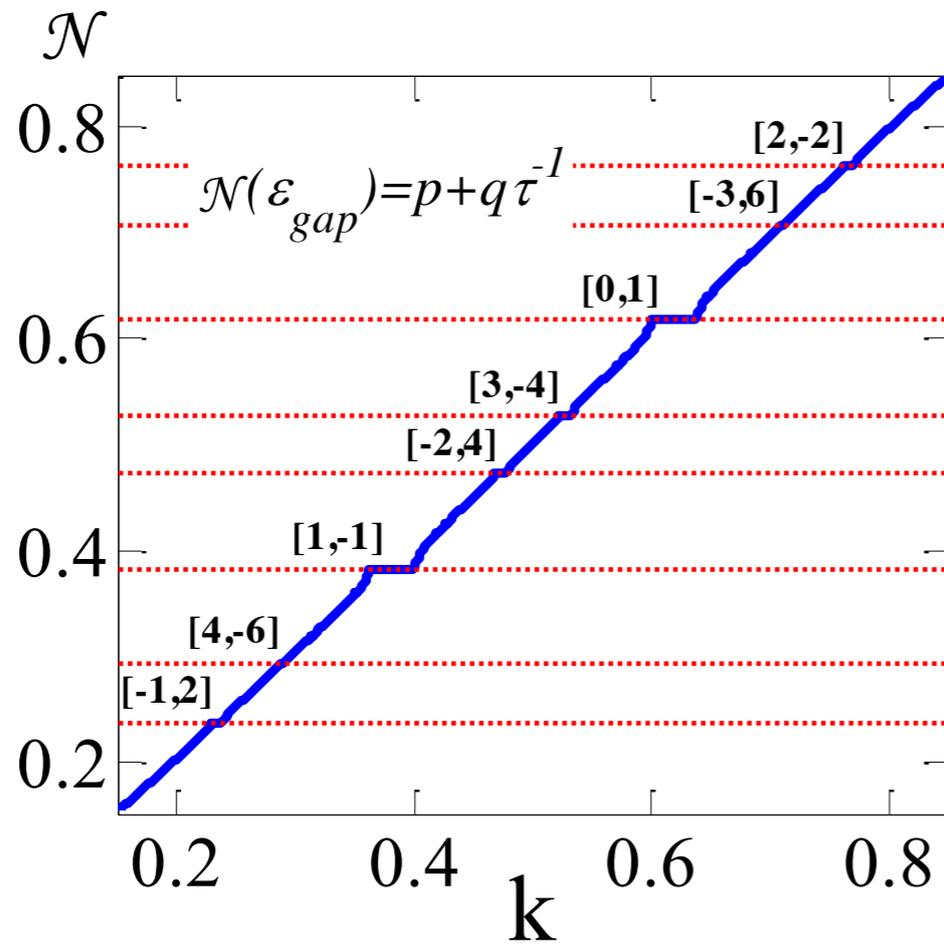
this is the GLT



Density of modes



IDOS-counting function



Gap Labeling Theorem (GLT)

- The \mathcal{S} -matrix is a 2×2 unitary matrix (in $1d$) : $\mathcal{S} \sim \begin{pmatrix} e^{i\gamma_1} & 0 \\ 0 & e^{i\gamma_2} \end{pmatrix}$
 - Uniquely identified by 2 phases
 - That can be written **universally**

- 1 A ϕ -independent spectral total phase shift

$$\delta(\nu) = \frac{1}{2} (\gamma_1 + \gamma_2) = \frac{1}{2} \text{Im} \log \det \mathcal{S}(\nu, \phi)$$

with $\mathcal{N}(\nu) - \mathcal{N}_0(\nu) = \frac{1}{\pi} \delta(\nu)$

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- 2 A ϕ -dependent spectral chiral phase by

$$\alpha(\nu_{\text{gap}}, \phi) = \gamma_1 - \gamma_2 = \text{Im} \text{Tr} [\sigma_z \log \mathcal{S}(\nu_{\text{gap}}, \phi)]$$

Where there is a ϕ -dependent phase – there is a **winding!**

Spectral Winding

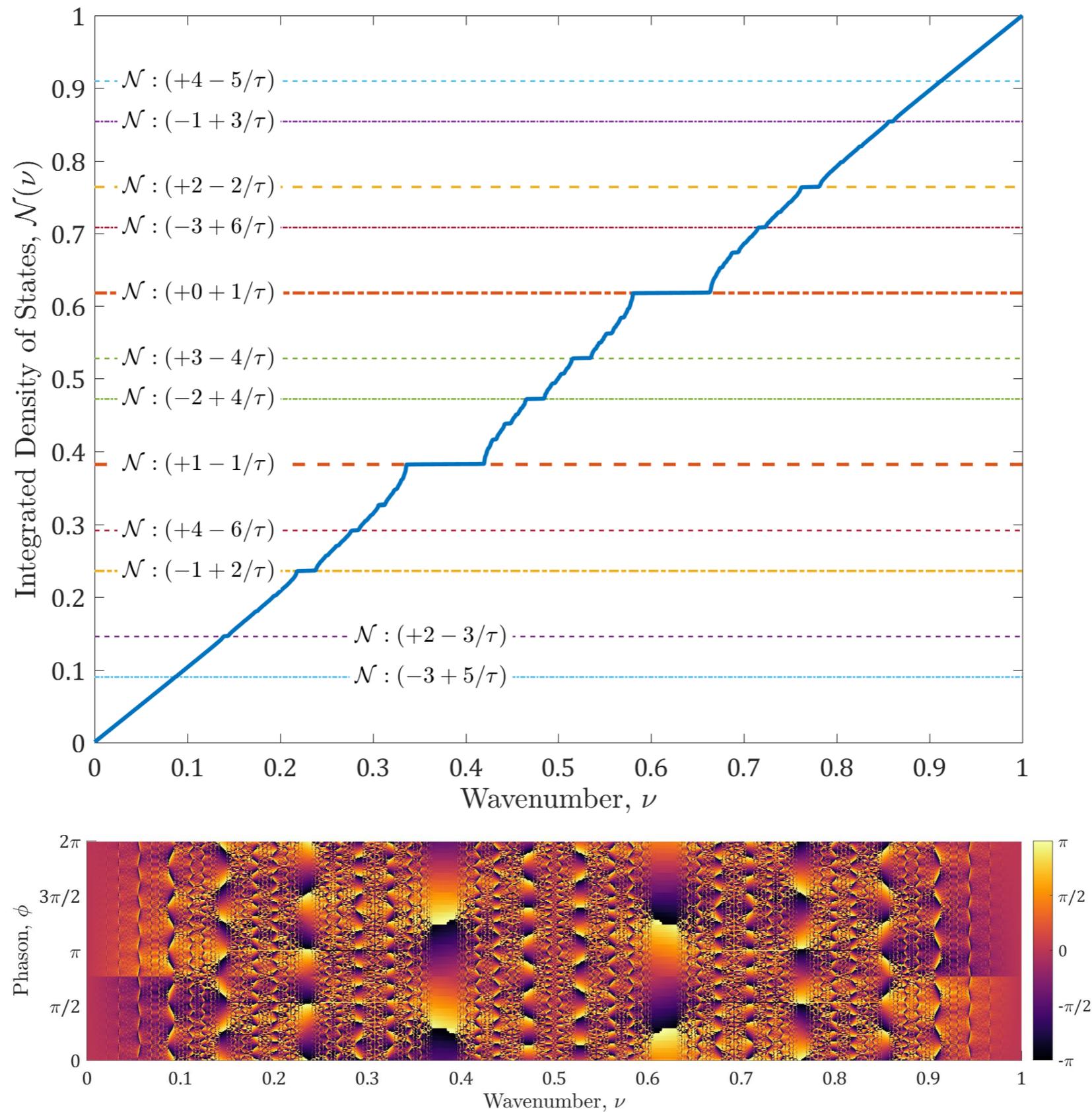
- To each gap $\mathcal{N}_{\text{gap}} = q s \pmod{1}$, count the winding

$$\mathcal{W}_{\phi}[\alpha] = \frac{1}{2\pi} \int_0^{2\pi} \frac{\partial \alpha(\nu_{\text{gap}}, \phi)}{\partial \phi} d\phi$$

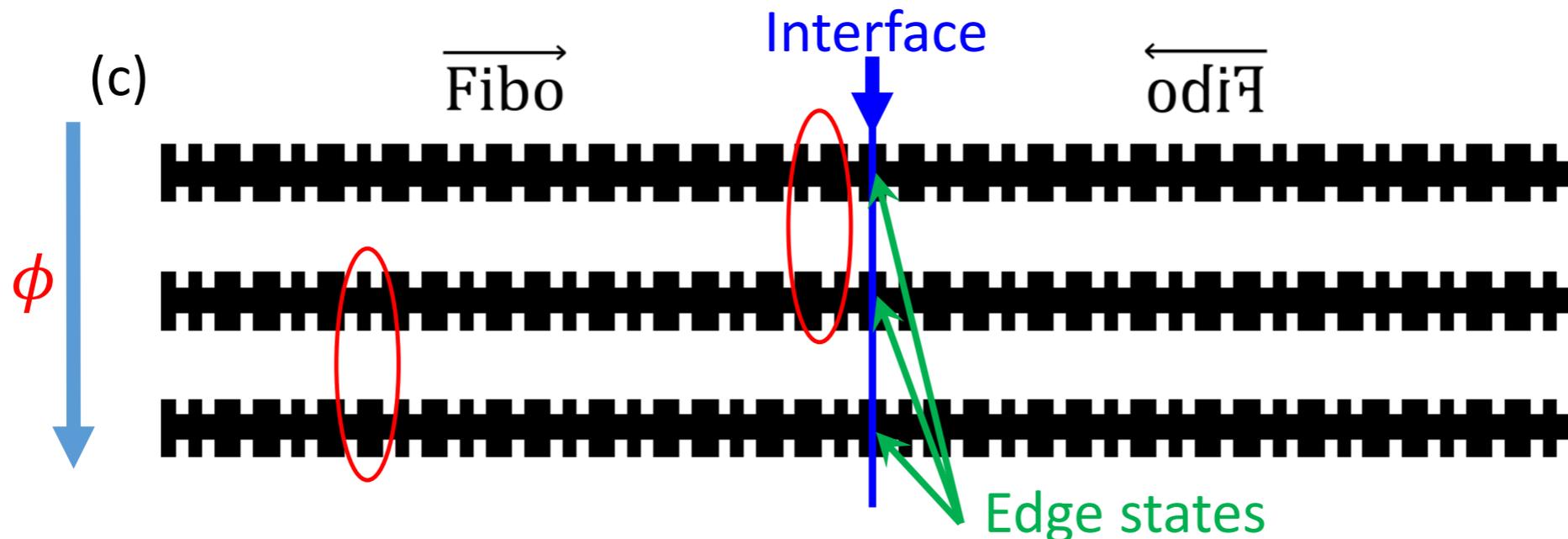
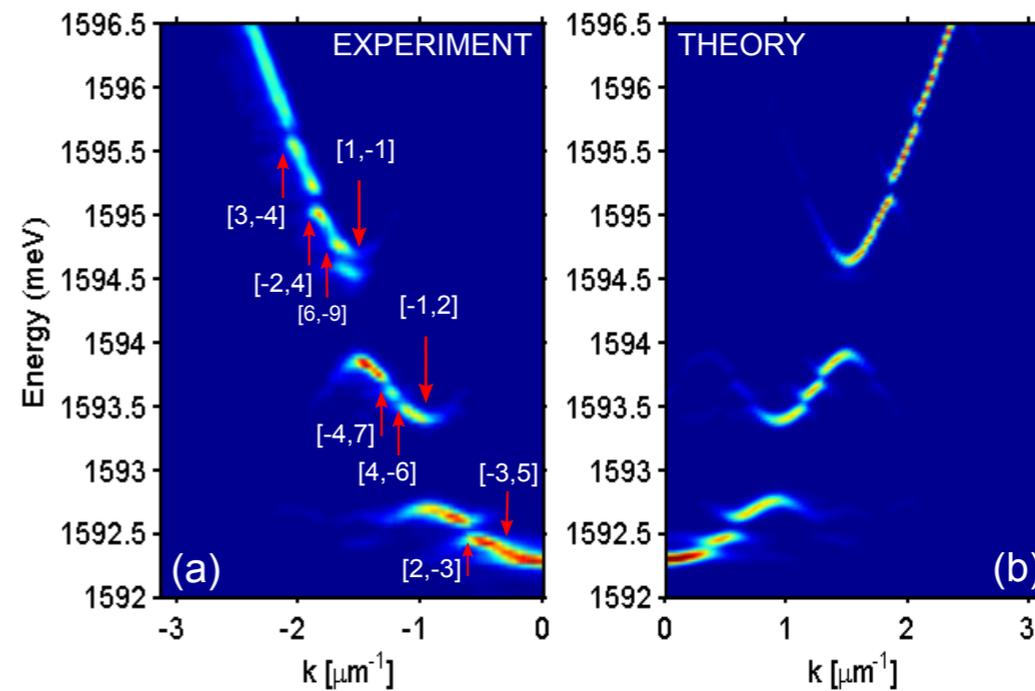
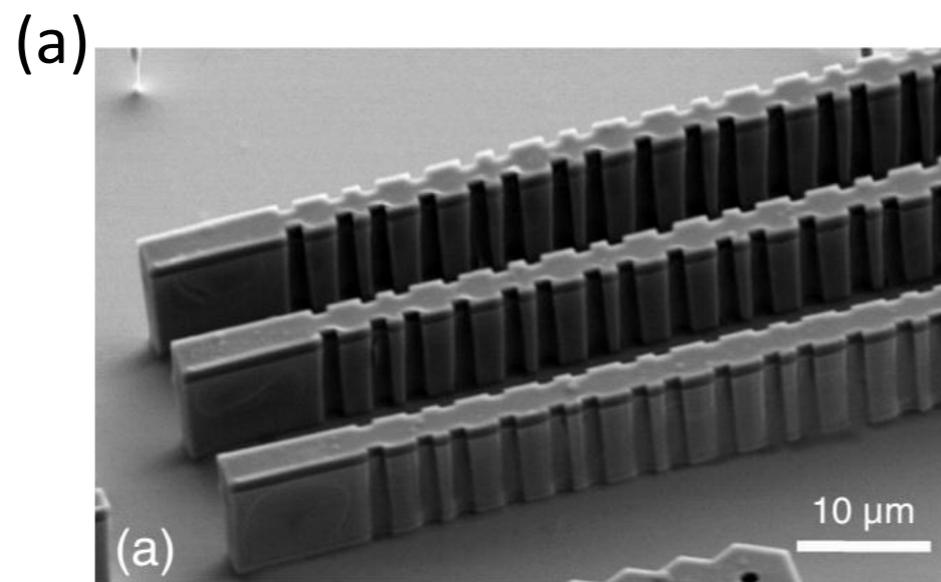
- Numerical calculation yields

$$\mathcal{W}_{\phi}[\alpha] = 2q$$

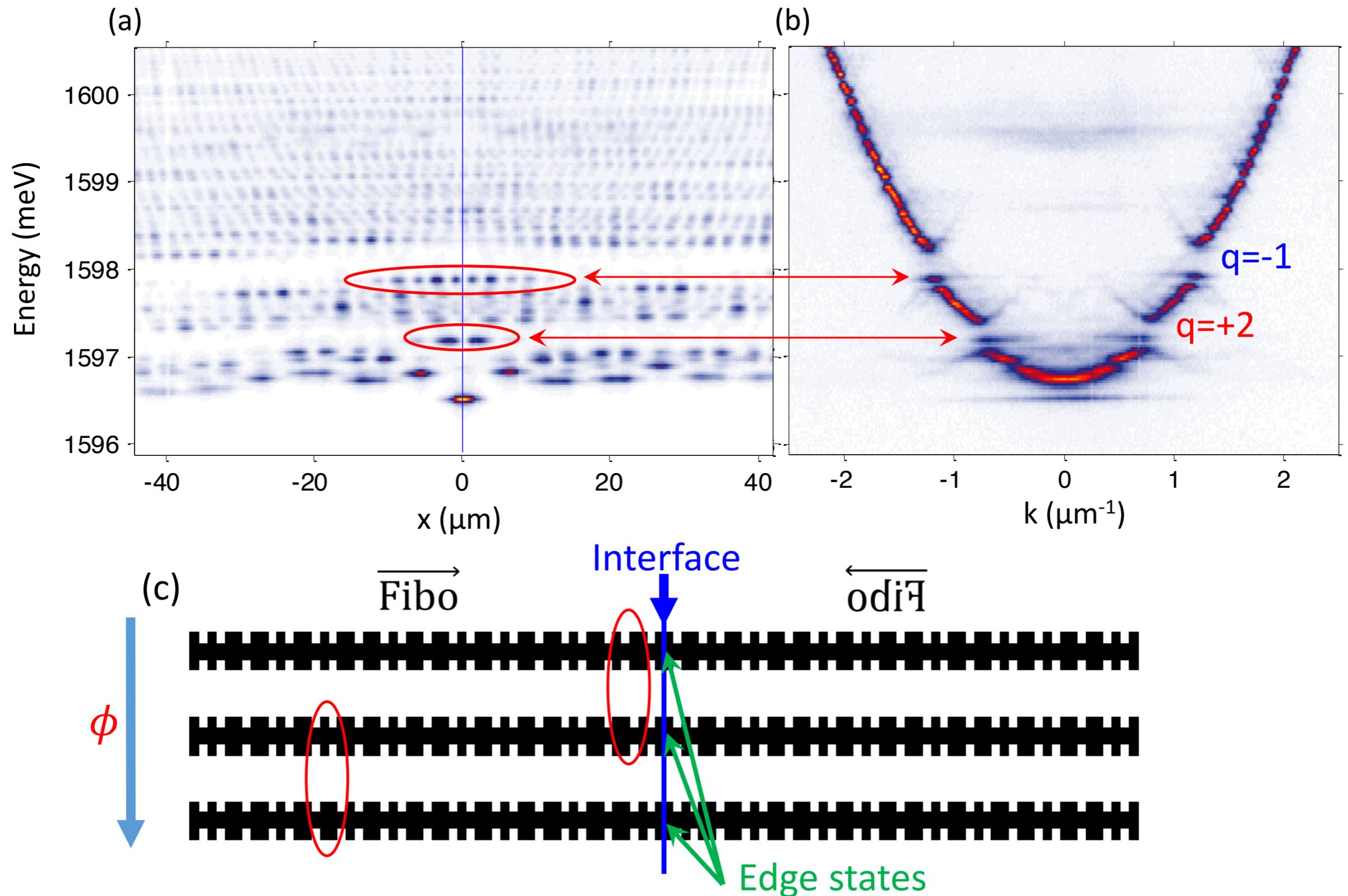
IDOS and Chiral Phase (Fibonacci)



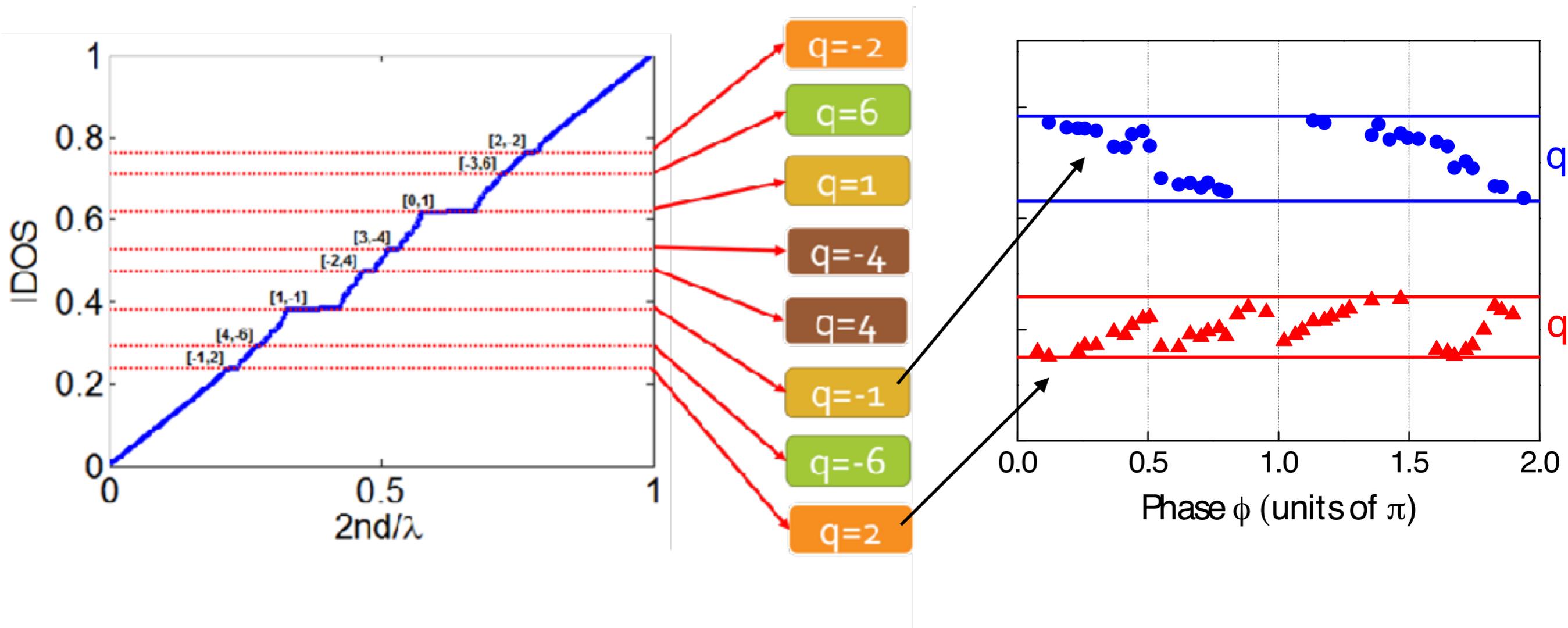
Measurement using cavity polaritons



Measurement using cavity polaritons



Measurement using cavity polaritons



F. Baboux, E. Levy, J. Bloch, E.A, 2016

Winding Relations

Two windings dependent on the same phason ϕ

Is there a relation?

Winding Relation

Structural – Spectral

$$2\mathcal{W}_\phi[\Theta] = 2q = \mathcal{W}_\phi[\alpha]$$

A Bloch Theorem ?

Two Phases - Winding numbers

- The Structural phase

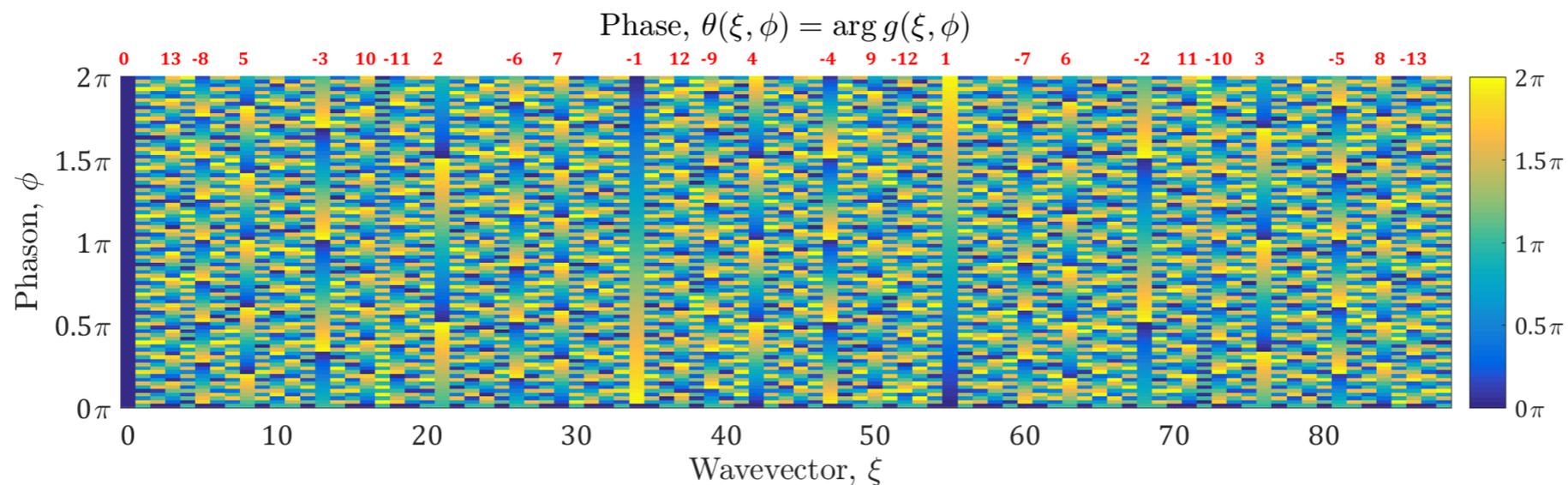
$$S(\xi, \phi) = |g(\xi, \phi)|^2, \quad \theta(\xi, \phi) = \arg g(\xi, \phi)$$

Two Phases - Winding numbers

- The Structural phase

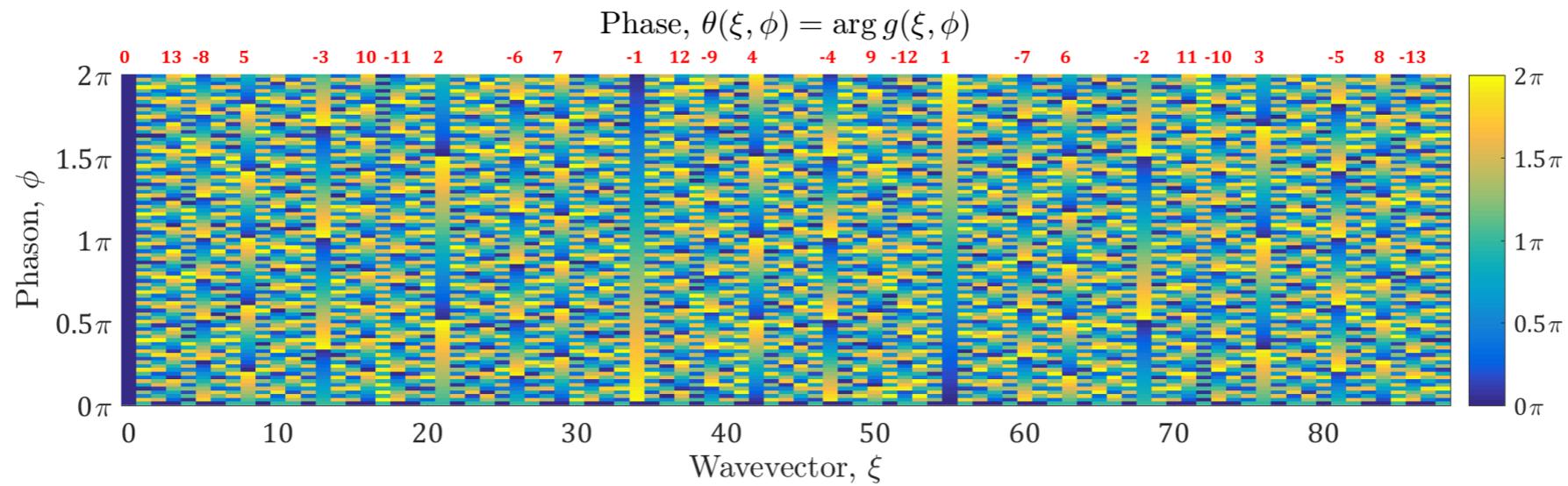
Winding number at a Bragg peak $\xi = \xi_0$

$$W_{\xi_0} = \frac{1}{2\pi} \int_0^{2\pi} \frac{\partial \theta (\xi = \xi_0, \phi)}{\partial \phi} d\phi$$



Two Phases - Winding numbers

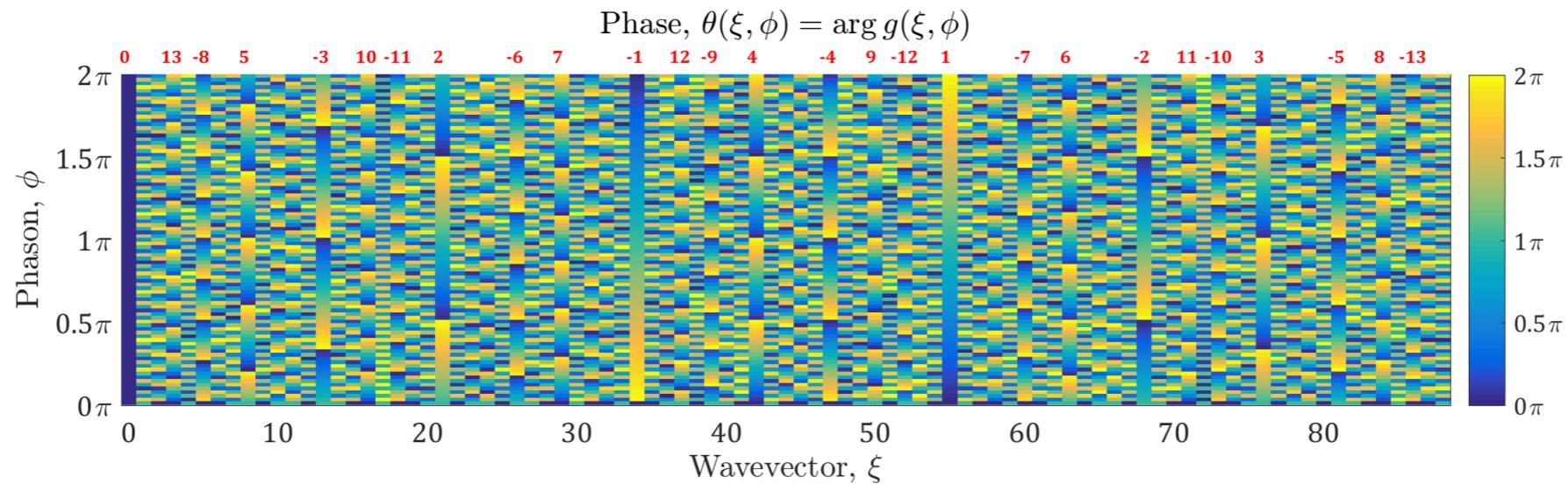
- The Structural phase



- The Chiral phase

Two Phases - Winding numbers

- The Structural phase



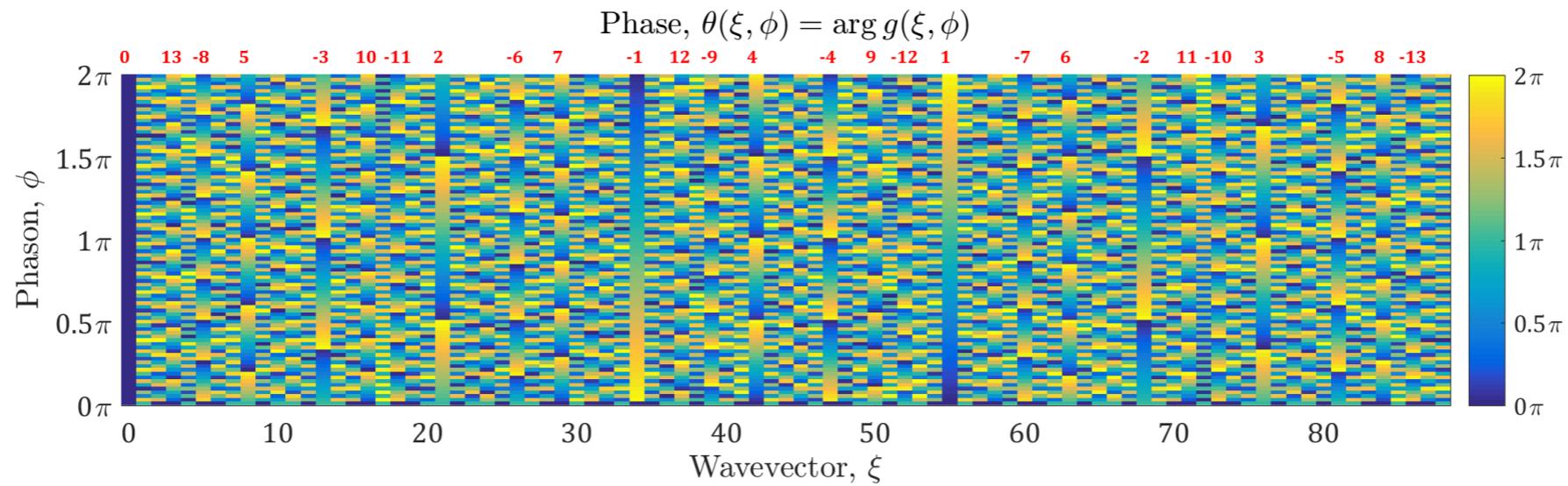
- The Chiral phase

Winding number at a spectral gap $k_{p,q}$

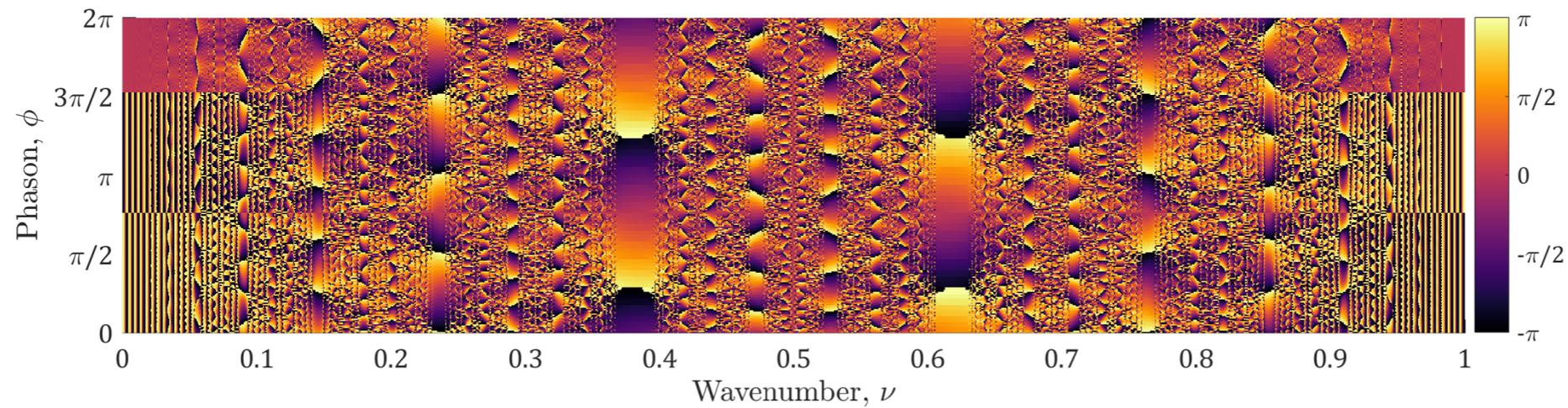
$$W_{\alpha_g} = \frac{1}{2\pi} \int_0^{2\pi} \frac{\partial \alpha(k = k_{p,q}, \phi)}{\partial \phi} d\phi$$

Two Phases - Winding numbers

- The Structural phase

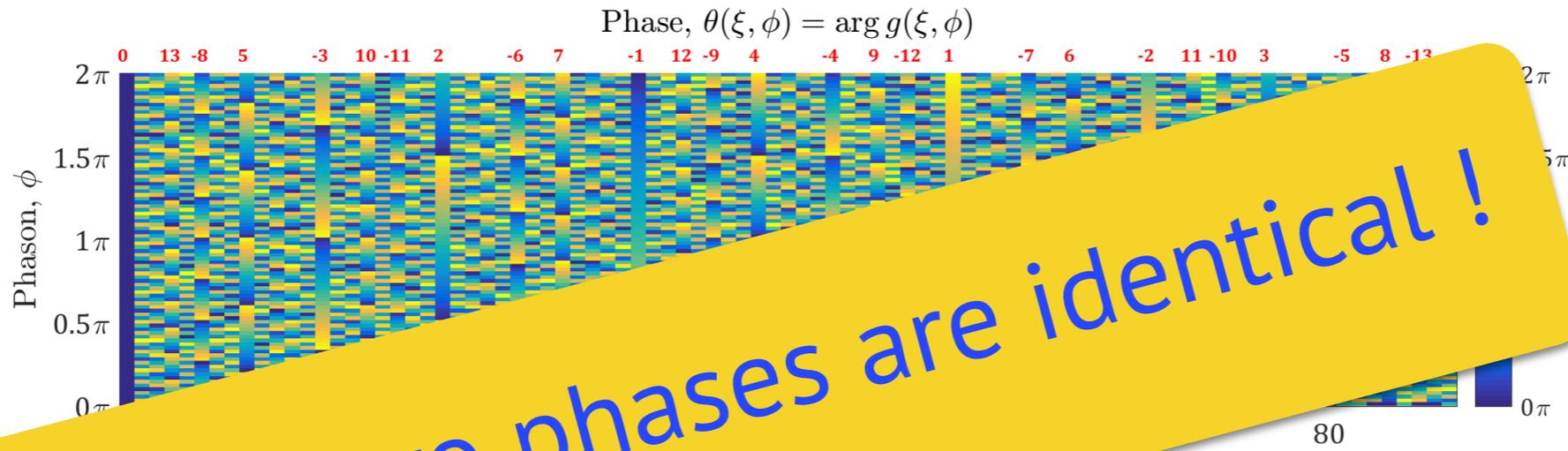


- The Chiral phase

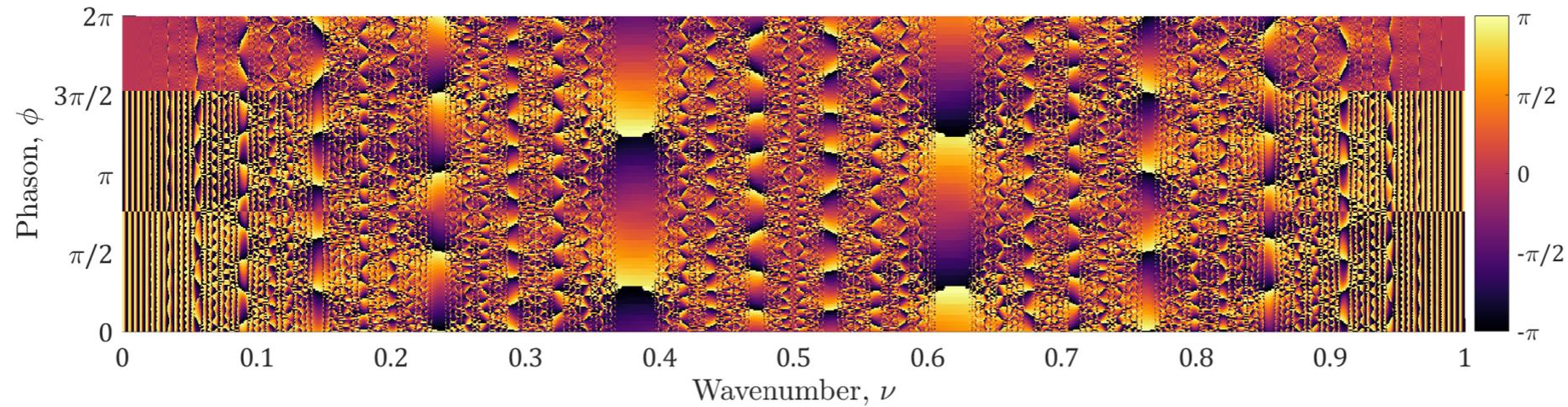


Two Phases - Winding numbers

- The Structural phase

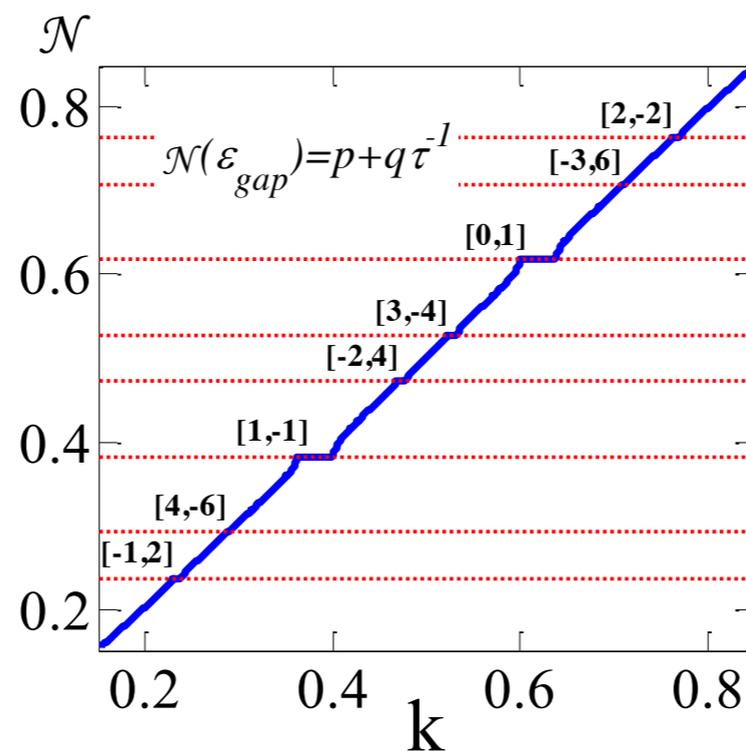
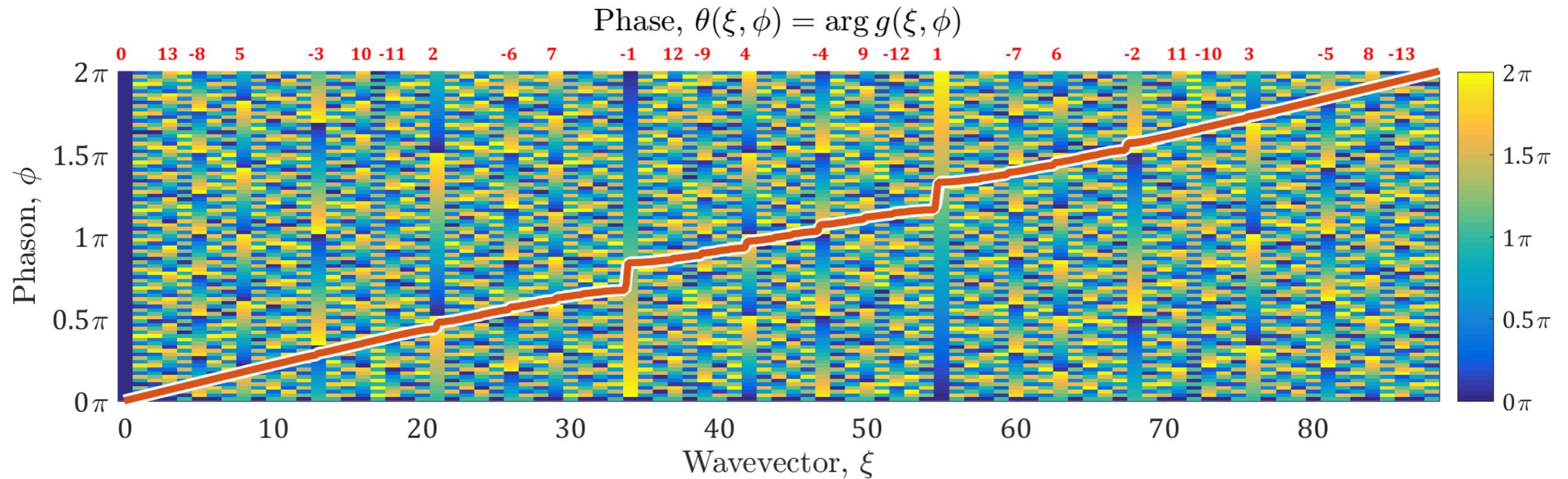


- The Structural phase



In case you are not yet convinced...

- The Structural phase



This establishes a relation between structure and spectrum.

A Bloch theorem for aperiodic tilings

$$k_b/k_0 = p + qs = \mathcal{N}(E_g), \quad p, q \in \mathbb{Z}.$$

$$2\mathcal{W}_\phi[\Theta_d] = \mathcal{W}_\phi[\Theta_s] = 2q$$

Structure

Spectrum

$$\begin{array}{ccc}
 \mathbb{Z} \cong \mathcal{W}_\phi[\Theta] & \xleftarrow{\|\cdot\|} & \mathcal{W}_\phi[\alpha] \cong \mathbb{Z} \\
 \uparrow & & \uparrow \\
 \mathbb{Z}^2 \cong \check{H}_{C\&P}^1 & \xleftarrow{\quad} & K_0^{C\&P} \cong \mathbb{Z} \oplus \mathbb{Z} \\
 \downarrow \tau_*^{\check{H}} & & \downarrow \tau_*^K \\
 \mathbb{Z} \oplus s\mathbb{Z} & \xleftrightarrow{=} & \mathbb{Z} \oplus s\mathbb{Z}
 \end{array}$$

- There is a topological content
 - independent of ϕ
- Our topological invariant has a name: the Čech Cohomology \check{H}^1
- Computable for many different tilings

Outline

- 1 Prologue
- 2 Cut and Project Tilings and Windings
- 3 Substitution Tilings and Čech Cohomology**
- 4 Bloch Theorem for Aperiodic Tilings
- 5 Topological Phase Transitions
in Fractals and Random Tilings
- 6 Epilogue

Substitutions

Periodic



- A simple rule:
$$\begin{cases} \sigma(a) = ab \\ \sigma(b) = ab \end{cases}$$

- ⊗ Resulting in...
 $a \mapsto ab \mapsto abab \mapsto$
 $abab abab \mapsto$
 $abab abab abab abab \mapsto \dots$

Fibonacci



- A simple rule:
$$\begin{cases} \sigma(a) = ab \\ \sigma(b) = a \end{cases}$$

- ⊗ Resulting in...
 $a \mapsto ab \mapsto aba \mapsto$
 $abaab \mapsto abaab aba \mapsto$
 $abaab aba abaab \mapsto \dots$

- Define substitution rules by

$$\begin{cases} \sigma(a) = a^\alpha b^\beta \\ \sigma(b) = a^\gamma b^\delta \end{cases} \Leftrightarrow \begin{cases} a \mapsto a^\alpha b^\beta \\ b \mapsto a^\gamma b^\delta \end{cases}$$

with $\alpha, \beta, \gamma, \delta \geq 0$.

- Acting on a word $w = l_1 l_2 \dots l_k$ is a **concatenation**

$$\sigma(w) = \sigma(l_1) \sigma(l_2) \dots \sigma(l_k)$$

- Associated occurrence matrix

$$M = \begin{pmatrix} \alpha & \beta \\ \gamma & \delta \end{pmatrix}$$

- Consider only **primitive** matrices (substitutions)
 - $\exists N_0$ such that $\forall N > N_0$ all entries of M^N are strictly positive
 - Largest eigenvalue $\lambda_1 > 1$ (Perron-Frobenius)
 - Left and right first eigenvectors are strictly positive

Examples

| Name | Rule | | M | λ_* |
|----------------------------|-----------------|----------------|--|-------------------------------|
| Periodic | $a \mapsto ab$ | $b \mapsto ab$ | $\begin{pmatrix} 1 & 1 \\ 1 & 1 \end{pmatrix}$ | 2 |
| Thue-Morse | $a \mapsto ab$ | $b \mapsto ba$ | $\begin{pmatrix} 1 & 1 \\ 1 & 1 \end{pmatrix}$ | 2 |
| Fibonacci | $a \mapsto ab$ | $b \mapsto a$ | $\begin{pmatrix} 1 & 1 \\ 1 & 0 \end{pmatrix}$ | $\tau = \frac{\sqrt{5}+1}{2}$ |
| Fibonacci ² | $a \mapsto aab$ | $b \mapsto ab$ | $\begin{pmatrix} 2 & 1 \\ 1 & 1 \end{pmatrix}$ | τ^2 |
| Non-Fibonacci ² | $a \mapsto aab$ | $b \mapsto ba$ | $\begin{pmatrix} 2 & 1 \\ 1 & 1 \end{pmatrix}$ | τ^2 |

- Representatives of 3 families:

① Periodic ② Quasiperiodic ③ Aperiodic

Occurrence Matrix M

What can be done?

- Gap labeling Thm.

What cannot be done?

- Diffraction

How to calculate the Čech Cohomology \check{H}^1 ?

Supertiles ($1D$)

Infinite tiling: $w_\infty = \sigma^\infty(a)$

Supertiles (words): $\Gamma_n = \{w \in w_\infty \mid |w| = n\}$

Supertile rule: $\sigma_n : \Gamma_n \rightarrow \Gamma_n^{\mathbb{N}}$

Occurrence mat.: M_n (all with the same λ_*)

Example: Fibonacci, $n = 2$

$$\Gamma_1 = \{a, b\}$$

$$\sigma_1 = \begin{cases} a \mapsto ab \\ b \mapsto a \end{cases}$$

$$M_1 = \begin{pmatrix} 1 & 1 \\ 1 & 0 \end{pmatrix}$$

$$w_\infty^{(1)} = abaab\ aba\ abaab$$

$$\Gamma_2 = \{A, B, C\} = \{aa, ab, ba\}$$

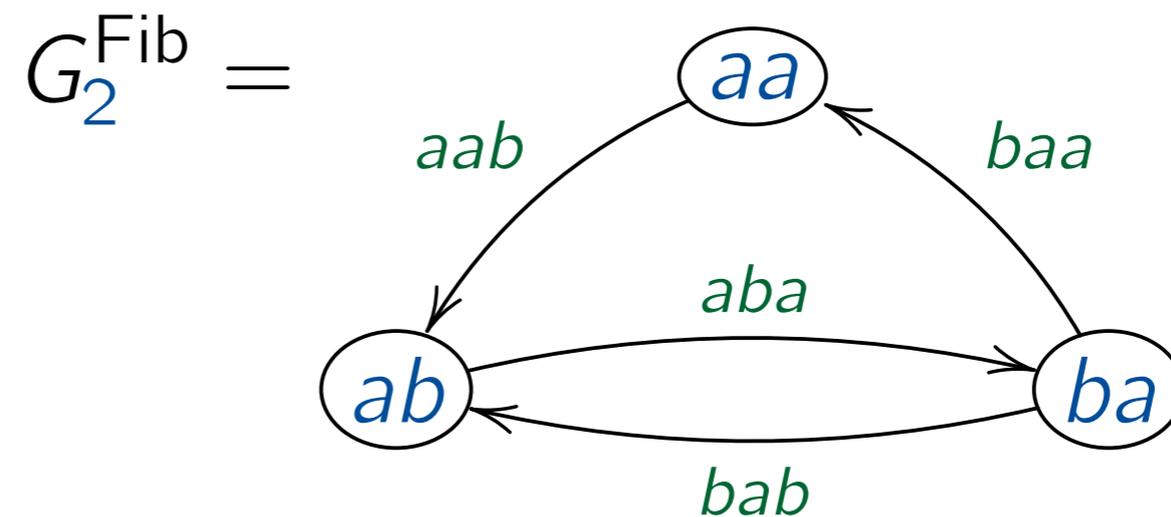
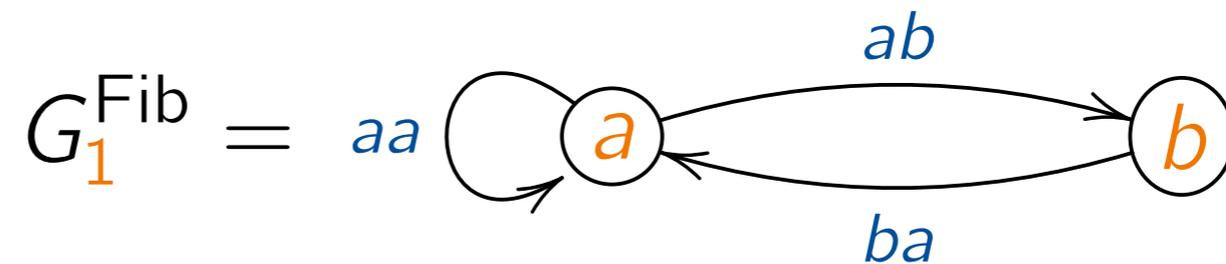
$$\sigma_2 = \begin{cases} A \mapsto BC \\ B \mapsto BC \\ C \mapsto A \end{cases}$$

$$M_2 = \begin{pmatrix} 0 & 1 & 1 \\ 0 & 1 & 1 \\ 1 & 0 & 0 \end{pmatrix}$$

$$w_\infty^{(2)} = BCABC\ BCA\ BCABC$$

Supertiles ($1D$)

- Shift-maps $\gamma_n(L_i) = L_j$ if $L_i L_j$ exists in $w_\infty^{(n)}$
- Representation by planar graphs G_n
(called Bratteli diagrams)



How to calculate the Čech Cohomology \check{H}^1 ?

$$\zeta(z) = \frac{\det(I - zA_0^\top)}{\det(I - zA_1^\top)} \doteq \frac{p_0(z)}{p_1(z)}.$$

$$p_k(z) = \prod_{i=1}^I (1 - c_i z) \prod_{j=1}^J (1 - d_j z - e_j z^2) \quad c_i, d_j, e_j \in \mathbb{Z}$$

$$\begin{aligned} \check{H}^k &\cong \bigoplus_{i=1}^I \mathbb{Z}[1/c_i] \oplus \bigoplus_{j=1}^J \mathbb{Z}^2[1/e_j] \\ &= \mathbb{Z}[c_1^{-1}] \oplus \cdots \oplus \mathbb{Z}[c_I^{-1}] \oplus \mathbb{Z}^2[e_1^{-1}] \oplus \cdots \oplus \mathbb{Z}^2[e_J^{-1}] \end{aligned}$$

$$\mathbb{Z}[1/c] = \{n/c^m \mid n, m \in \mathbb{Z}\}$$

How to calculate the Čech Cohomology \check{H}^1 ?

- Computable for many different tilings
- Distinguishes between families

| Family | \check{H}^1 | Diffraction peaks | Gap labeling |
|---------------|---|--------------------------------------|---|
| Periodic | \mathbb{Z} | $k_b = n/2$ | $\mathcal{N}_g = 1/2$ |
| Quasiperiodic | \mathbb{Z}^2 | $k_b = p + q \varrho_b$ | $\mathcal{N}_g = q \varrho_b$ |
| Thue-Morse | $\mathbb{Z} \oplus \mathbb{Z} \begin{bmatrix} 1 \\ 2 \end{bmatrix}$ | $k_b = \frac{1}{2n+1} \frac{m}{2^N}$ | $\mathcal{N}_g = \frac{1}{3} \frac{m}{2^N}$ |

Gap Labeling Theorem

- In 1D aperiodic substitutions, the possible gaps are

$$\mathcal{N}_{\text{gap}} \in \tau_*^K [K_0(\mathcal{B})]$$

- Explicitly ($k, N \in \mathbb{N}$),

$$\mathcal{N}_{\text{gap}} = \frac{1}{a} \frac{k}{\lambda_*^N} \pmod{1}$$

- The normalization factor a is inferred by $\mathbf{v}_*, \mathbf{v}_*^{(2)}$
- In C&P tilings ($p, q \in \mathbb{Z}$)

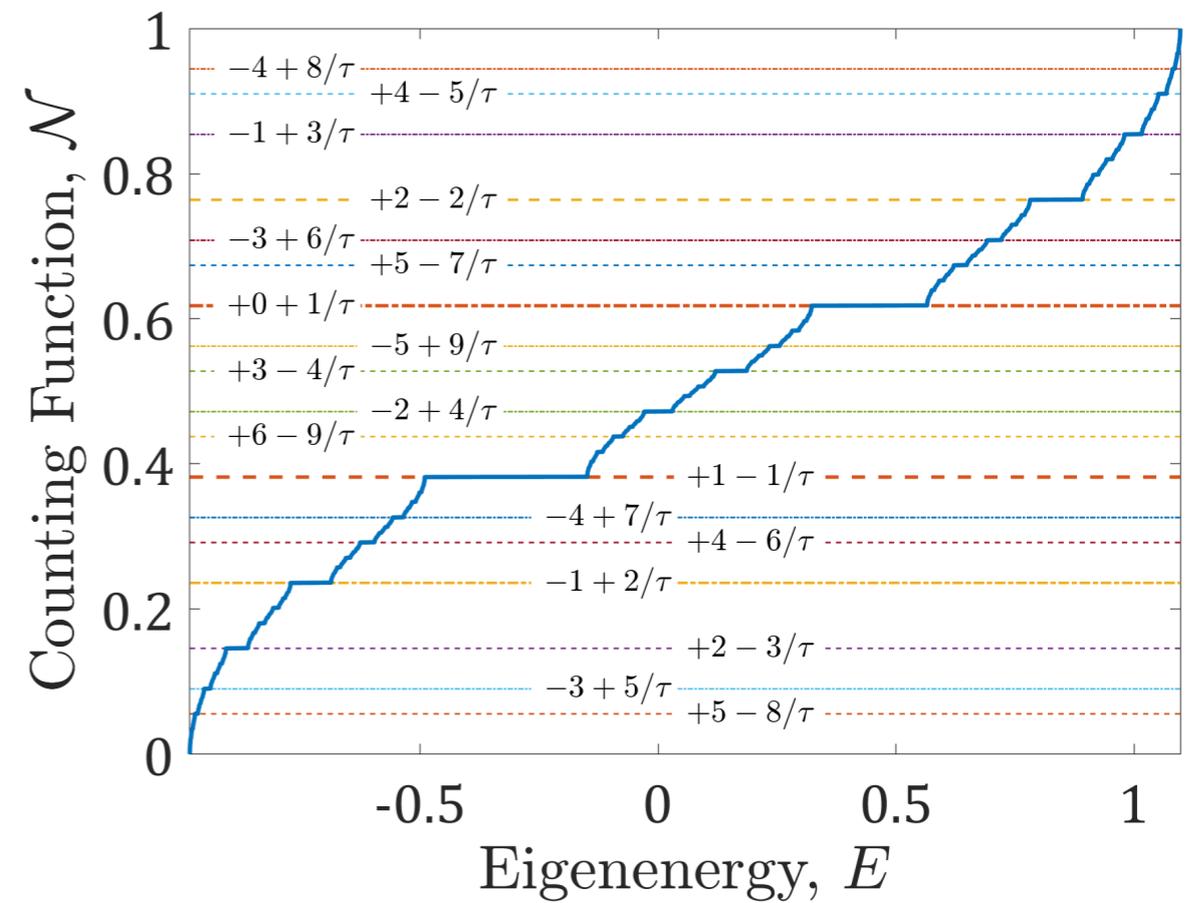
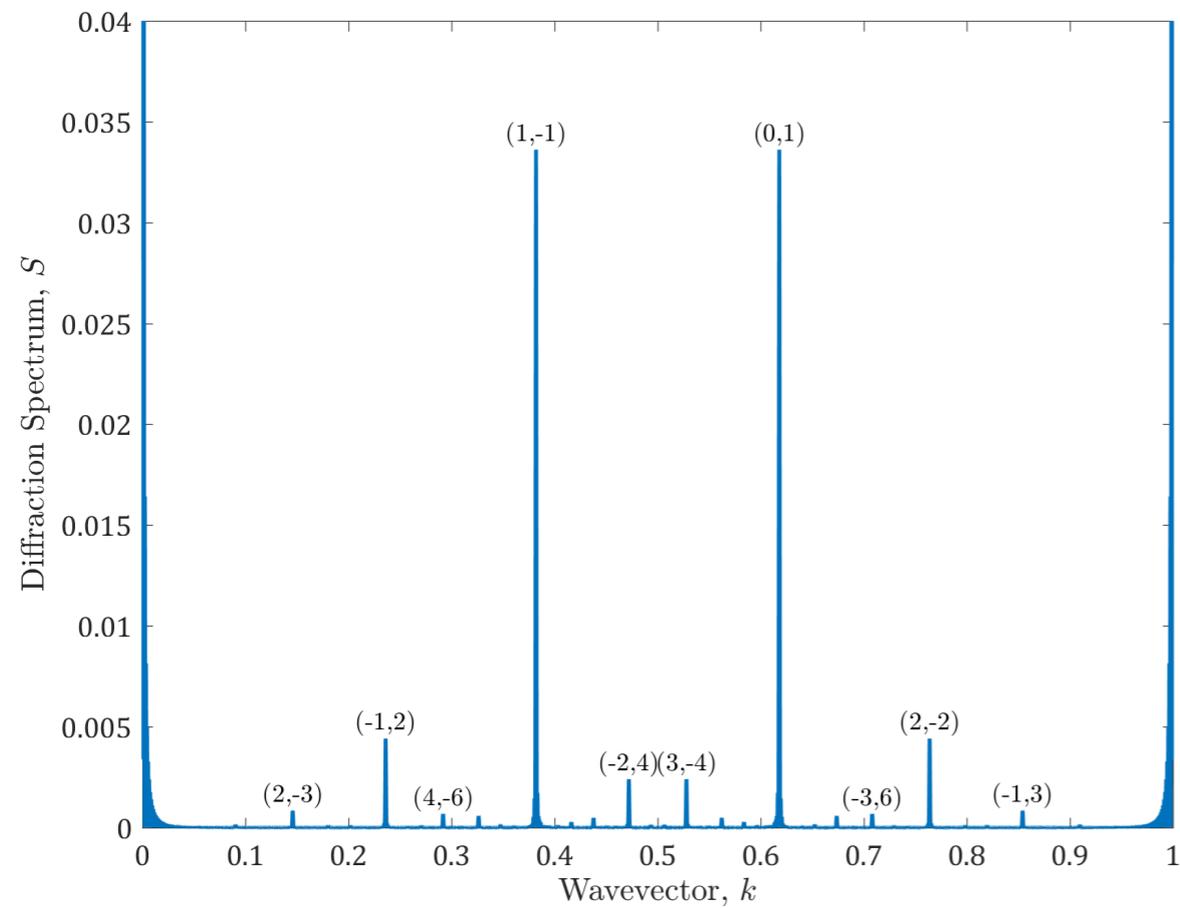
$$\mathcal{N}_{\text{gap}} = p + q s \pmod{1}$$

Outline

- 1 Prologue
- 2 Cut and Project Tilings and Windings
- 3 Substitution Tilings and Čech Cohomology
- 4 Bloch Theorem for Aperiodic Tilings**
- 5 Topological Phase Transitions
in Fractals and Random Tilings
- 6 Epilogue

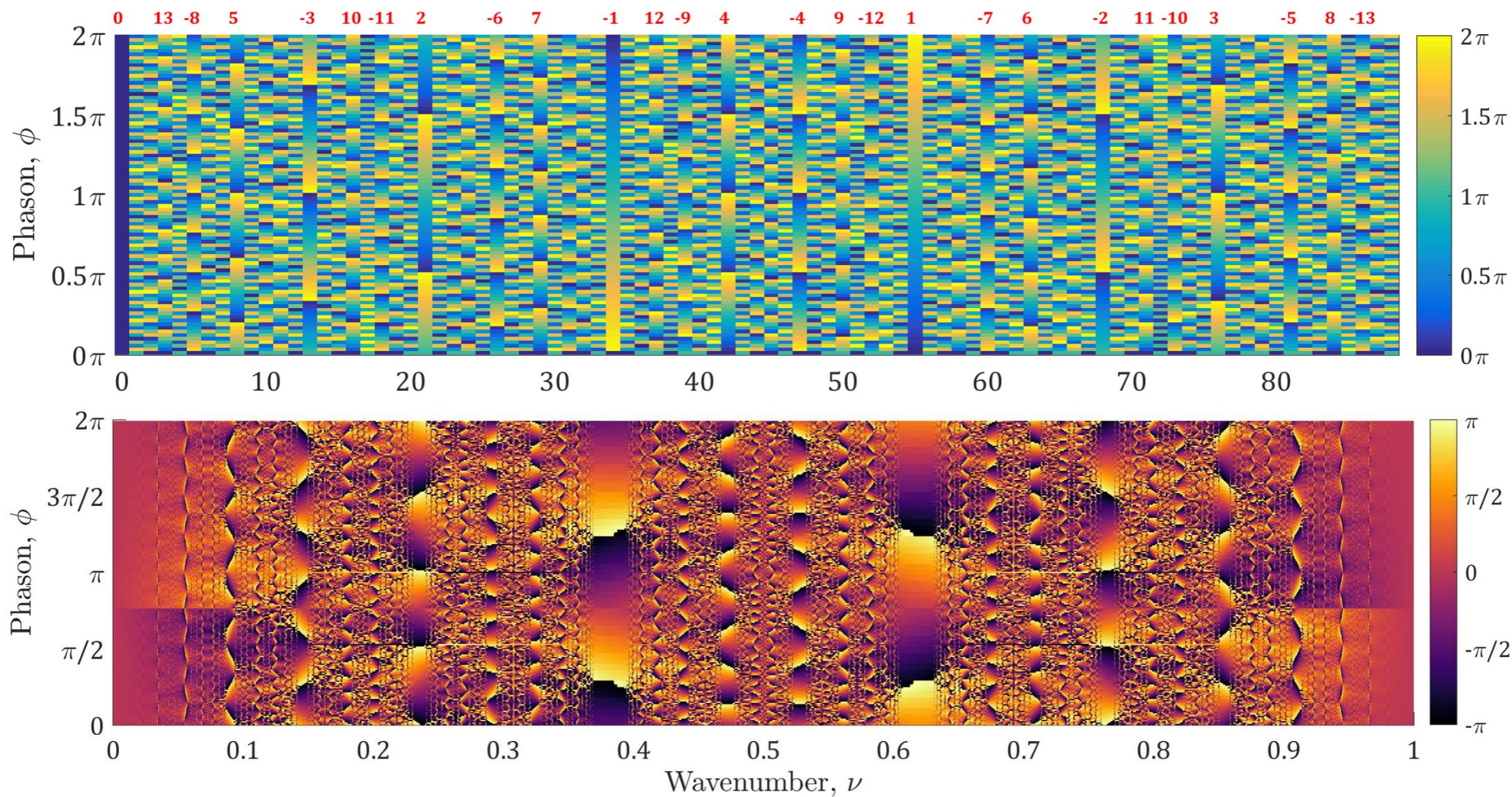
Symmetry:

$$k_b = p + q s = \mathcal{N}(E_g)$$



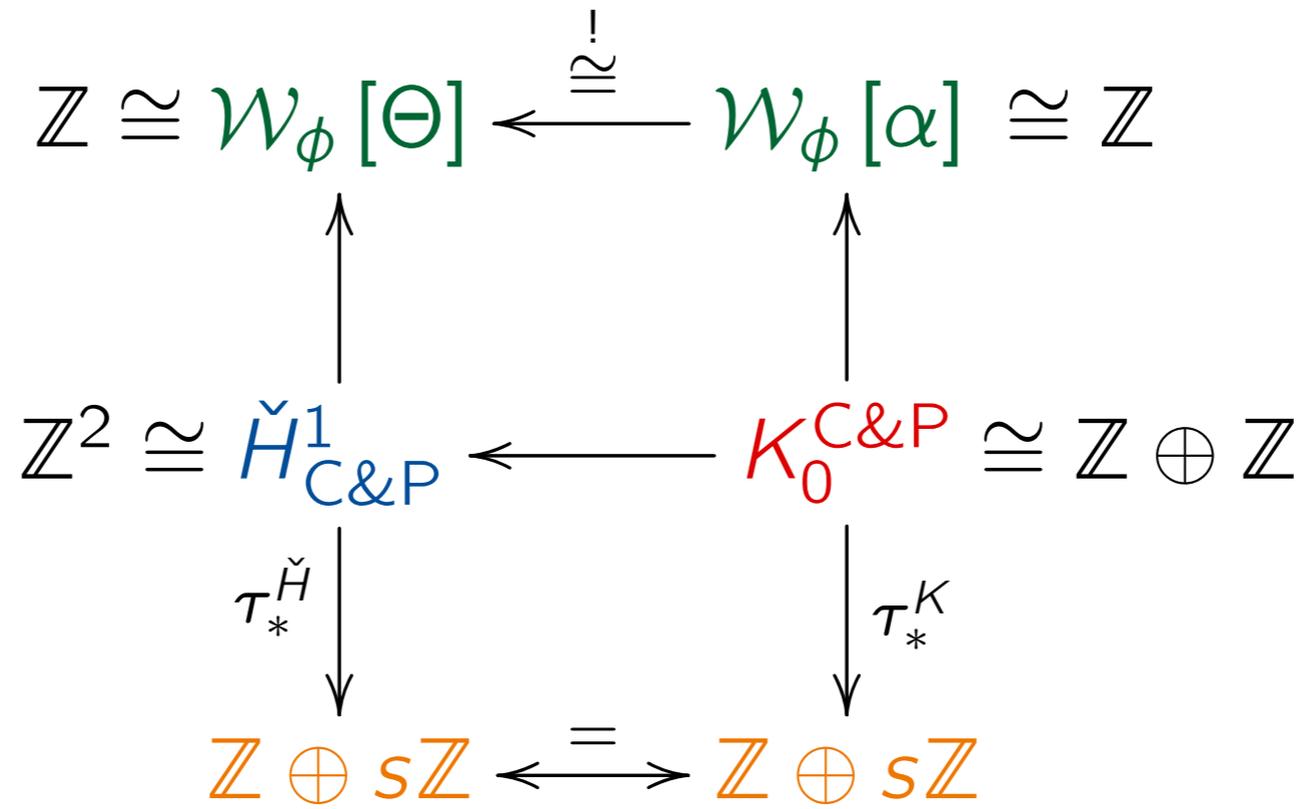
Windings:

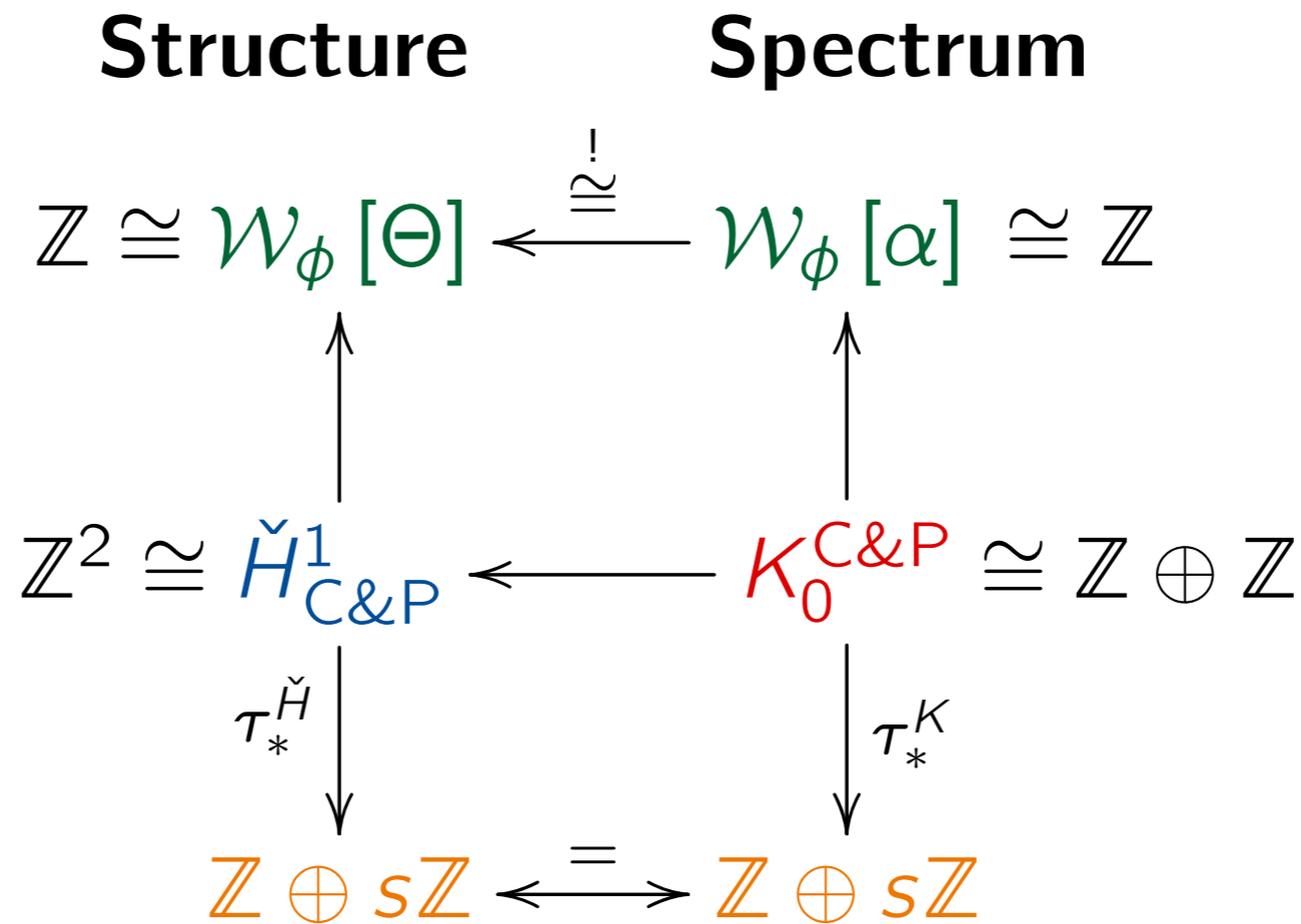
$$2\mathcal{W}_\phi [\Theta] = 2q = \mathcal{W}_\phi [\alpha]$$



Structure

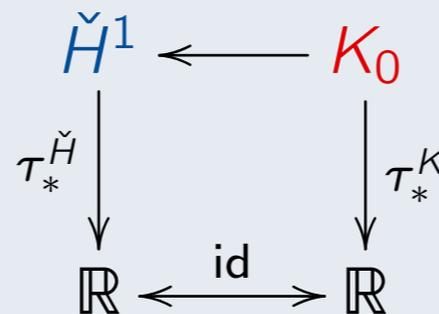
Spectrum





Theorem (Generalized Bloch)

For finite local complexity tilings with finitely many tile orientations, the following diagram commutes in dimensions $D \leq 3$



1. E. Akkermans, Y. Don, J. Rosenberg and C. L. Schochet, *Relating Diffraction and Spectral Data of Aperiodic Tilings: Towards a Bloch theorem*, *J. Geom. Phys.* **165**, 104217 (2021).

Implications

- Since the traces $\tau_*^{\check{H}}(\check{H}^1)$ and $\tau_*^K(K_0)$ commute, \check{H}^1 represents both spectral and structural properties
 - recall $\tau_*^K(K_0)$ is the GLT
- The trace $\tau_*^{\check{H}}(\check{H}^1)$ does not describe diffraction $S(k)$
 - except when $S(k)$ consists of Bragg peaks only

Generalized Bloch – Summary

| Family | \check{H}^1 | Diffraction peaks | | $\tau_*^{\check{H}}(\check{H}^1)$ | Spectral Gaps |
|-----------------|--|--|-------|---------------------------------------|---|
| Periodic | \mathbb{Z} | $k_n = n$ | PP | \mathbb{Z} | $\mathcal{N} = \text{const}$ |
| Fibonacci | \mathbb{Z}^2 | $k_{p,q} = p + q/\tau$ | PP | $\mathbb{Z} + \tau^{-1}\mathbb{Z}$ | $\mathcal{N}_q = q/\tau$ |
| Thue-Morse | $\mathbb{Z} \oplus \mathbb{Z}[\frac{1}{2}]$ | $k_{n,m,N} = \frac{1}{2n+1} \frac{m}{2^N}$ | SC+PP | $\frac{1}{3} \mathbb{Z}[\frac{1}{2}]$ | $\mathcal{N}_{m,N} = \frac{1}{3} \frac{m}{2^N}$ |
| Period Doubling | $\mathbb{Z} \oplus \mathbb{Z}[\frac{1}{2}]$ | $k_{m,N} = \frac{m}{2^N}$ | PP | $\frac{1}{3} \mathbb{Z}[\frac{1}{2}]$ | $\mathcal{N}_{m,N} = \frac{1}{3} \frac{m}{2^N}$ |
| Rudin-Shapiro | $\mathbb{Z} \oplus \mathbb{Z}[\frac{1}{2}] \oplus \mathbb{Z}^2[\frac{1}{2}]$ | N/A | AC | $\mathbb{Z}[\frac{1}{2}]$ | $\mathcal{N}_{m,N} = \frac{m}{2^N}$ |

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Summary I

- C&P tilings exhibit topological content by **winding numbers**
 - a diff. peaks and spectral gaps are analogous: $k_b = p + q s = \mathcal{N}(\nu_g)$
 - b structural and spectral windings are related: $2\mathcal{W}_\phi[\Theta] = 2q = \mathcal{W}_\phi[\alpha]$
 - c verified **experimentally**
- Aperiodic tilings are fully characterized by a topological invariant – the **Čech cohomology** \check{H}^1
 - a the right mathematical tool to answer physical questions
 - b allowing to: characterize tilings and count tiles; enumerate **Bragg** peaks; label **spectral gaps**
- The **Bloch theorem** is generalized to FLC tilings by \check{H}^1
 - a highlights the connection b/w structural & spectral features of tilings
 - b furthermore, for C&P tilings, relates structural and spectral **windings**
 - c for non-C&P tilings, $\tau_*^{\check{H}}(\check{H}^1)$ is unrelated to diffraction $S(k)$

Summary II

- A new description of **diffraction** using Bratteli diagrams
 - a Can be calculated for tilings with **Bragg diffraction** spectrum
 - b Closely related to **windings** on Bratteli diagrams
- Diffraction of **Thue-Morse** tiling is carefully analyzed
 - a Characterization of **peaks** by their **growth rate**
 - b Inconclusive **experimental results**
- Innovative portrayal of **fractals** employing tilings and substitutions
 - a Novel **Gap Labeling Conjecture** for fractals is presented
- **Topological phase transitions** are found
 - a By **flux** in fractals
 - b By **random** substitution rules

Prospect

Future

- Bloch Theorem
 - Extend to include windings
 - Explore in dimensions ≥ 4
- Windings
 - Identify the topological numbers for all $1D$ tilings
 - Explore in $2D$ and $3D$
- Calculate diffraction $S(k)$ using Bratteli diagrams
 - for all $1D$ tilings
 - extend to $2D$ and beyond
- Fractals
 - Prove the Gap Labeling Conjecture
 - Identify the proper \check{H}^1

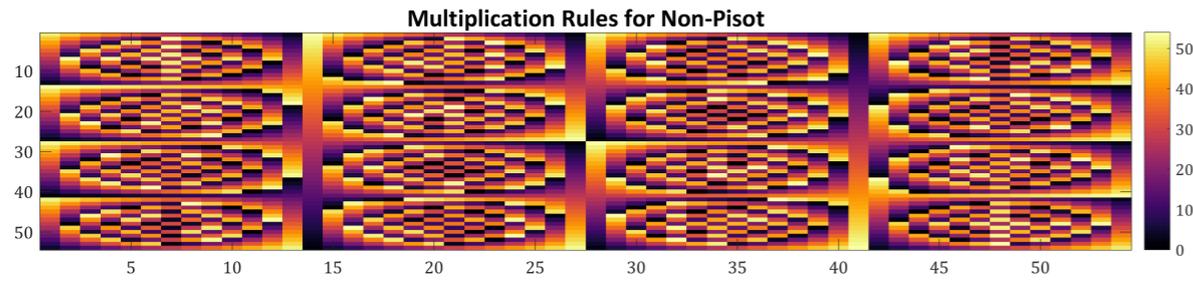


Fig. A.1: Multiplication matrix $\Lambda(q, r)$ for the Non-Pisot substitution.

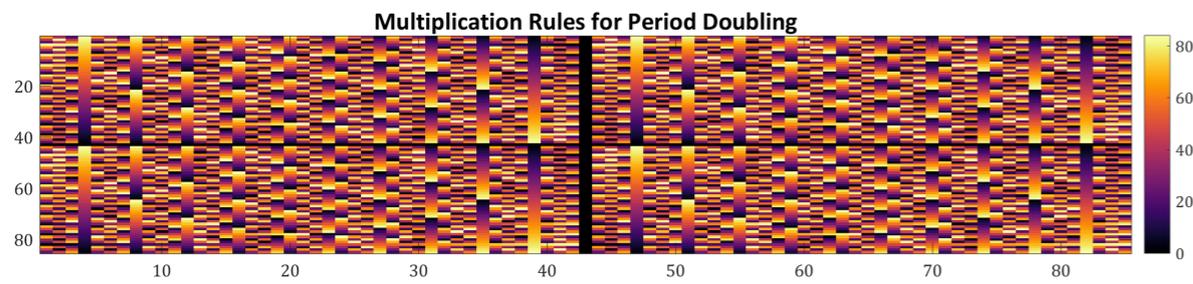


Fig. A.2: Multiplication matrix $\Lambda(q, r)$ for the Period Doubling substitution.

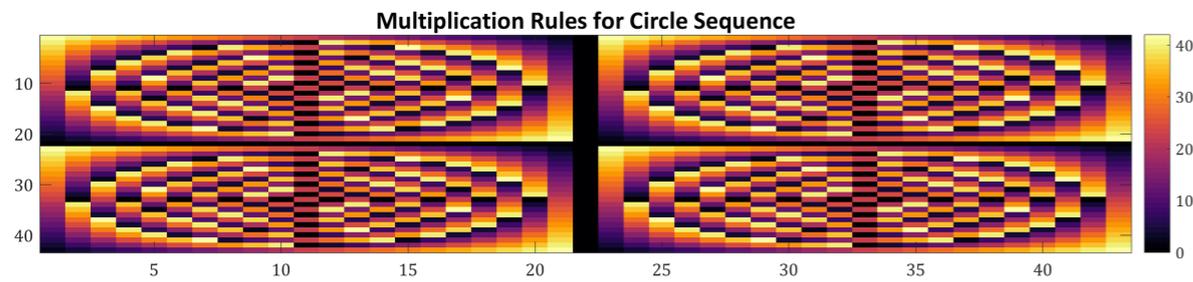


Fig. A.3: Multiplication matrix $\Lambda(q, r)$ for the Circle Sequence substitution.

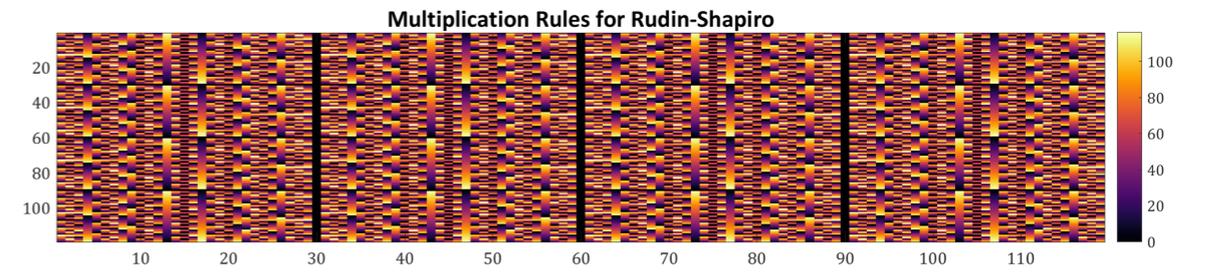


Fig. A.4: Multiplication matrix $\Lambda(q, r)$ for the Rudin-Shapiro substitution.

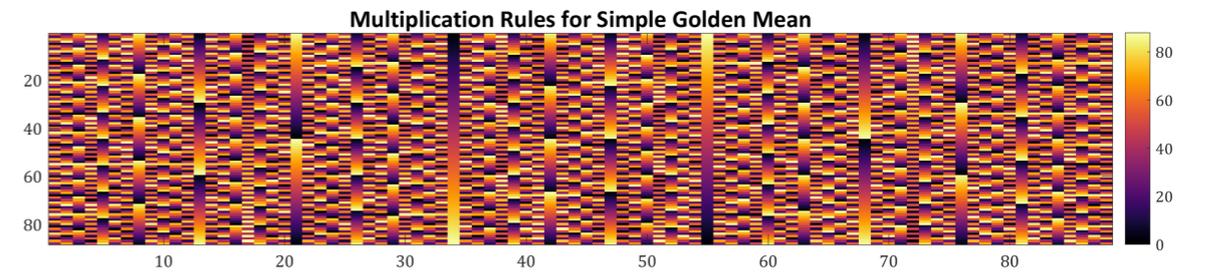


Fig. A.5: Multiplication matrix $\Lambda(q, r)$ for the Golden Mean substitution.

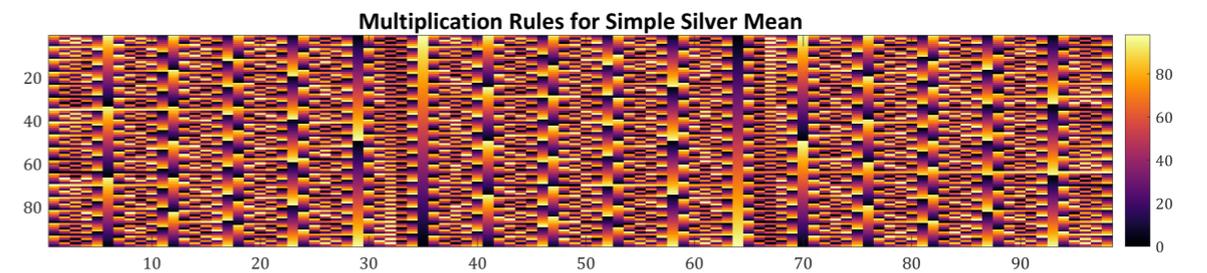


Fig. A.6: Multiplication matrix $\Lambda(q, r)$ for the Silver Mean substitution.

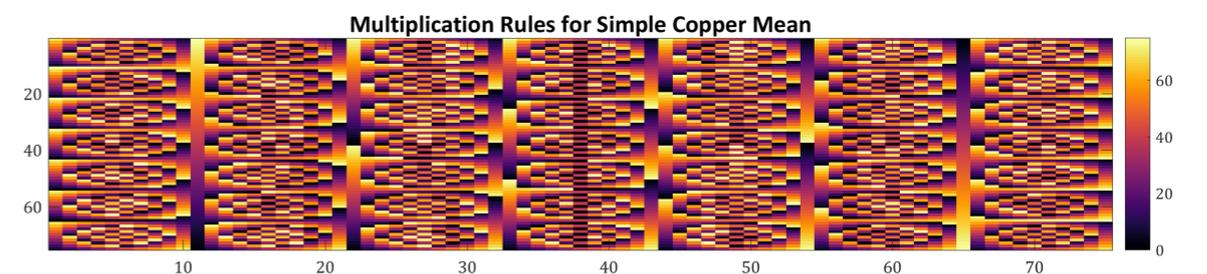


Fig. A.7: Multiplication matrix $\Lambda(q, r)$ for the Copper Mean substitution.

Thank you for your attention

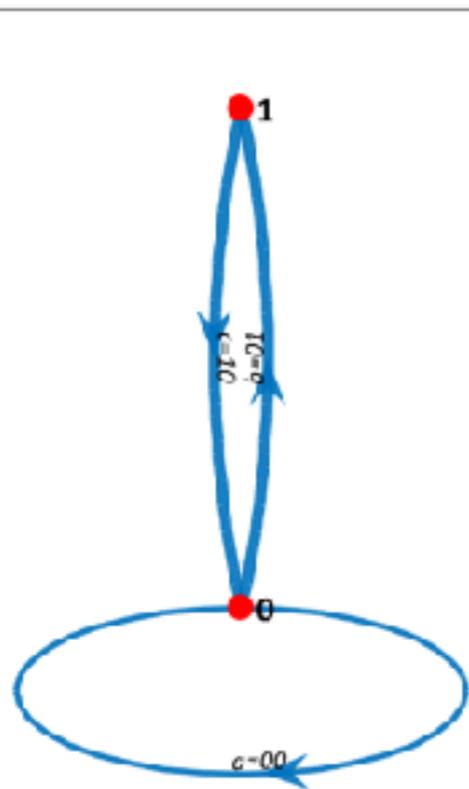
Most of results and details are
available at :

<https://phsites.technion.ac.il/eric/>

Summary - the untold part

- Given a topological meaning to the integers labelling the gaps of the fractal spectrum.
- Proposed a complete algebraic structure to account for the topological integers (Abelian group structure isomorphic to $\mathbb{Z}/F_N\mathbb{Z}$)
- This Abelian group is isomorphic to the cohomology group $H^{(1)}$ defined on (Bratelli) graphs associated to the quasi periodic structures.

Fibonacci for Letters

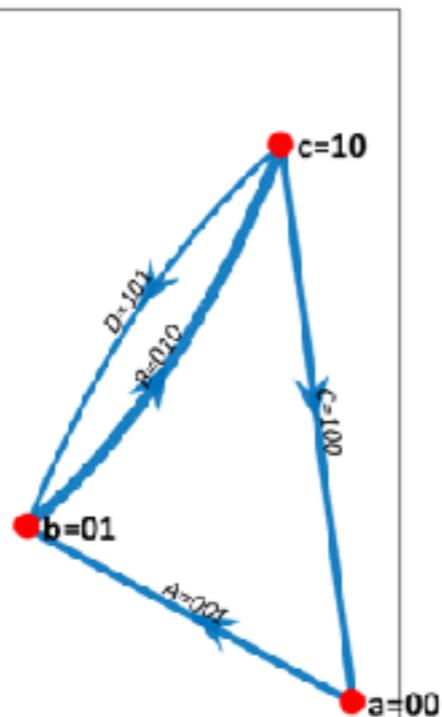


$$H^0 \cong \mathbb{Z}^1, H^1 \cong \mathbb{Z}^2$$

$$M_1 = \begin{pmatrix} 1 & 1 \\ 1 & 0 \end{pmatrix}, \quad \begin{array}{l} \sigma_1(0) = 01 \\ \sigma_1(1) = 0 \end{array}$$

$$\lambda_1^{(1)} = \tau, \quad v_1^{(1)} = (\tau - 1 \quad 2 - \tau)$$

Fibonacci for Nodes

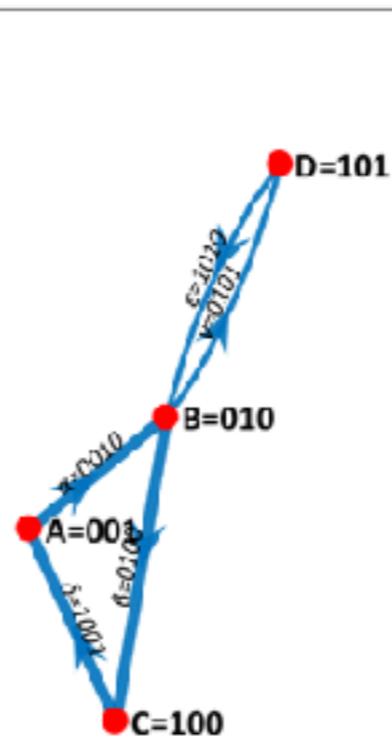


$$H^0 \cong \mathbb{Z}^1, H^1 \cong \mathbb{Z}^2$$

$$M_2 = \begin{pmatrix} 0 & 1 & 1 \\ 0 & 1 & 1 \\ 1 & 0 & 0 \end{pmatrix}, \quad \begin{array}{l} \sigma_2(a) = bc \\ \sigma_2(b) = bc \\ \sigma_2(c) = a \end{array}$$

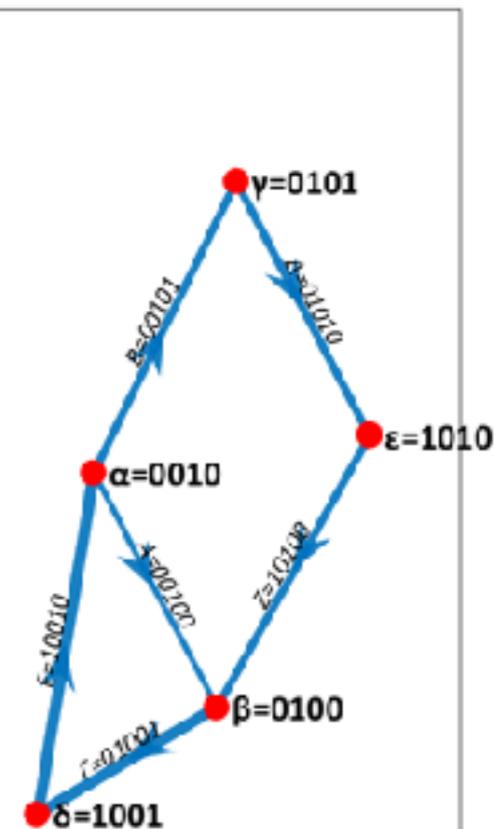
$$\lambda_1^{(2)} = \tau, \quad v_1^{(2)} = (2\tau - 3 \quad 2 - \tau \quad 2 - \tau)$$

Fibonacci for Edges



$$H^0 \cong \mathbb{Z}^1, H^1 \cong \mathbb{Z}^2$$

Fibonacci for Frames



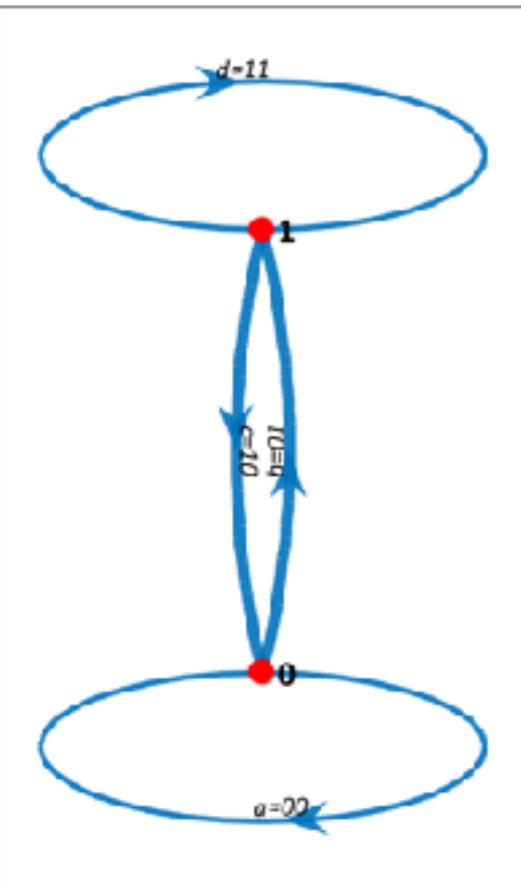
$$H^0 \cong \mathbb{Z}^1, H^1 \cong \mathbb{Z}^2$$

Gaps: $p + q \cdot \tau \cap [0, 1]$ for $p, q \in \mathbb{Z}$

Pisct ($|\lambda_2^{(1)}| = \tau - 1$), quasiperiodic

Inflation: $\zeta(z) = \frac{1 - z}{-z^2 - z + 1}$

Thue-Morse for Letters

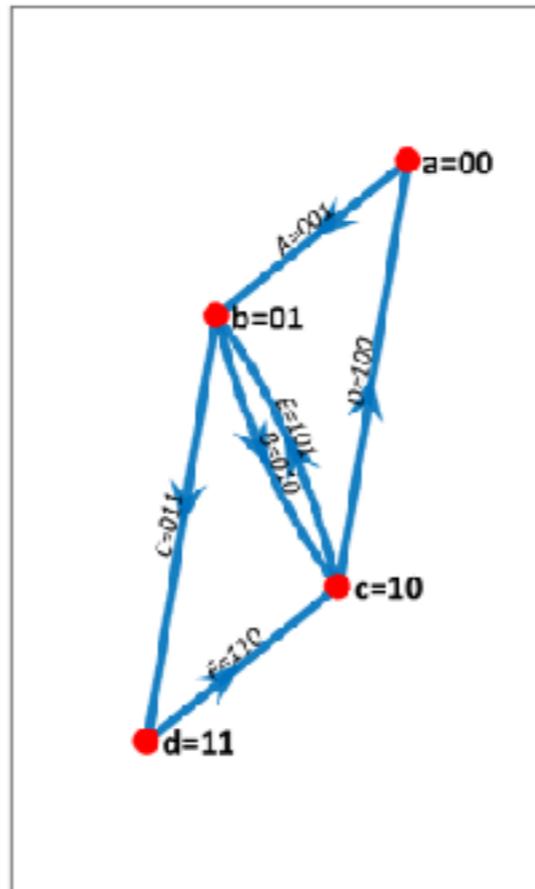


$$H^0 \cong \mathbb{Z}^1, H^1 \cong \mathbb{Z}^3$$

$$M_1 = \begin{pmatrix} 1 & 1 \\ 1 & 1 \end{pmatrix}, \quad \sigma_1(0) = 01, \quad \sigma_1(1) = 10$$

$$\lambda_1^{(1)} = 2, \quad v_1^{(1)} = \left(\frac{1}{2} \quad \frac{1}{2} \right)$$

Thue-Morse for Nodes

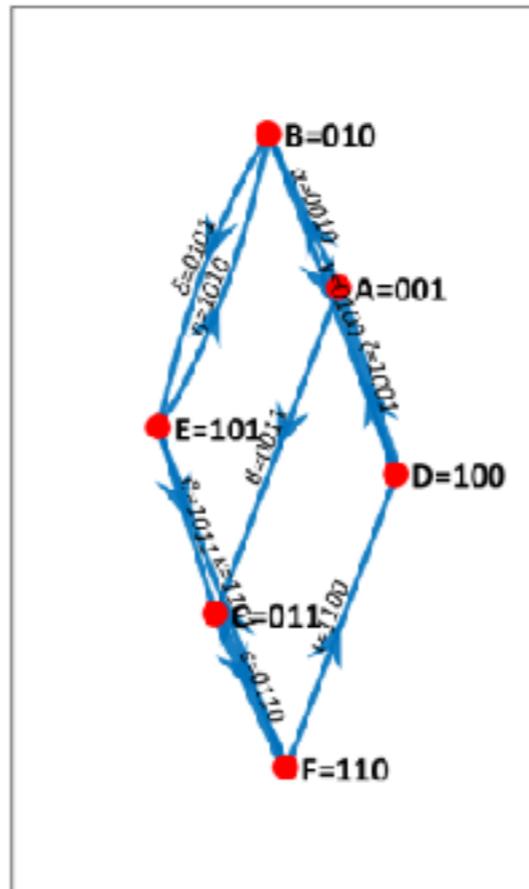


$$H^0 \cong \mathbb{Z}^1, H^1 \cong \mathbb{Z}^3$$

$$M_2 = \begin{pmatrix} 0 & 1 & 1 & 0 \\ 0 & 1 & 0 & 1 \\ 1 & 0 & 1 & 0 \\ 0 & 1 & 1 & 0 \end{pmatrix}, \quad \begin{aligned} \sigma_2(a) &= bc \\ \sigma_2(b) &= bd \\ \sigma_2(c) &= ca \\ \sigma_2(d) &= cb \end{aligned}$$

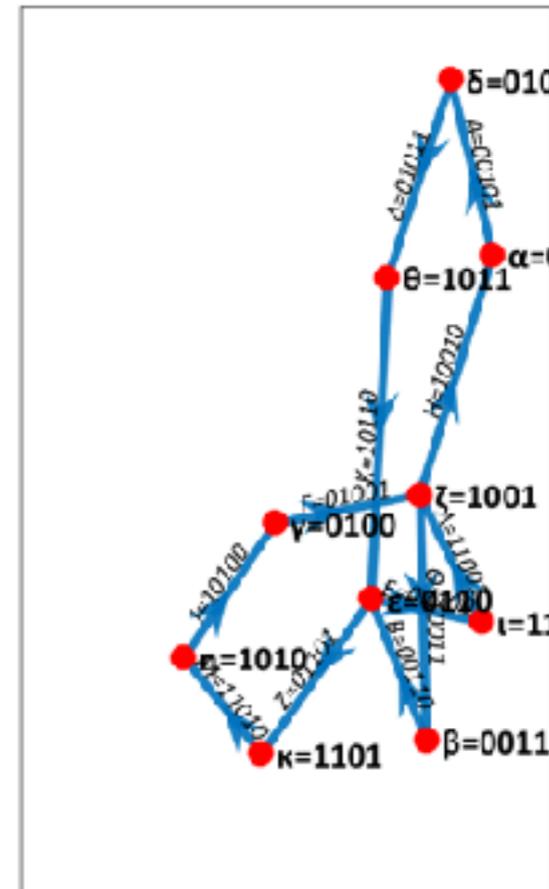
$$\lambda_1^{(2)} = 2, \quad v_1^{(2)} = \left(\frac{1}{6} \quad \frac{1}{3} \quad \frac{1}{3} \quad \frac{1}{6} \right)$$

Thue-Morse for Edges



$$H^0 \cong \mathbb{Z}^1, H^1 \cong \mathbb{Z}^5$$

Thue-Morse for Frames



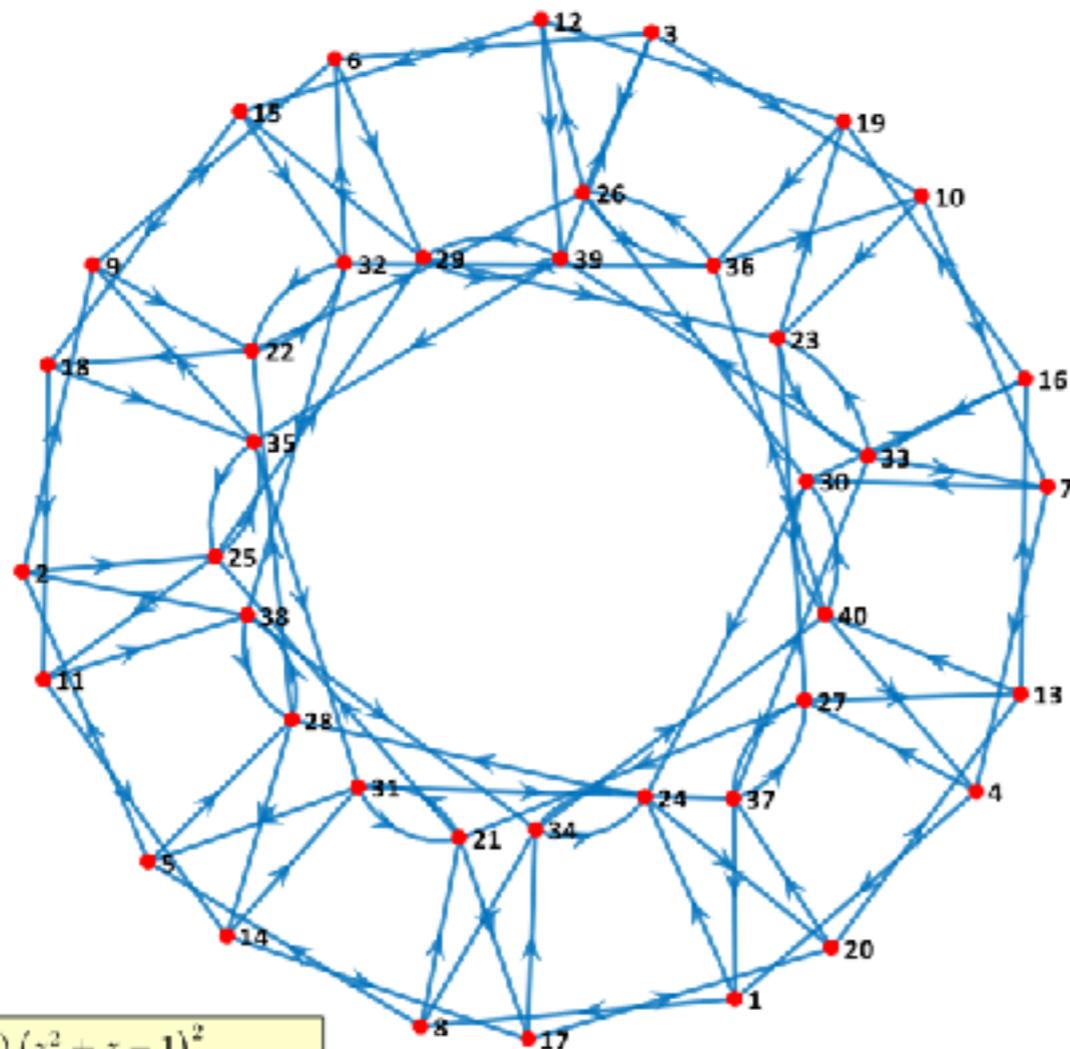
$$H^0 \cong \mathbb{Z}^1, H^1 \cong \mathbb{Z}^3$$

Caps: $\frac{k}{3 \cdot 2^N} \cap [0, 1)$ for $k, N \in \mathbb{Z}$

Pisot ($|\lambda_2^{(1)}| = 0$), aperiodic

Inflation: $\zeta(z) = \frac{1-z}{-(2z-1)(z+1)}$

Penrose Tiling

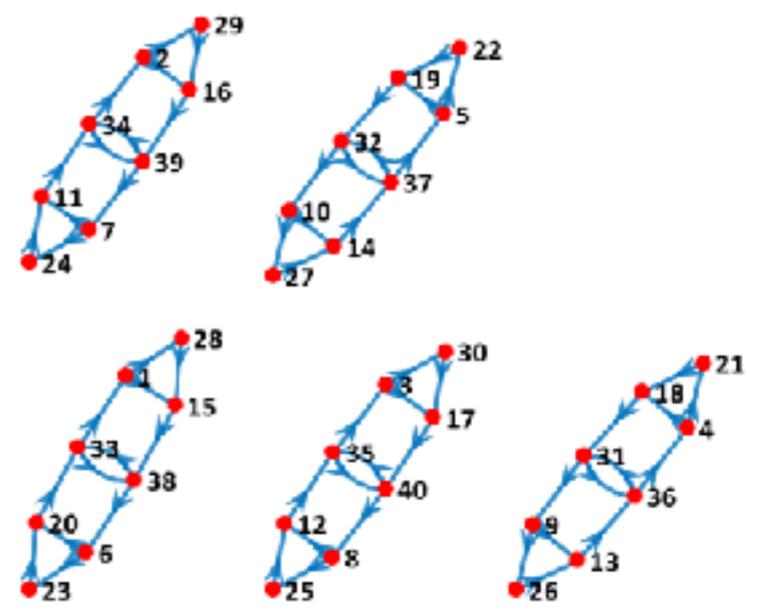


| | |
|--------------------------|----------------------------|
| $\lambda_1 = -1$ | $\lambda_{31} = -1/\tau$ |
| $\lambda_2 = -1$ | $\lambda_{32} = -1/\tau$ |
| $\lambda_3 = -1$ | $\lambda_{33} = 1/\tau$ |
| $\lambda_4 = -1$ | $\lambda_{34} = 1/\tau$ |
| $\lambda_5 = -1$ | $\lambda_{35} = 1/\tau$ |
| $\lambda_6 = -1$ | $\lambda_{36} = \tau$ |
| $\lambda_7 = -1$ | $\lambda_{37} = \tau$ |
| $\lambda_8 = -1$ | $\lambda_{38} = \tau$ |
| $\lambda_9 = -i$ | $\lambda_{39} = 1/\tau^2$ |
| $\lambda_{10} = -i$ | $\lambda_{40} = \tau^2$ |
| $\lambda_{11} = -i$ | $\lambda_{31} = -\omega$ |
| $\lambda_{12} = -i$ | $\lambda_{32} = -\omega$ |
| $\lambda_{13} = i$ | $\lambda_{33} = -\omega$ |
| $\lambda_{14} = i$ | $\lambda_{34} = -\omega$ |
| $\lambda_{15} = i$ | $\lambda_{35} = -\omega$ |
| $\lambda_{16} = i$ | $\lambda_{36} = -\omega^2$ |
| $\lambda_{17} = -\tau$ | $\lambda_{37} = -\omega^2$ |
| $\lambda_{18} = -\tau$ | $\lambda_{38} = -\omega^2$ |
| $\lambda_{19} = -\tau$ | $\lambda_{39} = -\omega^2$ |
| $\lambda_{20} = -1/\tau$ | $\lambda_{40} = -\omega^2$ |

$$\zeta(z) = -\frac{(z+1)(z^2+z-1)^2}{(z-1)(-z^2+z+1)^3(z^2-3z+1)}$$

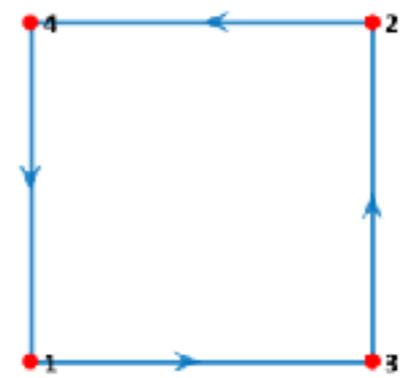
Inflation for Faces, $H^{(2)} = \mathbb{Z}^8$

| | |
|---------------------|----------------------------|
| $\lambda_1 = -1$ | $\lambda_{21} = -1/\tau$ |
| $\lambda_2 = -1$ | $\lambda_{22} = -1/\tau$ |
| $\lambda_3 = -1$ | $\lambda_{23} = -1/\tau$ |
| $\lambda_4 = -1$ | $\lambda_{24} = -1/\tau$ |
| $\lambda_5 = -1$ | $\lambda_{25} = -1/\tau$ |
| $\lambda_6 = -1$ | $\lambda_{26} = \tau$ |
| $\lambda_7 = -1$ | $\lambda_{27} = \tau$ |
| $\lambda_8 = -1$ | $\lambda_{28} = \tau$ |
| $\lambda_9 = -1$ | $\lambda_{29} = \tau$ |
| $\lambda_{10} = -1$ | $\lambda_{30} = \tau$ |
| $\lambda_{11} = -i$ | $\lambda_{31} = -\omega$ |
| $\lambda_{12} = -i$ | $\lambda_{32} = -\omega$ |
| $\lambda_{13} = -i$ | $\lambda_{33} = -\omega$ |
| $\lambda_{14} = -i$ | $\lambda_{34} = -\omega$ |
| $\lambda_{15} = -i$ | $\lambda_{35} = -\omega$ |
| $\lambda_{16} = i$ | $\lambda_{36} = -\omega^2$ |
| $\lambda_{17} = i$ | $\lambda_{37} = -\omega^2$ |
| $\lambda_{18} = i$ | $\lambda_{38} = -\omega^2$ |
| $\lambda_{19} = i$ | $\lambda_{39} = -\omega^2$ |
| $\lambda_{20} = i$ | $\lambda_{40} = -\omega^2$ |



Inflation for Edges, $H^{(1)} = \mathbb{Z}^5$

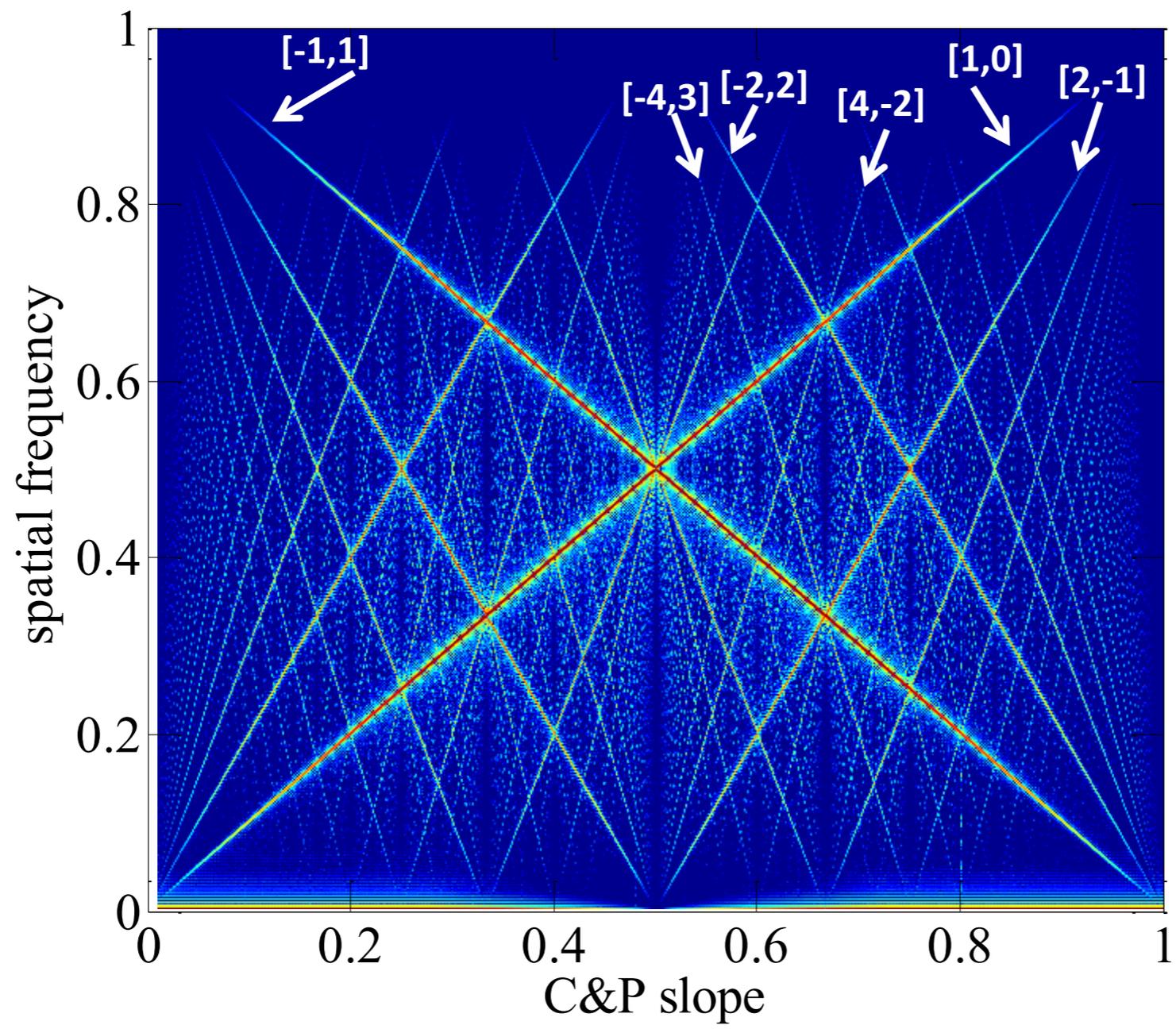
| |
|------------------|
| $\lambda_1 = -1$ |
| $\lambda_2 = 1$ |
| $\lambda_3 = -i$ |
| $\lambda_4 = i$ |

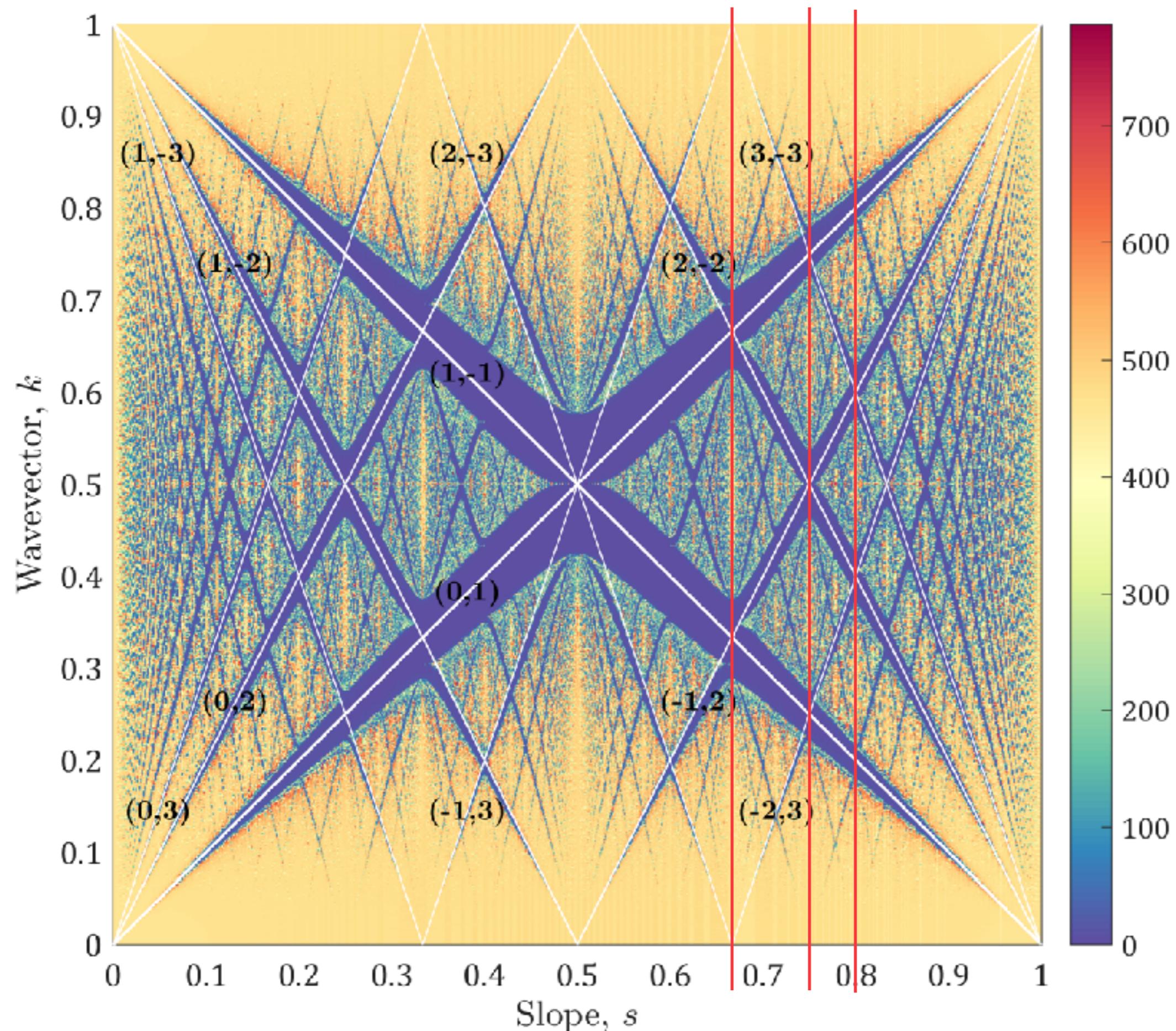


Inflation for Vertices, $H^{(0)} = \mathbb{Z}^1$

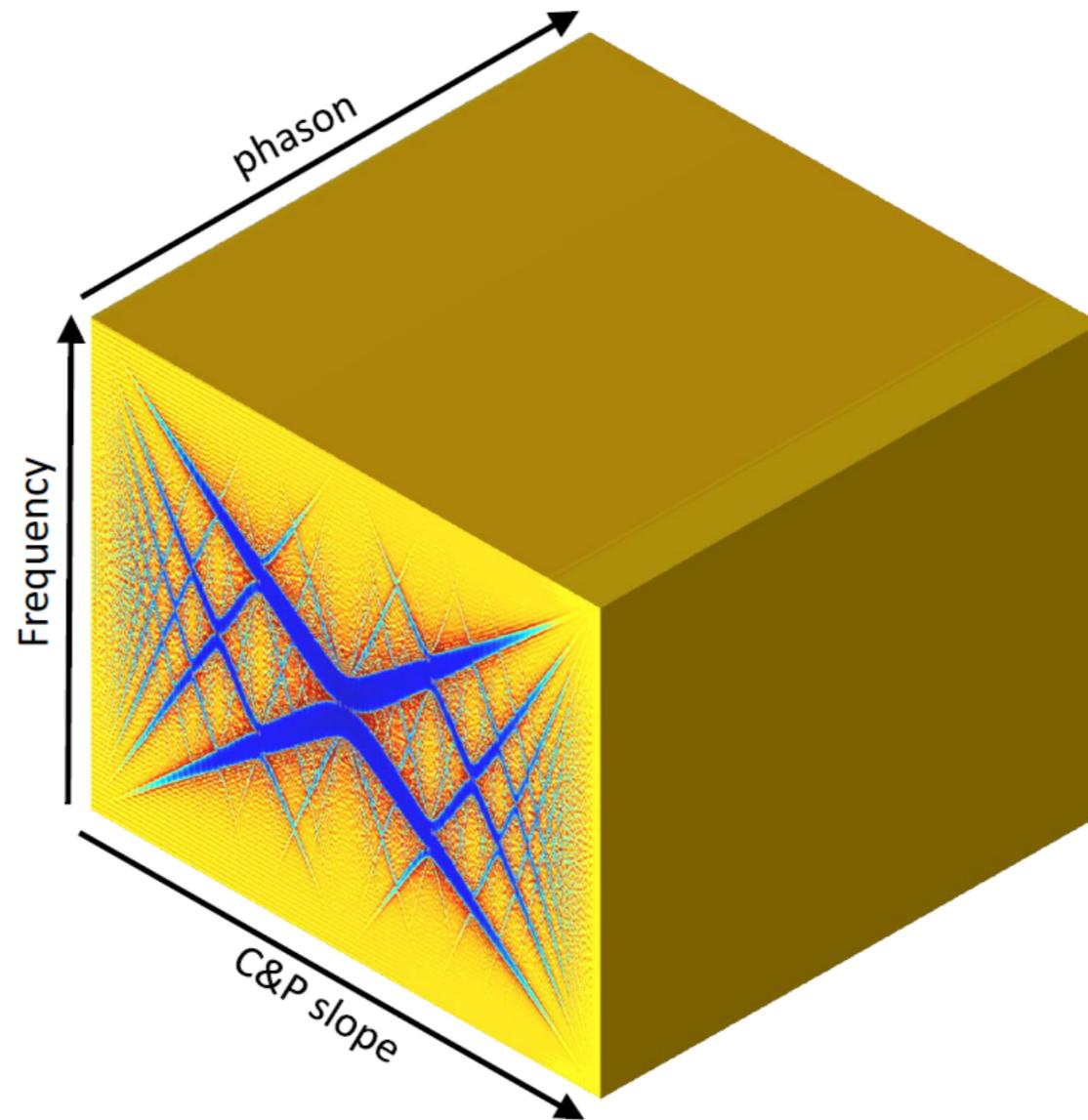
Summary - the untold part

- Given a topological meaning to the integers labelling the gaps of the fractal spectrum.
- Proposed a complete algebraic structure to account for the topological integers (Abelian group structure isomorphic to $\mathbb{Z}/F_N\mathbb{Z}$)
- Generalisation to other substitutions
- Generalized Cut&Project

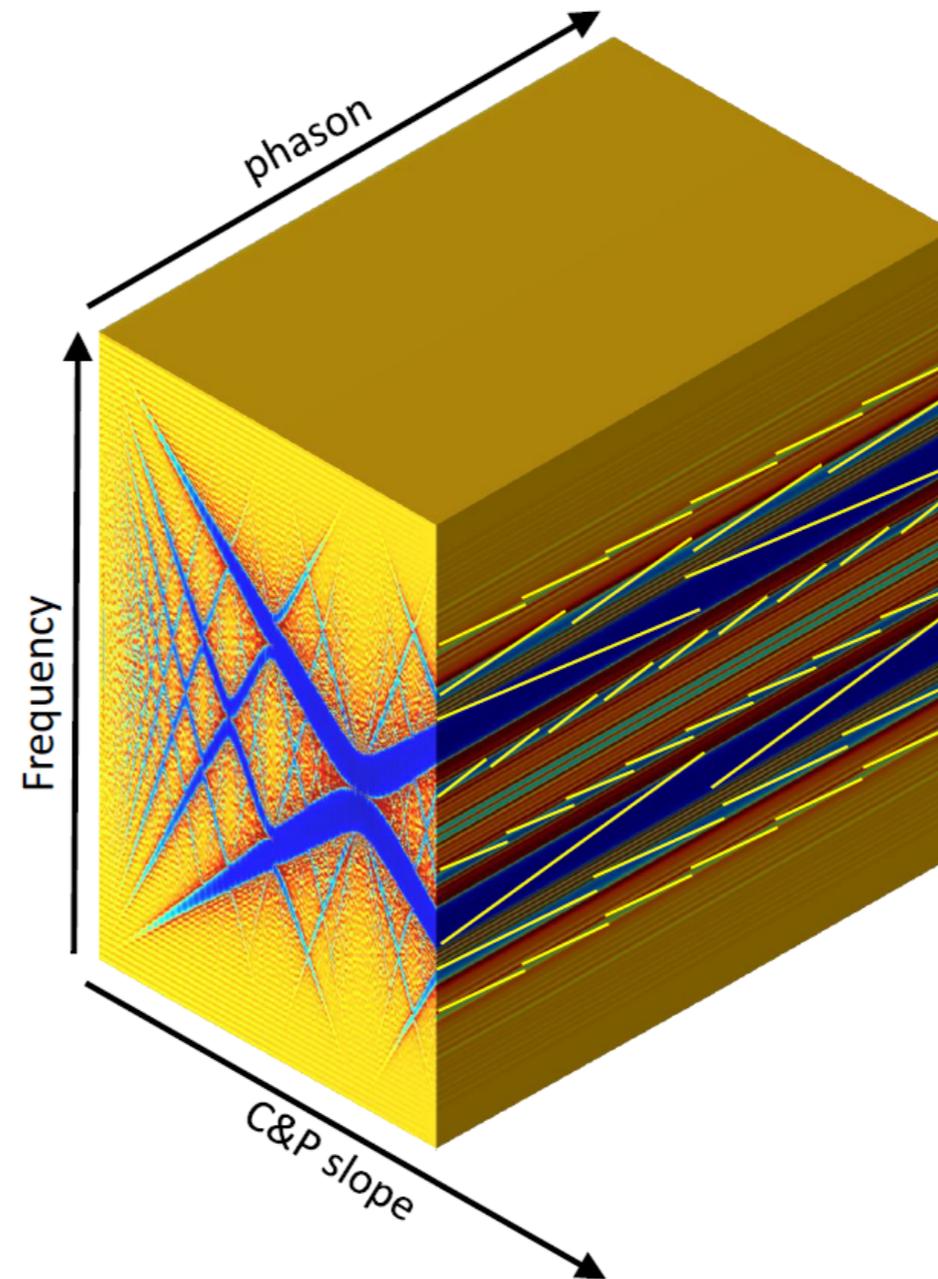




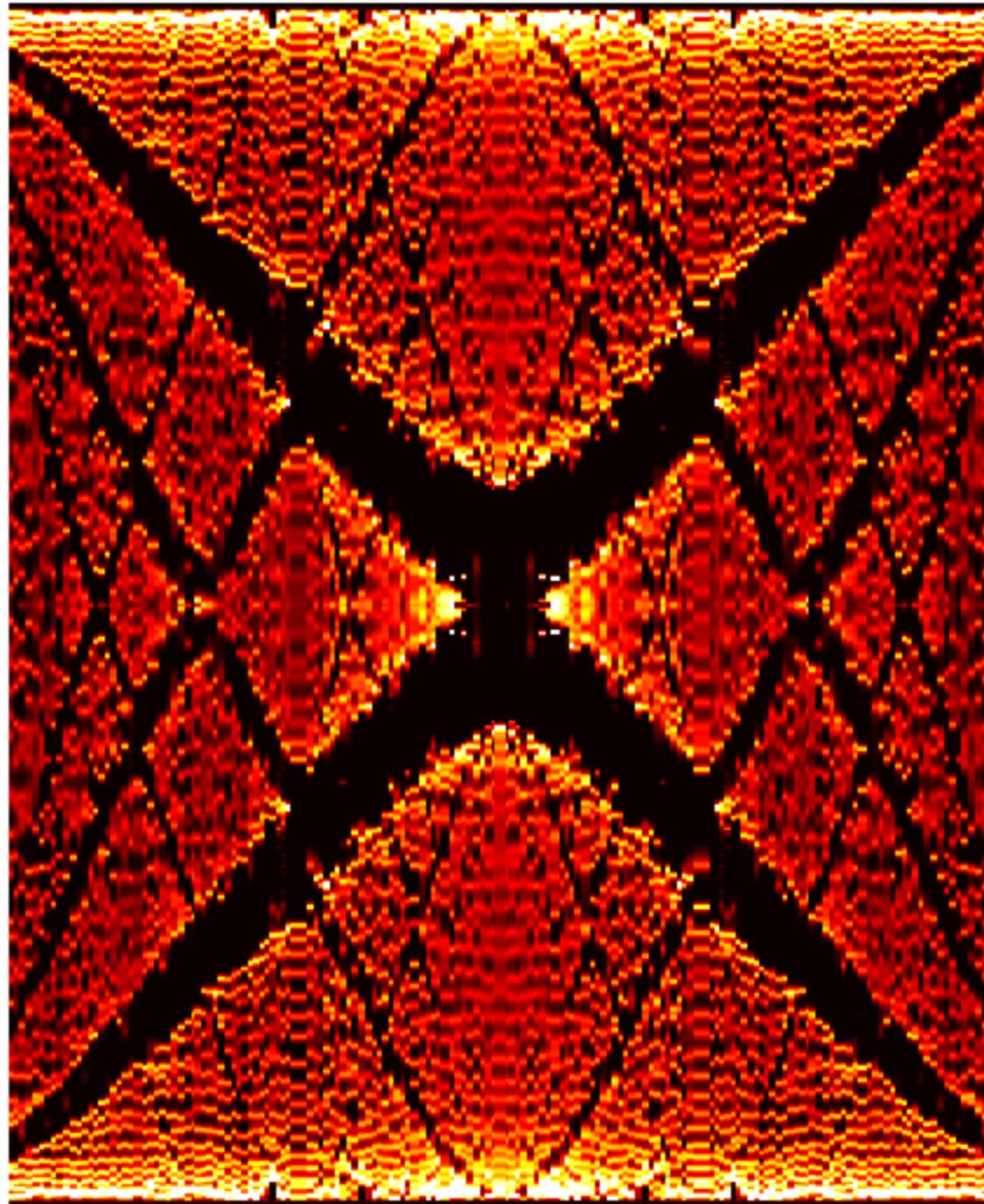
- The λ_1 is fixed per substitution
- How to change λ_1 continuously? By **Cut and Project**



- The λ_1 is fixed per substitution
- How to change λ_1 continuously? By **Cut and Project**



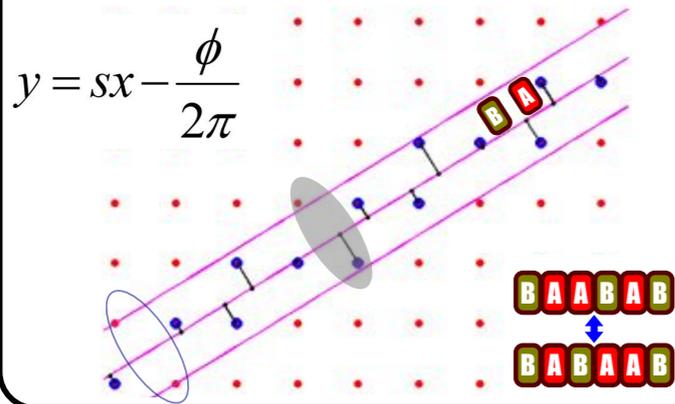
Wannier Butterfly



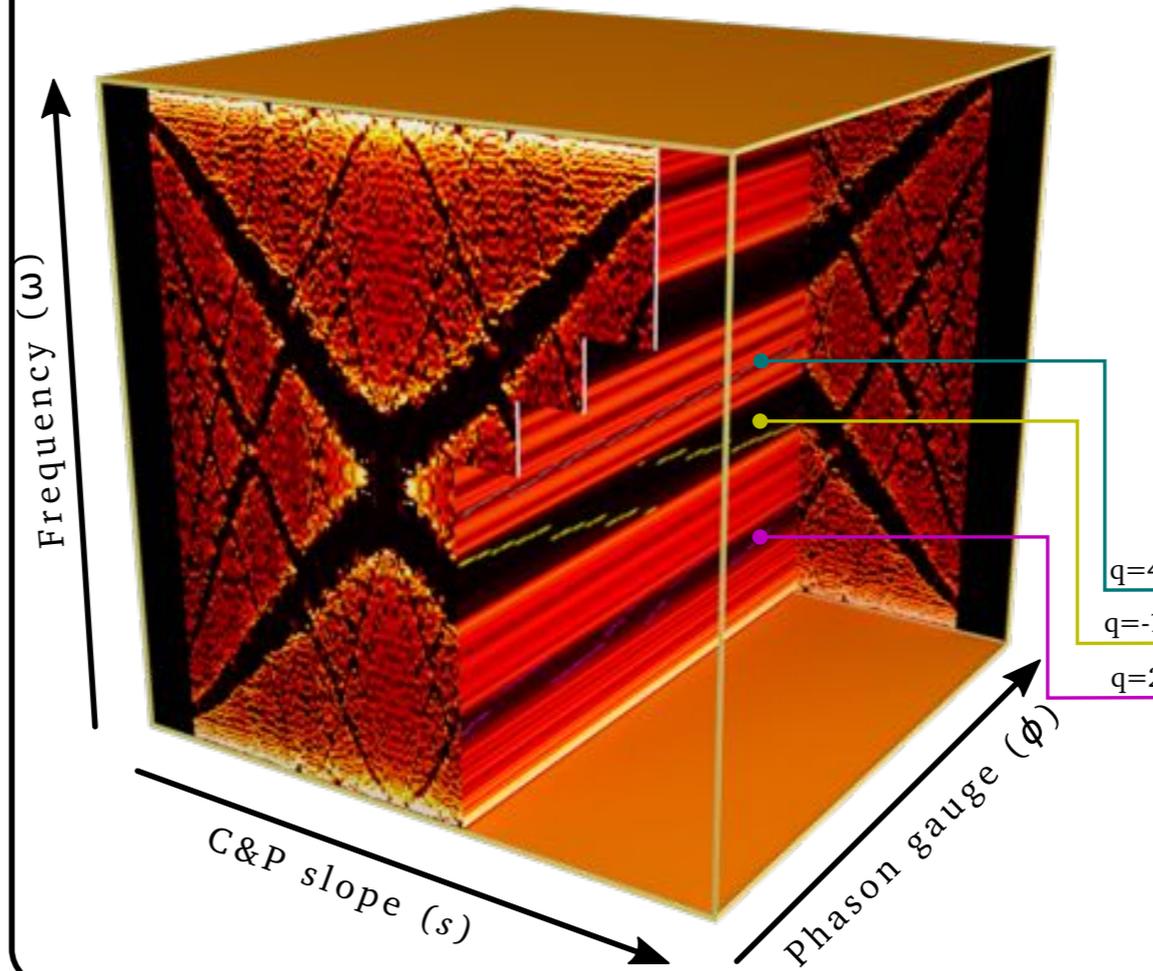
E. Levy, F. Mortessagne, U. Kuhl, E.A, 2016

Wannier Butterfly

Cut & Paste



Phase space

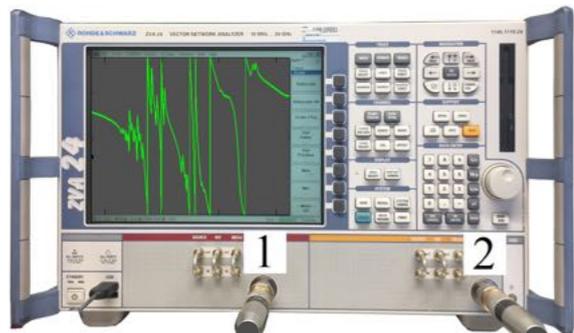


Gap labeling theorem

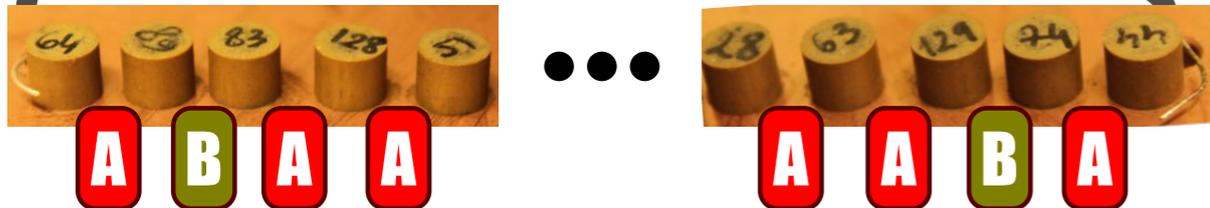
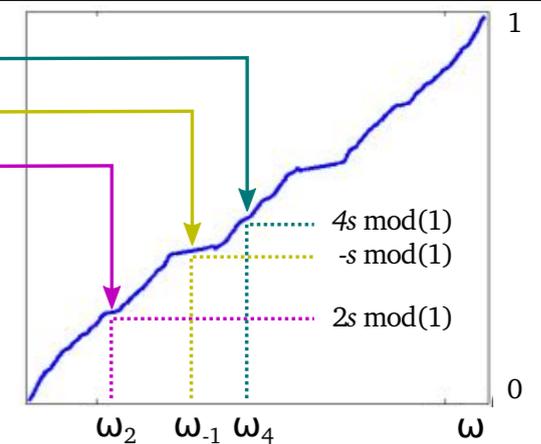
$$IDOS(\omega_q) = \frac{1}{\pi} \delta(\omega_q) = qs \pmod{1}$$

$q \in \mathbb{Z}$

Scattering phases: $\delta(\omega)$, $\alpha(\omega)$



IDOS (ω)



Topological winding

$$\mathcal{W}(\alpha_q) = \frac{1}{2\pi} \int_0^{2\pi} d\phi \frac{d\alpha_q(\phi)}{d\phi} = 2q$$

Ambiguity for winding numbers

